Solid State NMR

- AVANCE Solids

User Manual

Version 003

Innovation with Integrity

NMR

This manual was written by

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Document Number: Z4D10641B

P/N: Z31848

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1 Introduction

This manual is intended to help users set up a variety of different experiments that have become more or less standard in solid state NMR.

In previous versions of this manual, the hardware was described in some detail, and also basic setup procedures. Armed with this knowledge, it was assumed the users would be in a position to manage the setup of even complicated experiments themselves.

In this version, the hardware is not discussed in detail, since there is no longer much hardware which is specific to solid-state NMR. There are still transmitters with higher power, and preamps and probes that take this power, but for the purposes of experimental setup, detailed knowledge is not required, since the setup does not generally depend on the details of the hardware. Thus this manual is now much more specific to the type of experiment which is to be executed, and includes advice on how to set the experiment up properly for best performance. If any special hardware (or software) knowledge is required, it is included within the experiment details.

The manual covers some of the most frequently used solid-state NMR experiments. The manual was written primarily for use with Bruker AVANCE III instruments, but the experiments are identical, or similar, for AVANCE I and AVANCE II instruments. For example, pulse programs will have slightly different names, differing usually in the pulse program name extension. If you do not find the experiment/pulse program that you are looking for, contact your nearest applications scientist. Users with older instruments (DSX, DMX, DRX) should refer to the Solids User Manual delivered within the Help system at **Help** | **Other topics** | **Solids User Manual**. Even though the pulse programs may look similar, they will not run on these instruments.

The first five chapters deal with basic setup procedures, subsequent chapters are dedicated to specific types of experiments. There may be many different „sub“ experiments within a given type, since the same information can often be obtained with pulse sequences differing by subunits only, or in using a totally different principle. The experiments outlined here are usually the most important ones and/or the ones that were common at the time when the manual was written.

The manual consists of largely self-contained units rather than being a comprehensive single volume. This was done in order to be more flexible in updating/replacing individual chapters. The individual chapters are written by different people, so there may be some differences in style and composition.

Note Concerning Future TopSpin Releases

TopSpin is continuously under development, thus some user interfaces, settings and routines described in this manual will differ from the latest TopSpin version.

1.1 Disclaimer

Any hardware units mentioned in this manual should only be used for their intended purpose as described in their respective manual. Use of units for any purpose other than that for which they are intended is taken only at the users own risk and invalidates any and all manufacturer warranties.

Service or maintenance work on the units must be carried out by qualified personnel.

Only those persons schooled in the operation of the units should operate the units.

Read the appropriate user manuals before operating any of the units mentioned. Pay particular attention to any safety related information.

1.2 Safety Issues

Please refer to the corresponding user manuals for any hardware mentioned in this manual for relevant safety information.

2 Test Samples

Nucleus	Sample	Method	O1P	Remarks
^{3}H				
^{1}H	Silicone paste	1HMAS	0	Setup proton channel, shim, set field.
	Silicone rubber	1HMAS	0	Setup proton channel, set field.
	Adamantane	1HMAS	0	Setup proton channel, set field, shim under CRAMPS conditions.
	Glycine	CRAMPS	-3	Setup CRAMPS.
	Malonic Acid	CRAMPS	-3	Resolution CRAMPS, d1=60 seconds.
^{19}F	PVDF	19FMAS	106	Direct observe ^{19}F.
		CP		CP $^{1}H/^{19}F$, $^{1}H/^{13}C$, $^{19}F/^{13}C$ (low sensitivity).
	PTFE	19FMAS	126	Direct observe.
^{3}He				
$^{203,209}Tl$				
^{31}P	$(NH_4)H_2PO_4$	$^{1}H/^{31}PCP$	0	Powdered sample, piezoelectric, 4 seconds.
^{7}Li	LiCl	MAS		
$^{117,119}Sn$	Sn (cyclohexyl)$_4$	CP		5 milliseconds contact, d1>10 seconds.
	$Sm_2Sn_2O_7/SnO_2$	MAS		VT shift thermometer, d1<1 seconds.
				$Sm_2Sn_2O_7$,>60 seconds SnO_2 (temp. independent).
^{87}Rb	$RbNO_3$, $RbClO_4$	MQMAS	0	0.5 seconds repetition.
^{11}B	BN	MAS		
	Boric Acid	MQMAS		>5 seconds repetition.
^{65}Cu	Cu-metal powder	wideline		Knight shift +2500 ppm.
^{71}Ga	Ga_2O_3	hahn echo		CT 300 kHz wide.
^{129}Xe	as hydroquinon Clathrate	CPMAS	0	d1>5 seconds.
	gas in air		0	Single pulses overnight, 1 seconds.
^{23}Na	Na_2HPO_4	MQMAS	0	Dep. on crystal water 2-5 lines.
	$Na_3P_3O_9$	MQMAS		
^{51}V	NH_4VO_4			
^{123}Te				
^{27}Al	AlPO-14	MQMAS	0	d1 05-1seconds, 4 lines.
^{13}C	Adamantane	CP, DEC	50	HH setup, shim.
	α-glycine	CP	110	Sensitivity, decoupling. Prep.: precipitate with acetone from aq. solution, C, N fully.
				Labelled for fast setup, recoupling, REDOR (10% in natrl. abundance).

^{79}Br	KBr	MAS	57	d1< 50 milliseconds, angle setting. Finely powdered, reduced volume.
^{59}Co	$Co(CN)_6$	MAS		Shift thermometer.
^{55}Mn	$KMnO_4$	MAS		>500 kHz pattern.
^{93}Nb				
^{207}Pb	$PbNO_3$ $Pb(p\text{-}tolyl)_4$	MAS CP	 -150	Shift thermometer, 0.753 ppm/degr. d1>10 seconds. 5 milliseconds, 15 seconds.
^{29}Si	Q_8M_8 DSS,TMSS	CPMAS CPMAS	-50 0	d1>5 seconds, reference sample 12.6/-108 ppm. Reference sample 0 ppm.
^{77}Se	H_3SeO_3 $(NH_4)_2SeO_4$	CPMAS CPMAS	1800 -200	HH setup, 8 milliseconds contact, d1>10 seconds. 3 milliseconds, d1>4 seconds.
^{113}Cd	$Cd(NO_3)_2*4H_2O$	CPMAS	350	15 milliseconds contact, d1>8 seconds.
^{195}Pt	$K_2Pt(OH)_6$	CPMAS	-12000	1 milliseconds contact, d1>4 seconds.
^{199}Hg	$Hg(acetate)_2$ Hexakis (dimethyl sulphoxide)Hg(II) trifluoromethansulfonate	CPMAS	2500-2313	5 milliseconds contact, d1>10 seconds 30-35 milliseconds contact, d1>10 seconds *.
^{2}H	d-PMMA d-PE $d\text{-}DMSO_2$	WL WL WL	0 0 0	Wideline setup d1 5 seconds. Wideline setup d1 0.5 seconds/10 seconds amorphous/crystalline exchange expt. at 315K.
^{6}Li	LiCl, Li (org.)			Make sure it is not 6Li depleted, d1>60 seconds.
^{17}O	D_2O		0	Pulse determination, 100 scans, 0.5 seconds.
^{15}N	α-glycine	CP	50	Sensitivity, 4 milliseconds contact, 4 seconds. Labelled for fast setup.
^{35}Cl	KCl	WL, MAS	0	Pulse determination, 100 scans.
^{33}S	K_2S	MAS	0	100 scans in a >=500 MHz instr.
^{14}N	NH_4Cl	MAS, WL	0	100 scans, narrow line.
$^{47/49}Ti$	Anatas	MAS		
^{39}K	KCl	MAS, WL	0	100 scans.
^{109}Ag	$AgNO_3$ $AgSO_3CH_3$	MAS CPMAS	 70	1scan, 500 seconds, finely powdered. 50 milliseconds contact, 10 seconds repetition, 1 scan.
^{89}Y	$Y(NO_3)_3*6H_2O$	CPMAS	-50	10 milliseconds contact, d1>10 seconds.
* Literature: J.M. Hook, P.A.W. Dean and L.C.M. van Gorkom, Magnetic Resonance in Chemistry, 33, 77 (1995).				

Table 2.1: Setup Samples for Different NMR Sensitive Nuclei

3 General Hardware Setup

Avance instruments are constructed in a way to minimize the requirements to reconnect or readjust hardware for different experiments. Probe changes may sometimes be necessary, and require some manual operations. This chapter deals with connections that need to be done by the operator, and also with other manipulations that are required to set up the instrument in an optimum way.

Since the RF pathways are under software control up to the preamplifier, and under operator control between the preamplifier and probe, both setups are considered separately.

All remaining connections (heater cable, thermocouple, gas flow, spin rate cable, PICS cable) are not under software control, so the operator is responsible for proper wiring, cabling, and tubing! Since mistakes (especially in connection with compressed gas tubing) may cause rather expensive repairs, it is recommended to check the connection carefully before an experiment is started.

The following operations will be described and illustrated with suitable images, for WB and SB probes, when non-trivial differences exist.

3.1 Connections to the Preamplifier

For solids and liquids there should normally be different sets of preamplifiers. Liquids preamplifiers (HPPR, High Performance Preamplifiers) are not suitable for some of the requirements of solid state NMR. When CP/MAS applications are the only solids applications, it is possible however to use liquids preamplifiers for X-observation. Solids preamplifiers (HPHPPR, High Power High Performance Preamplifiers) are definitely required if high power ≥ 1 kW is used (liquids preamplifiers take max. 500W for X frequencies, 50W for proton and fluorine frequency). For the high frequency range ^{19}F and ^{1}H, two different types of solids preamps are available, the older HPHPPR ^{19}F /^{1}H and the recent replacement HPLNA (High Power Low Noise Amplifier) which is strictly frequency selective, either ^{19}F or ^{1}H.

The connections on the back) of the preamp stack should normally not be changed. For broadband high power preamplifiers, it is important to insert the appropriate matching box into the side of the preamp.

Figure 3.1: All Connections to the Back of the Preamplifier

The figure above shows RF cables from the transmitter, RS-485 control, DC voltages in, tune and lock RF in, RF signal out to receiver, and gate pulses for preamplifier control (multi-receive setup only). The orange colored cable is the high voltage supply for the HPLNA preamplifier.

Figure 3.2: Transmitter Cables (only) Wired to Back of the Preamplifier

In the figure above the lock preamplifier is located at the bottom of the stack; the transmitter cable carries the lock pulses. For solids, this preamp is normally not required. When the transmitter cables are rewired to different preamp modules, the changes must be entered into the edasp routing (type **edasp setpreamp**, NMR Administration password required).

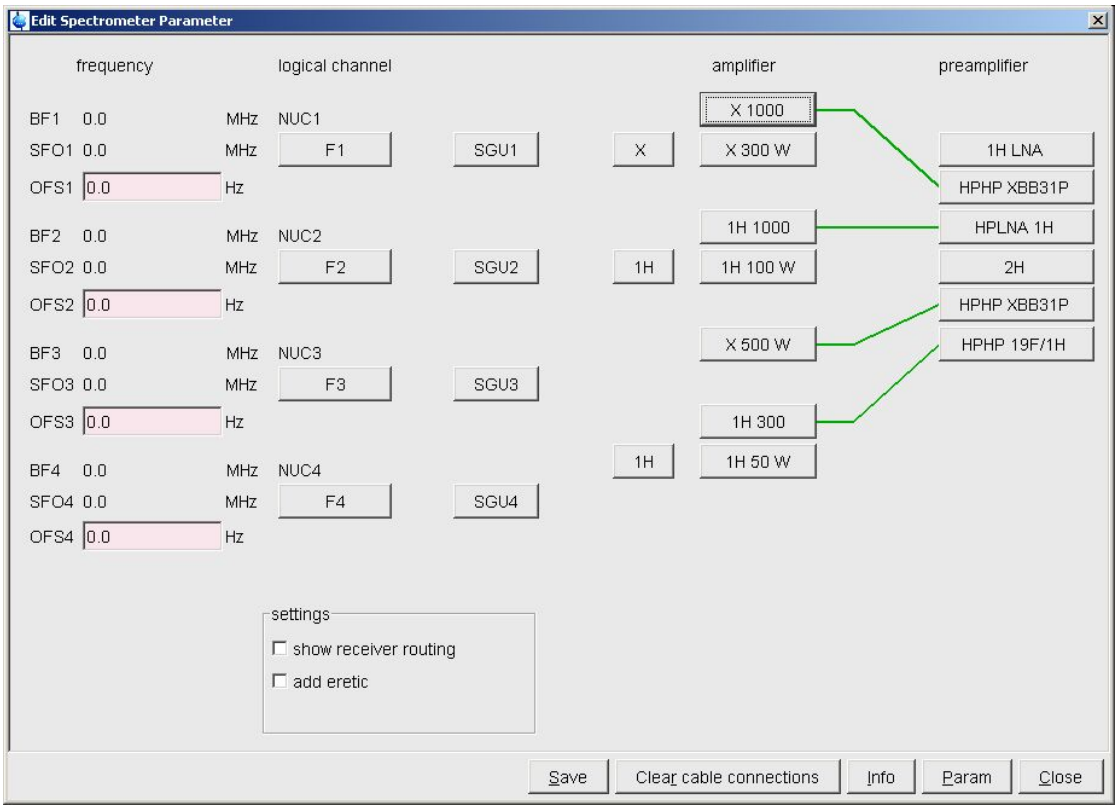


Figure 3.3: The edasp setpreamp Display

Note for the figure above: The transmitter to preamplifier wiring must reflect hardware connections!

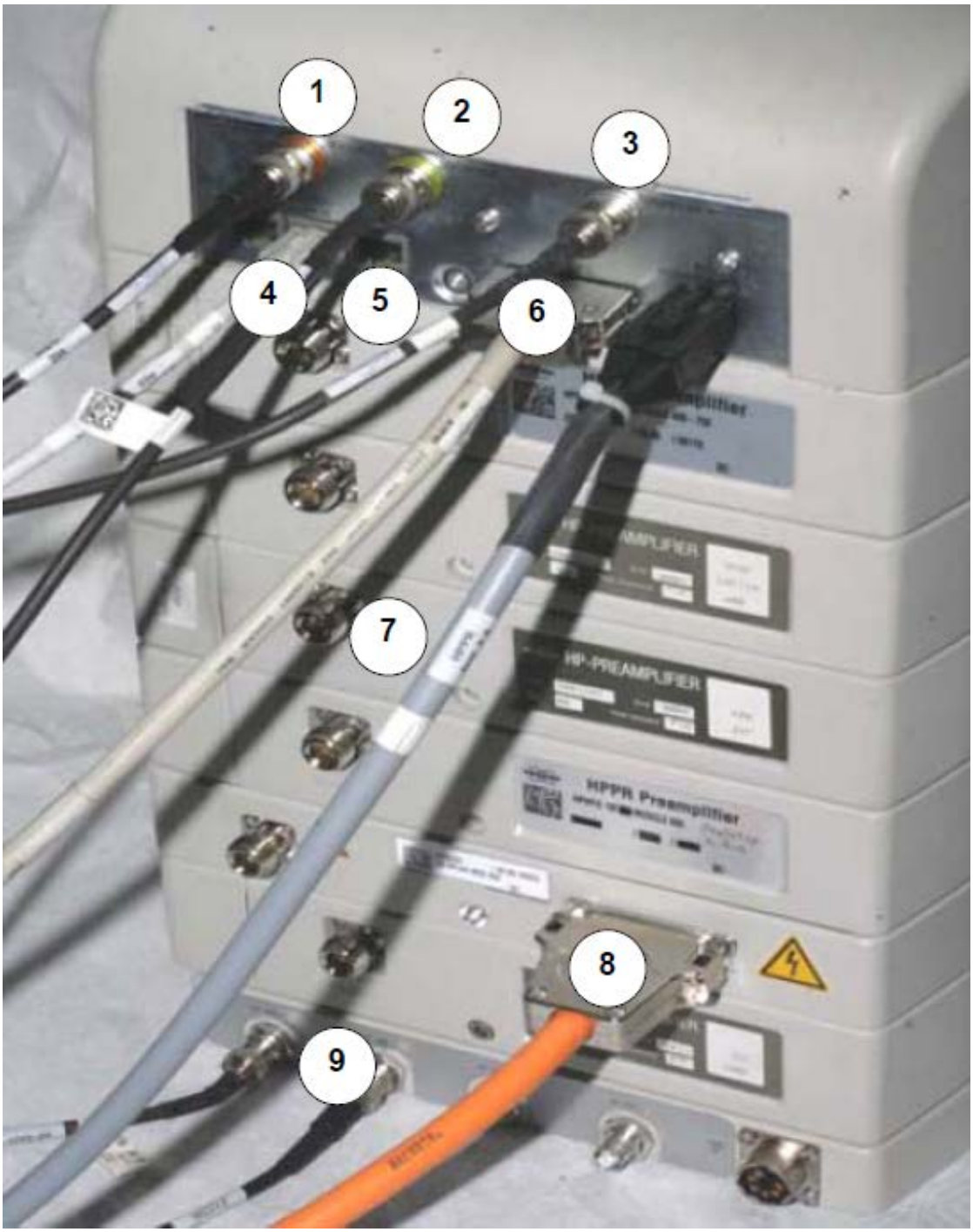


Figure 3.4: Additional Connections to the Preamplifier Stack

1	RF signal out to receiver	6	RS 485 control connection and DCin
2	Lock signal out to lock receiver	7	Additional DC supply for >3 preamps
3	Tune RF in (from SGU 2 aux out)	8	High voltage DC for HPLNA-preamp
4	PICS probe ID cable to probe	9	Additional controls for multi-receiver
5	ATMA and AUX connectors		

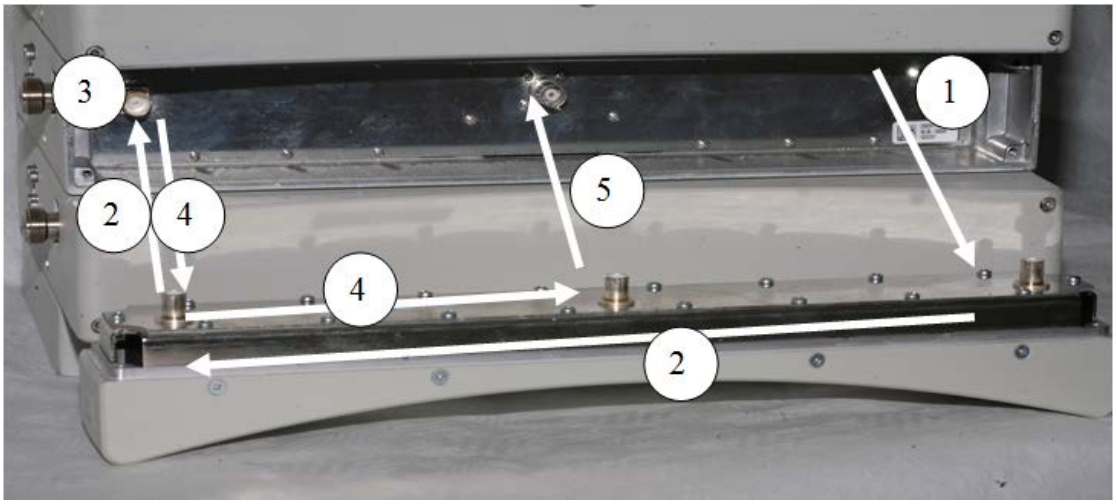


Figure 3.5: Matching Box Setup for High Power X-BB Preamplifiers

1	Pulse from TX	4	Signal from probe
2	Pulse pathway to probe	5	Signal to preamplifier
3	To probe		

The frequency of the observed nucleus must be within the bandwidth of the matching box: The matching box contains a low pass filter to suppress frequencies above the X-nucleus frequency range (^{1}H, ^{19}F) and a passive diode multiplexer, which directs the RF pulse into the probe, and the NMR signal into the preamplifier. High resolution preamplifiers use actively switched pin diodes for this purpose and are therefore broadbanded, so there is no exchangeable box.

Pulsing with high power into an RF circuit which is not properly set up to pass this frequency may result in damage to the RF circuit (in this case, the matching box) or to the transmitter. This applies to filters, preamplifiers and matching boxes. When liquids preamps are used for solids work, power limitations and frequency limitations must be strictly observed!

3.2 RF Connections Between Preamplifier and Probe

These connections must be made using high quality cable, with suitable length. It should be short, but not too short so that the cable is not severely bent. Higher quality cable is fairly stiff; the flexible ones are of less quality.

Note: It is extremely important that RF cables are not bent to a radius of less than 30 cm, and that no force is exerted on the RF connectors. Adapters should be avoided; since every connector may change the impedance to deviate from the required 50 Ω. Cables with loose connectors should be discarded, unless they can be repaired by a skilled RF engineer. BNC connectors should be avoided; they are usually off by 50 Ω.

As pulses in solids NMR can be rather long, and rather high powered, it is also necessary to consider the preamplifier's power limitations.

Proton high resolution preamps are unsuitable for high power pulses, especially for durations required for decoupling. High resolution X-BB preamplifiers are limited to 10 msec pulses at 300 (500) watts. If the back label does not say 500W, it is 300W max.

^{1}H/^{19}F high power preamps do not necessarily need to be bypassed, but may gradually deteriorate under many decoupling pulses. It is therefore recommended to bypass these for decoupling unless the experiment requires that the preamp remain in line.

Note: These preamps are not optimized for ^{19}F, so ^{19}F decoupling should never be done through the preamp. ^{1}H HPLNA preamplifiers need not be bypassed. HPLNA preamplifiers are strictly frequency selective; a ^{19}F pulse through a ^{1}H HPLNA will destroy it!

3.3 RF-Filters in the RF Pathway

RF filters are frequently required if more than one frequency is transmitted to the probe.

Without filtering, the noise and spurious outputs from the transmitter of one channel would severely interact with signal detection on another channel. One has to keep in mind that pulse voltages are in the order of hundreds of volts, but NMR signals are in the order of microvolts. In high resolution, where the selection of nuclei to run is rather limited, it is possible to apply the necessary filtering inside the preamplifier. For solids, this is not so easy, due to the wide range of possible detection frequencies, and to the additional dead time that filters may cause. So all filtering is done with external filters. If a single channel NMR experiment is run, no filters are required.

Usually, one filter per RF channel is required. Both filters should mutually exclude the frequency of the other channel(s). Usual attenuations of the frequency to be suppressed should be around > 80 dB, in special cases, when both frequencies are rather close, >140 dB may be necessary (as in the case $^{1}H/^{19}F$). More than 90 dB is usually hard to achieve with one filter.

Using external filters has three principal safety aspects:

1. Make sure you do not pulse into a filter with a frequency that is not rated for this filter.
2. Make sure the pulse power you apply does not exceed the power rating of this filter. Most modern Bruker filters will survive 1 kW pulses of 5 ms, but older filters (or non-Bruker filters) may not.
3. Remember that filters may attenuate the pulse RF voltage by as much as 1.5 dB (about 20%)!

The following figures illustrate the most common filter combinations.

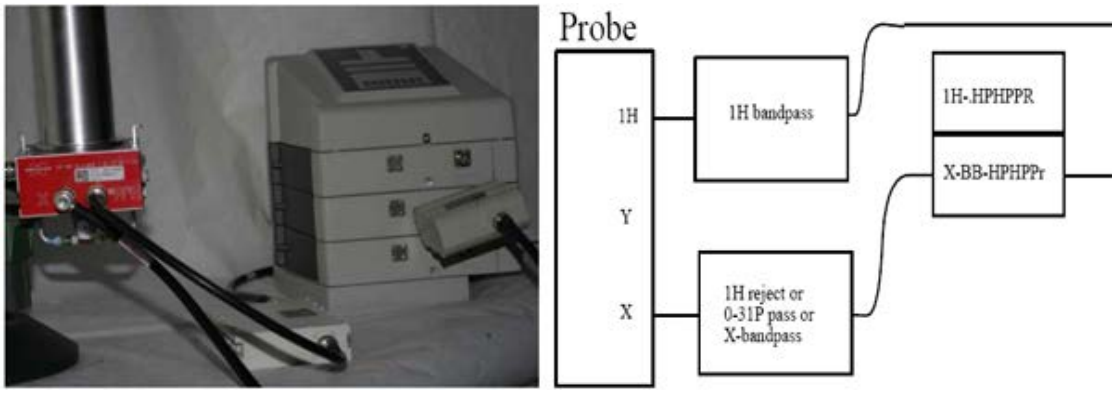


Figure 3.6: Standard Double Resonance CP Experiment, Bypassing the Proton Preamp

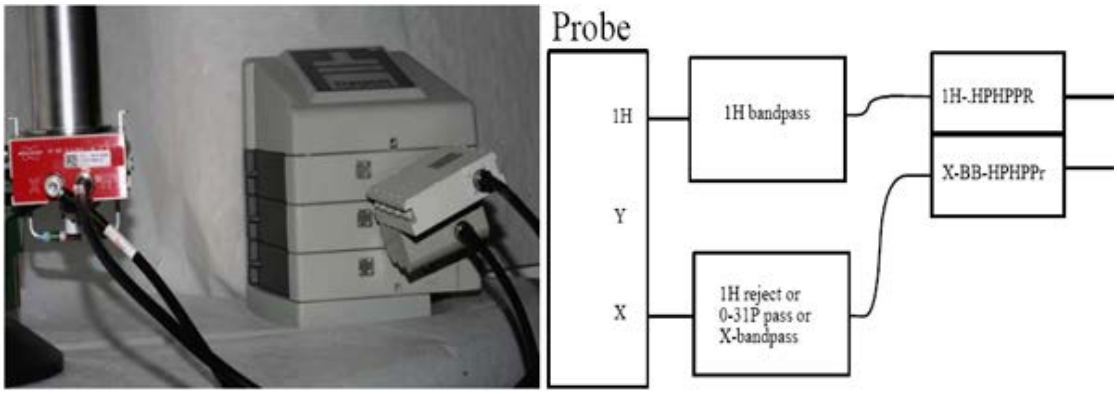


Figure 3.7: Standard CP Experiment, Proton Preamp in Line

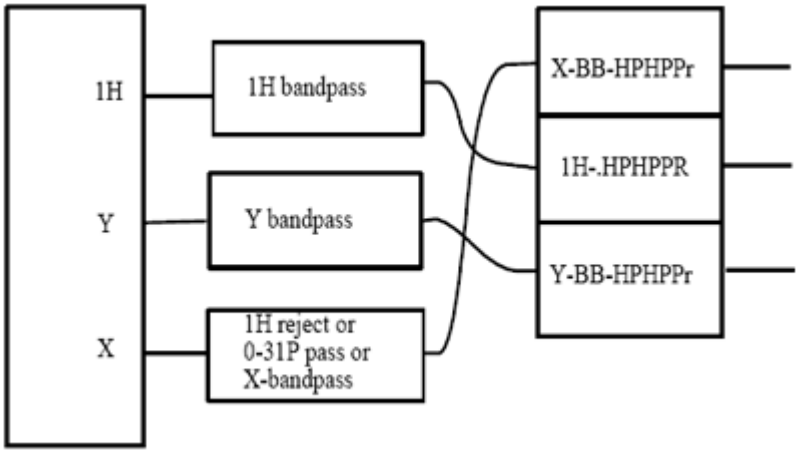


Figure 3.8: Triple Resonance Experiment, without X-Y Decoupling

Please note in the figure above, only high power preamps allow decoupling through the preamp.

In the figure above is a triple resonance experiment, without X-Y decoupling (one bandpass will suffice), note the preamp configuration. It is recommended not to put two preamps of the same kind next to each other in order to avoid incorrect wiring of probe and filters. For the X-channel, only the proton frequency needs to be filtered out if X or Y is not decoupled while Y or X is observed (protons are usually decoupled).

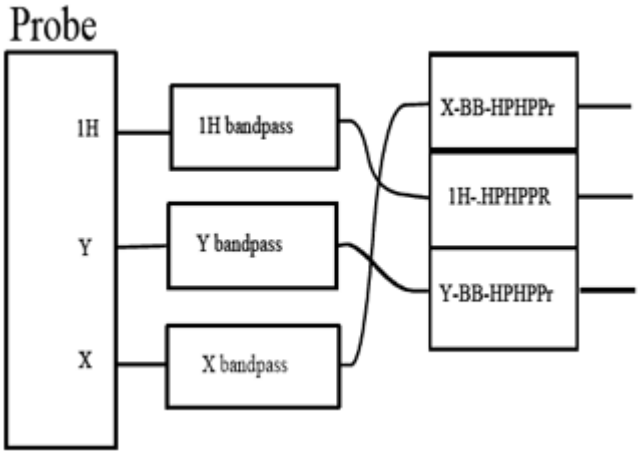


Figure 3.9: Triple Resonance Experiment, with X-Y Decoupling

In the figure above is a triple resonance experiment, with X-Y decoupling (two band passes required! Care should be taken that the two bandpass filters mutually exclude the other frequency as efficiently as possible. Low pass filters will not allow X or Y observe while Y or X is decoupled!).

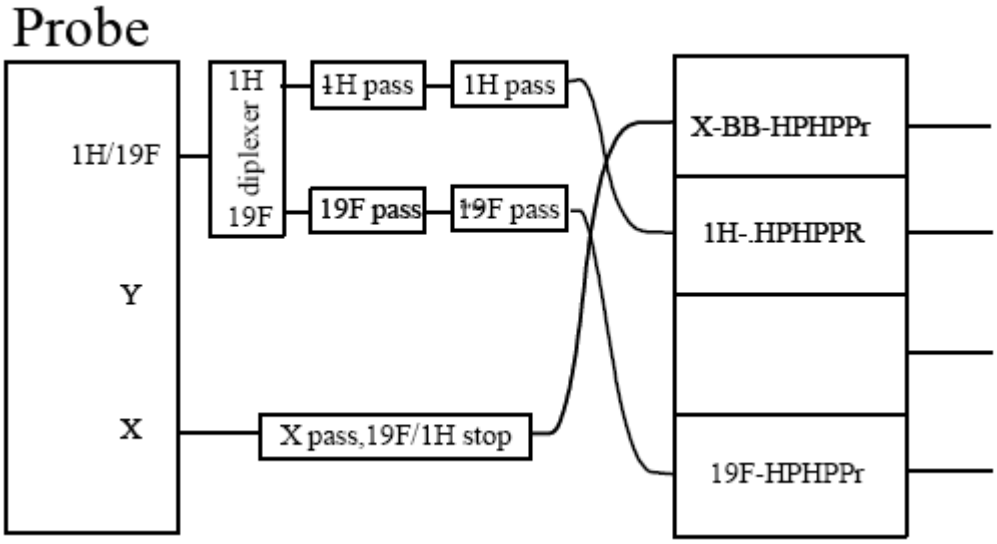


Figure 3.10: Triple Resonance 1H/19F-Experiment

In the figure above is a double/triple resonance HF-experiment, with 19F observation and 1H decoupling or X-observation with 19F and 1H decoupling. This is for WB probes ≥ 400 MHz only! For SB probes and <400 MHz different hardware is used. A set of 1H-transmitter/ bandpass/preamp and 19F-transmitter/bandpass/preamp is required.

Note: A standard 1H/19F preamplifier will not allow long 19F pulses to pass through it, but for short pulses it is okay. For decoupling it must be bypassed, or a dedicated 19F preamp must be used.

Figure 3.11: 19F/1H Combiner/Filter Set

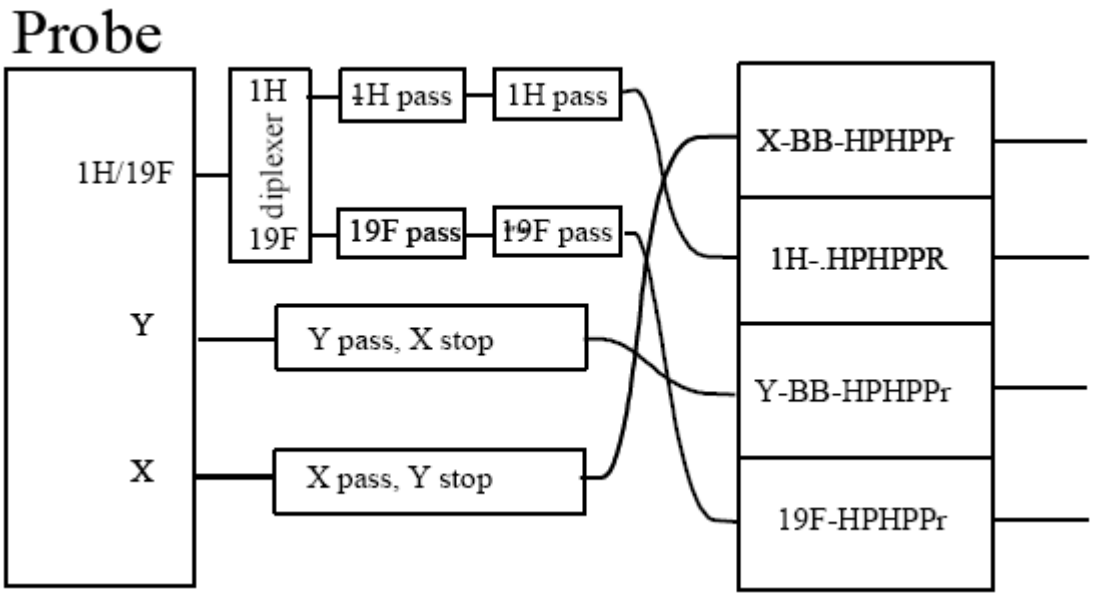


Figure 3.12: Quadruple Resonance HFXY Experiment (WB probes ≥ 400 MHz only!)

3.4 Connections for Probe Identification and Spin Detection

Most older solids probes were delivered without a Probe Identification System (PICS). Probes delivered since 2007 are equipped with PICS. Please refer to *Figure 3.4 [▶ 17]* to identify the PICS port at the preamplifier cover module. The probe connections for the spin rate cable and the PICS cable are shown in the following figures for a WB and for a SB probe.

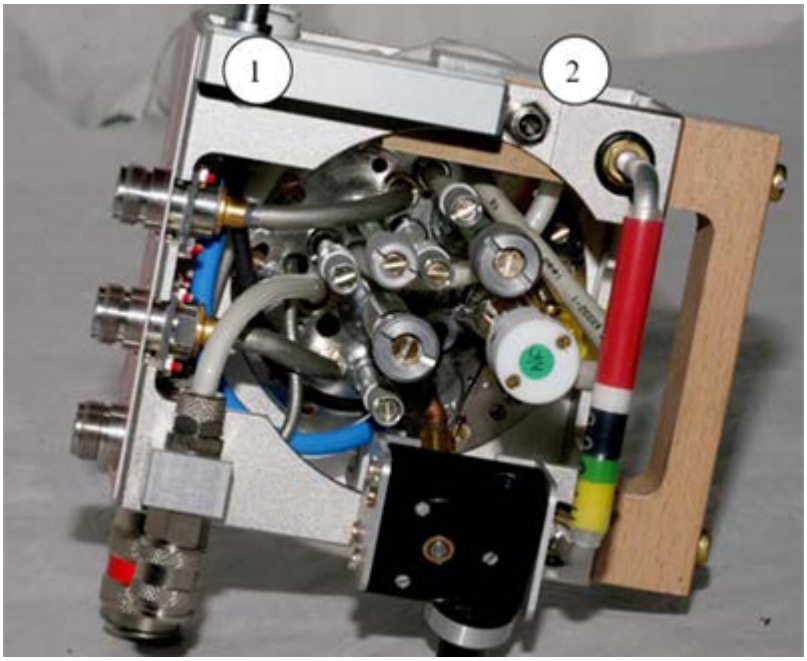


Figure 3.13: PICS Probe Connector and Spin Rate Monitor Cable on a WB Probe

1	PICS probe connector	2	Spin rate monitor cable

On SB probes the location may differ, but the connectors (if present in the case of PICS) can be easily identified by the type of connector.

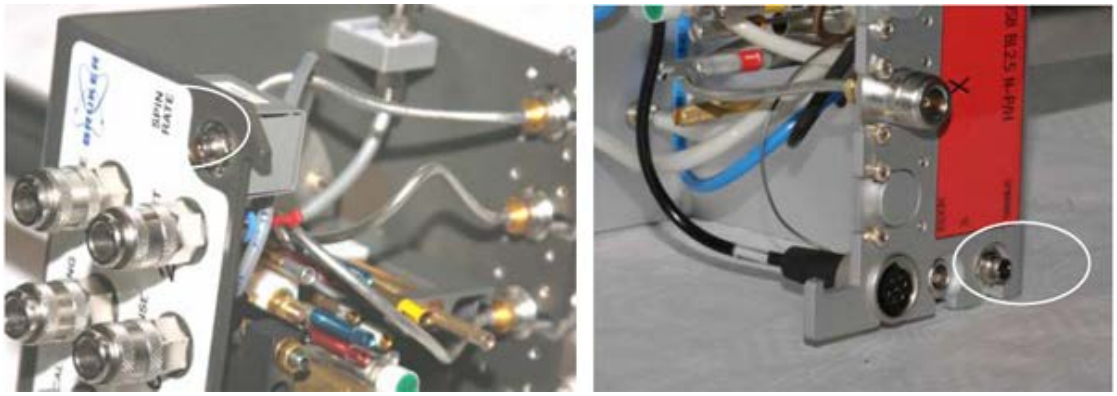


Figure 3.14: Spin Rate Monitor Cable Connector for 2 Different Types of SB Probes

The spin rate monitor cable has a probe side connector that is exactly the same as the power supply cable for the B-TO thermocouple oven used with many high resolution probes. If this cable is connected to the spin rate monitor assembly at the probe, the latter will be destroyed. Make sure the B-TO cable and the MAS cable (labelled “probe” at the probe side) are labelled such that they cannot be mistaken!

3.5 MAS Tubing Connections

For any type of fast spinning probe, compressed gas is used to provide the spinner bearing and drive gas. Please refer to the installation or site planning manuals to learn about the gas requirements. The most important parameters are:

- Mains pressure (should be at least maximum required pressure +1 bar, to provide pressure regulation range).
- Bearing pressure: up to 4.5 bar (as of February 2008)
- Drive pressure: up to 4.5 bar (as of February 2008)

This means that at least 5.5 bars of pressure should be available at the outlet. If the pressure droop along the supply tube is substantial, the internal pressure may drop below 5 bars, whereas the MAS unit stops to regulate and gives a warning.

Consequently, we recommend a primary (inlet) pressure of min. 6 bar, but preferably 7-8 bar (maximum 10 bar) and a low loss gas line (8 mm inner diameter, distance ≤ 5 meters) between the instrument and the gas supply. This will assure trouble free operation even under conditions of high gas throughput. The maximum throughput depends on the experimental conditions and the probe type.

The following gas requirements exist:

1. At room temperature or higher: dew point min. -30 °C, compressed air will do.
2. At temperatures 200 °C or higher (suitable probe required!): nitrogen is required to prevent coil oxidation.
3. At temperatures between room temperature and -50 °C (using a B-CUX cryo cooler with -80 °C exchanger temperature): nitrogen or compressed air with a dew point ≤ -100 °C.
4. At temperatures below -50 °C (using liquid nitrogen and heat exchanger): boil-off nitrogen with a dew point -196 °C.

Any compressed gas used in NMR probes must be free of any liquid droplets or of oil (from compressor lubrication). Oil (even smallest amounts) will especially lead to probe arcing, and/or spinning problems and potentially expensive repair. Using boil-off nitrogen should be carefully considered, since it is by far the most reliable, stable and trouble free source of compressed gas, to be used at any temperature!

3.5.1 Connections

MAS tubing connections are quite different between different types of probes (for stationary, non spinning probes, only frame flush and VT gas are required:

1. WB probes, VTN, WVT and DVT probes (VTN: VT-normal range, WVT: VT-wide range). These probes have a diameter of 72 mm and are longer than SB probes. Probe lengths are the same up to 400 MHz, the same for 500 and 600 MHz, and longer for higher fields.
2. SB probes, VTN and DVT type probes, also major differences between older and more modern probes. Furthermore, probes with sample insert/eject and probes without insert/eject exist.

The major difference between DVT and VTN/WVT probes is that for VTN/WVT probes the bearing gas is used for temperature control, whereas for DVT probes, bearing, drive and VT gas are separate.

3.5.1.1 Wide Bore (WB) Magnet Probes

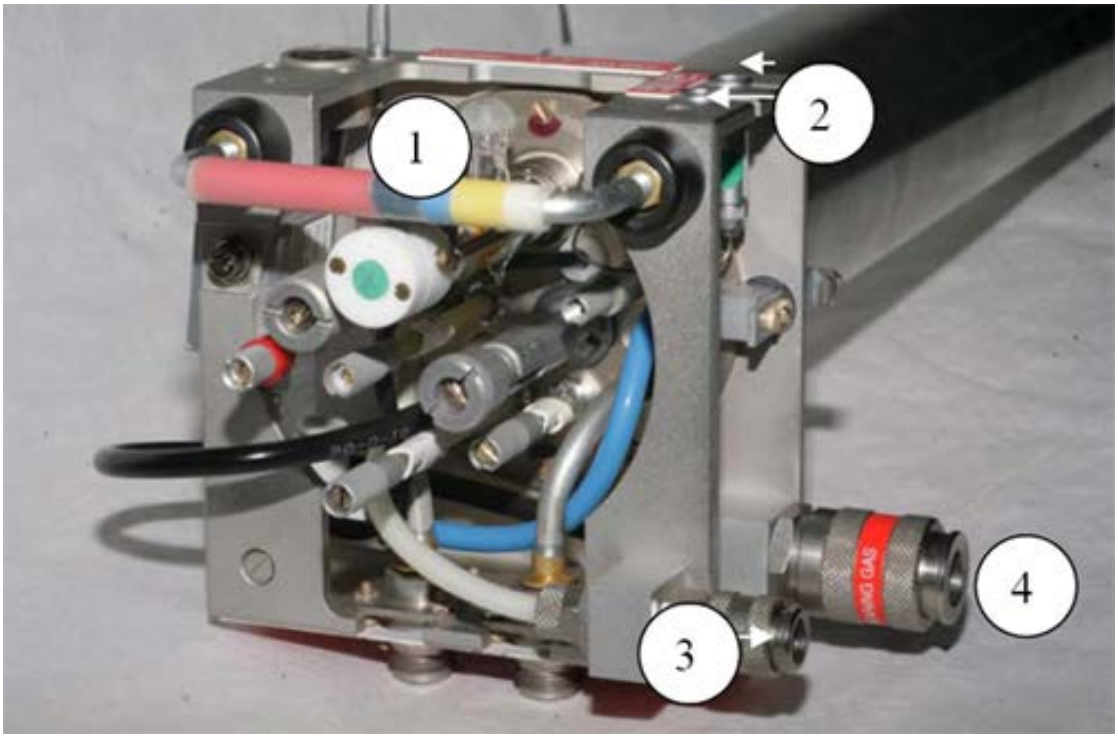


Figure 3.15: WB DVT Probe MAS Tubing Connections

1	VT gas only input into dewar	3	Drive gas in
2	Two thermocouple connectors	4	Bearing gas in

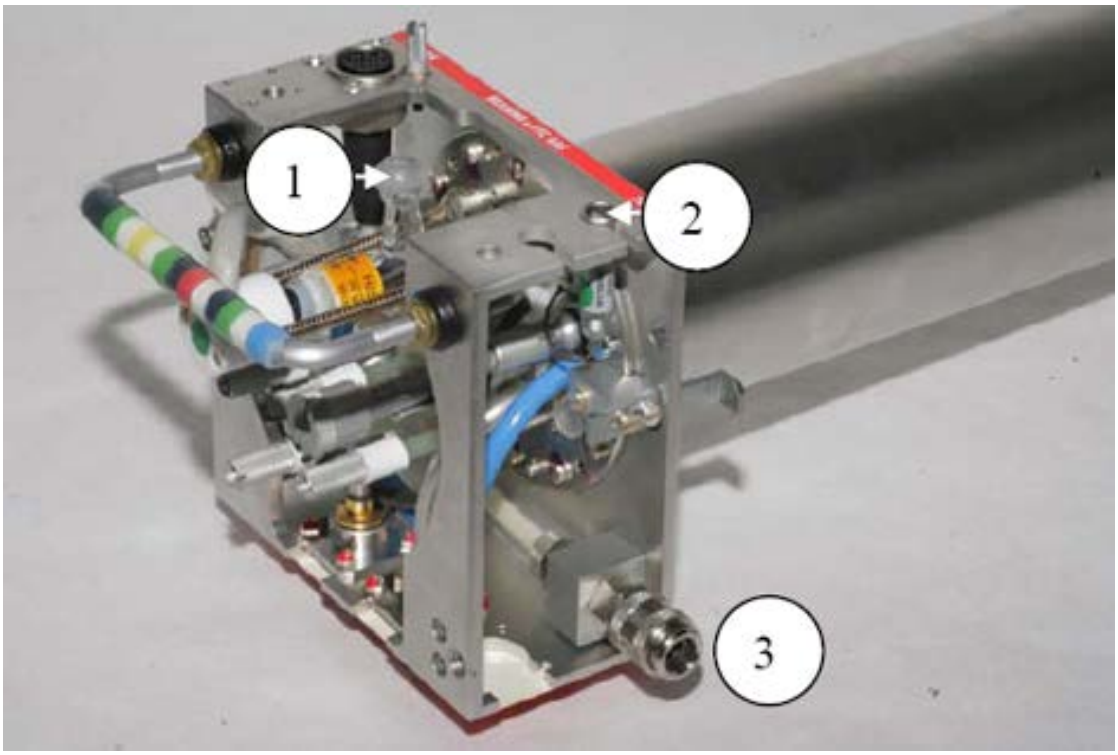


Figure 3.16: VTN Probe MAS Tubing Connections Note: WVT Probes are VTN-Type Probes

1	VT plus bearing gas	3	Drive gas
2	One thermocouple connector at stator inlet		

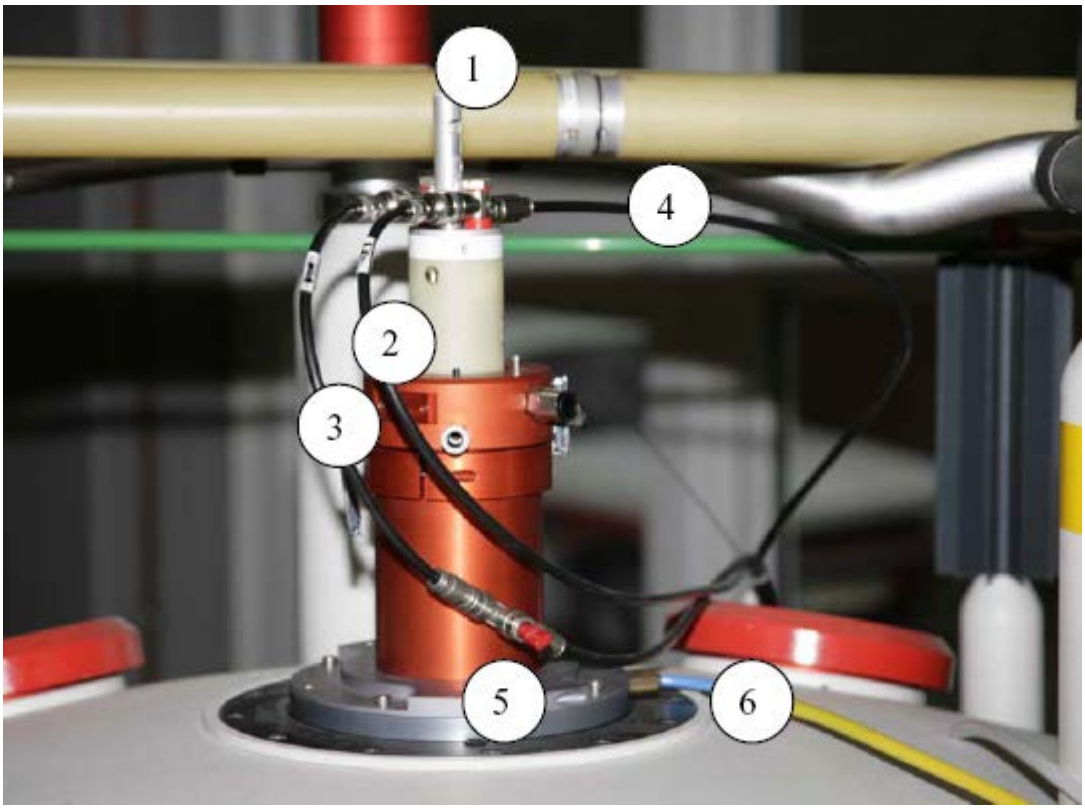


Figure 3.17: WB Probes: Eject/Insert Connections

1	Sample insert/eject	4	Flush gas for transfer tube
2	Eject gas in	5	T-piece to insert tube (allows the flush gas to be fed in at low temperature to avoid ice formation on spinner cap)
3	Insert gas in	6	Shim stack flush connection

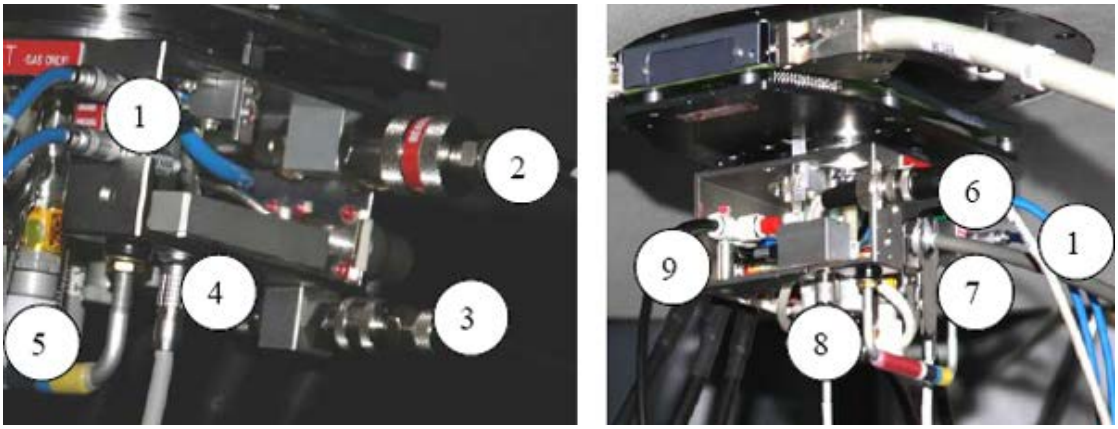


Figure 3.18: WB Probes: DVT, Probe Connections for RT and HT Measurements

1	Thermocouple(s)	6	Heater cable in
2	Bearing gas in	7	VT gas in
3	Drive gas in	8	Spin rate cable
4	PICS cable	9	Flush gas in
5	Heater		

3.5.1.2 Standard Bore (SB) Magnet Probes

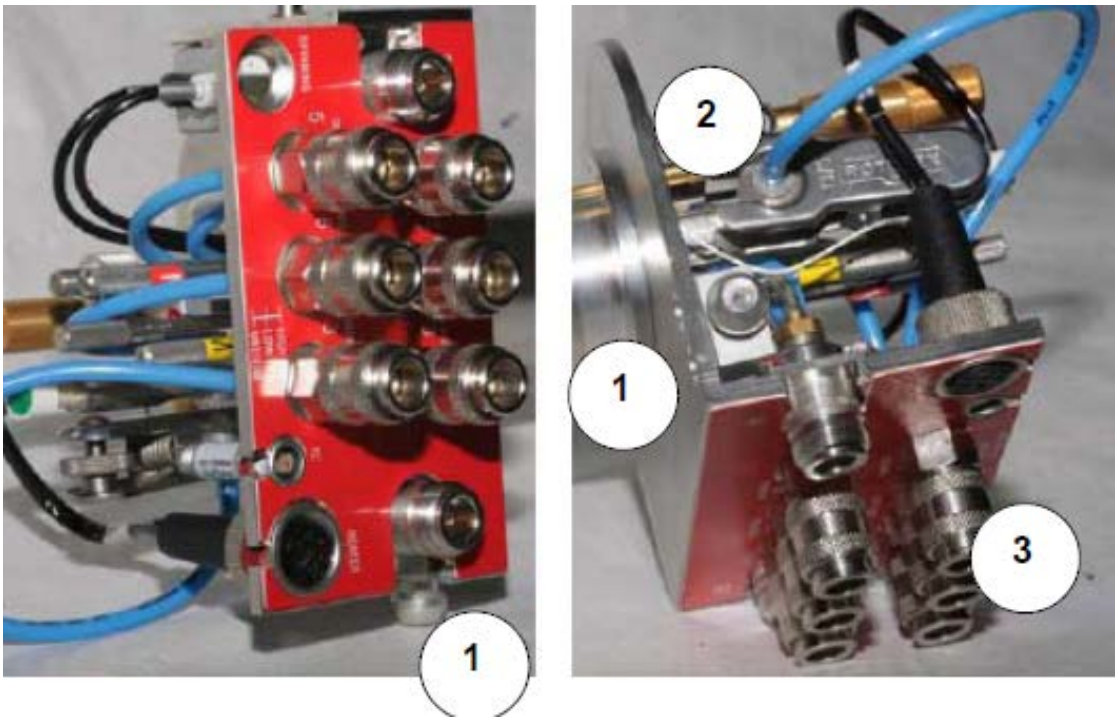


Figure 3.19: SB VTN Probe MAS Connections

1	Frame flush for VT.	3	Bearing connector for ambient temperature gas.
2	Ball joint takes bearing gas from the Quickfit connector at the front into the heater dewar.		

With the standard bore VTN probe, quick fit connectors include:

- Bearing (3).
- Eject (2).
- Drive (5).
- Vertical (7) to the tilt stator for eject.
- Magic Angle (8) to tilt stator into the magic angle.
- Bearing sense (to supervise bearing pressure, shut down in case of a pressure loss).

For LT experiments, the ball joint at the heater dewar must be opened and the transfer line of the heat exchanger or the cooling unit must be connected.

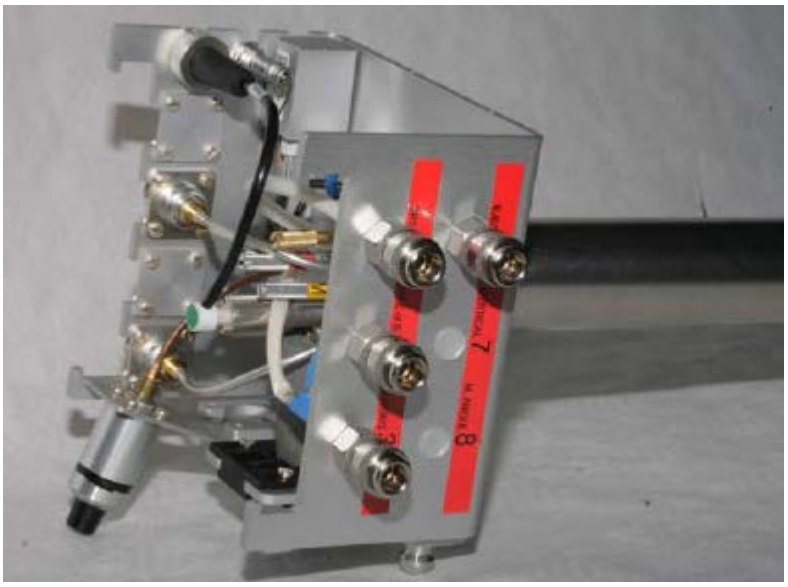


Figure 3.20: SB DVT probe MAS connections.

The numbered connectors are the same as the VTN probe. Connectors 7 and 8 are not present since the probe does not tilt the stator for eject (not required for 2.5 mm probes).

3.6 Additional Connections for VT Operation

If Variable Temperature experiments must be run, there are a few additional connections that need to be made which are not required for room temperature experiments.

First of all, there must be a flow of VT control gas. For MAS probes, this flow can be the bearing gas (VTN) or it can be separate (DVT). In any case, the VT control gas will flow through a dewar which contains a heater. There may be at least one thermocouple which senses the temperature as close as possible to the sample. Several requirements must be fulfilled to obtain precise and stable temperature readings that are as close as possible to the real sample temperature. This topic will however be part of a different chapter.

Some connections are required to control the temperature of the probe/sample, others are necessary to protect the probe and the magnet from extreme temperatures. MAS probes are usually not as well insulated as, for instance, a wideline or PE stationary probe is. Therefore the probe outer shell warms/cools down during the experiment. The heat transfer between heater and probe electronics/probe environment must be kept at a safe level.

Safety precautions involve flushing the probe frame, this serves to keep the tuning elements at acceptable temperatures. Furthermore, the magnet must be kept at legal temperatures to prevent freezing of O-rings or excessive expansion of the inner bore tube. The shim stack must be kept at temperatures below 70 °C, or else the shim coils may be damaged. With MAS probes you must ensure that no wet air is sucked into the eject tube, which would lead to formation of ice on low temperature runs. This requires maintaining some overpressure above the spinner (usually by applying some flow to the insert gas line).

Heat exchangers must be dried with a flow of dry gas before use, so there is no water left in which will ice up the exchanger loop. After use, they must be warmed up and dried with a dry gas flow so that there is no water present, which can lead to corrosion.

The following figures show the various connections to different probes.

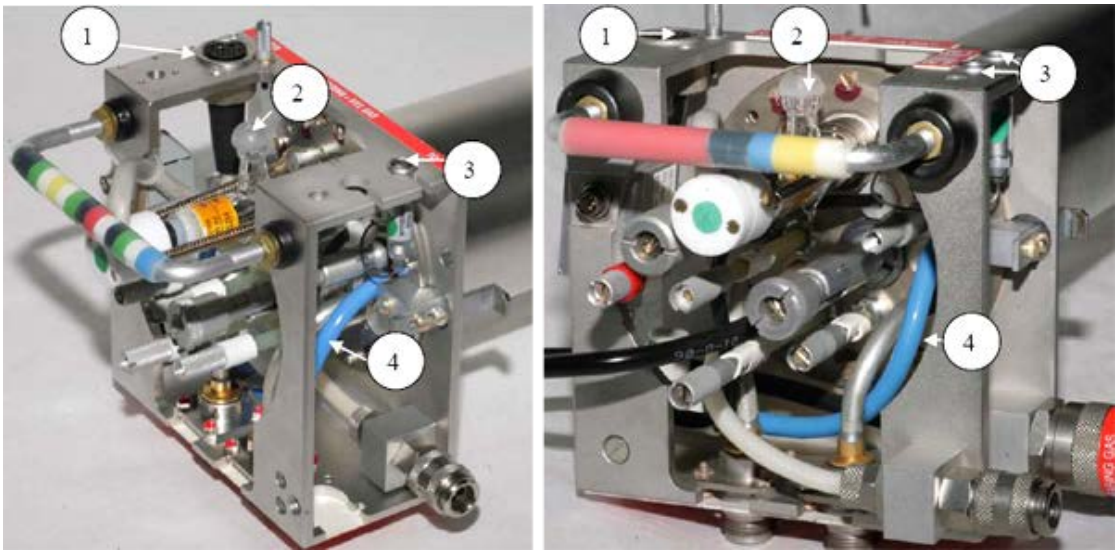


Figure 3.21: WB Probe MAS VTN and WVT, and DVT Probe Connections

1	Probe heater connector	3	TC connector (s)
2	VT gas in	4	Frame flush gas

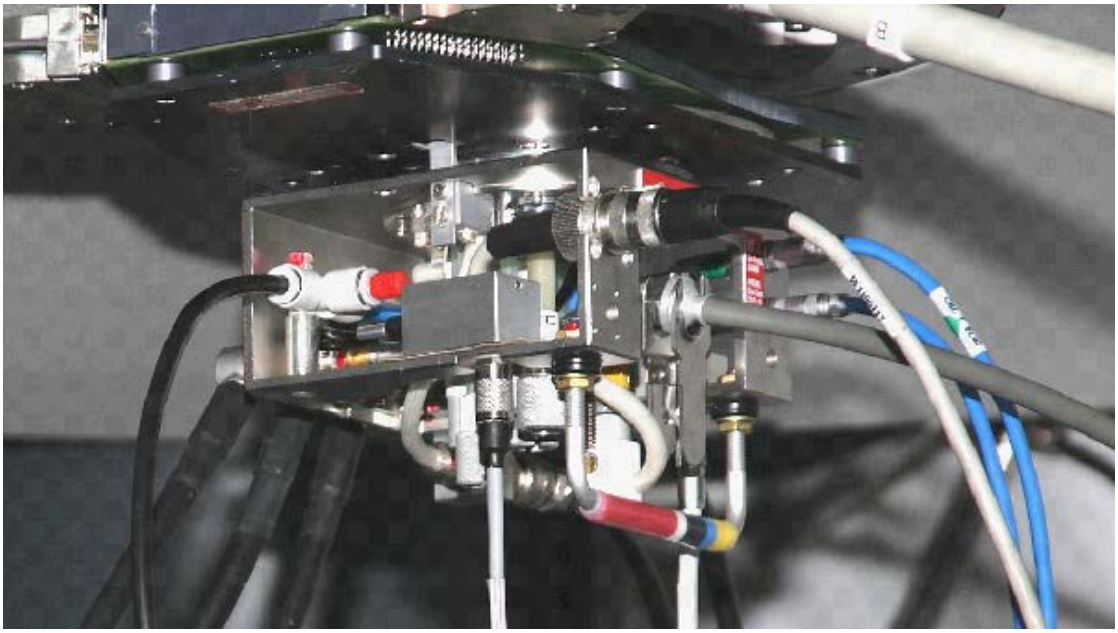


Figure 3.22: WB Probe MAS DVT Connections

1	Probe heater connector	3	TC connector (s)
2	VT gas in	4	Frame flush gas

In the figure above the upper thermocouple connector (read), located at stator out, lower thermocoupler connector (regul), located at stator inlet are connected. In order to read more than one temperature, the VT unit must have the auxiliary sensor module. Only the TC labelled "regul" is used for regulation.

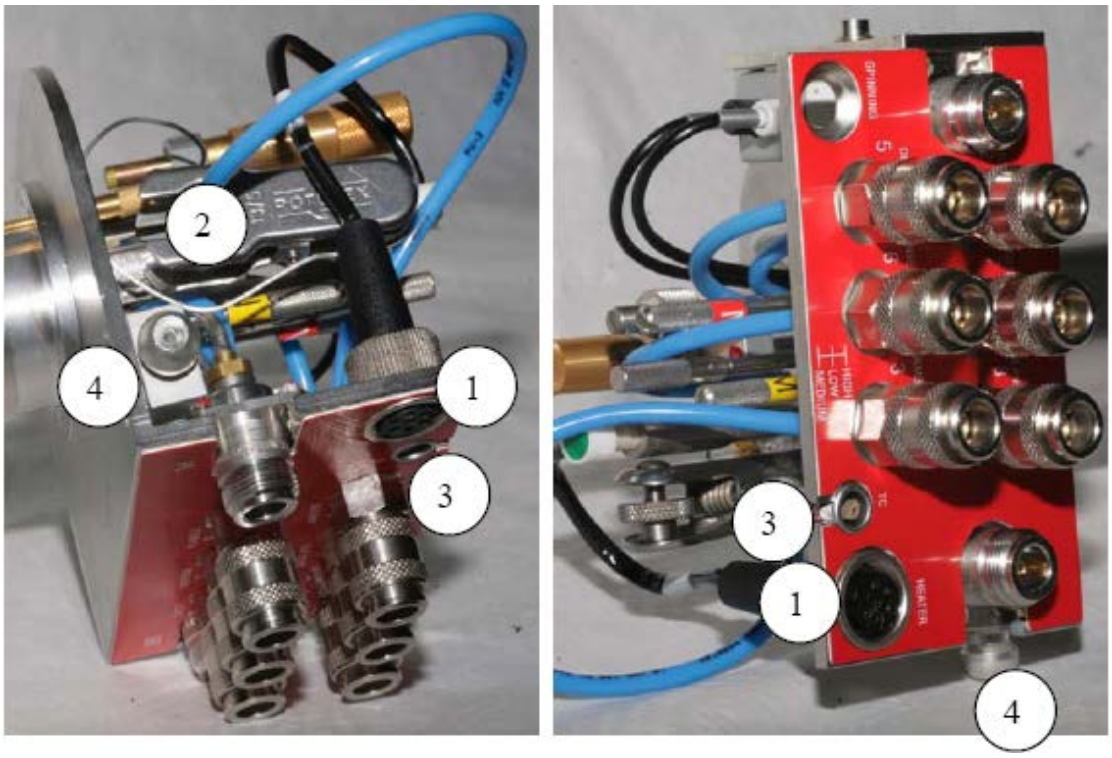


Figure 3.23: SB Probe MAS VTN

1	Probe heater connector	3	TC connector (s)
2	VT gas in	4	Frame flush gas

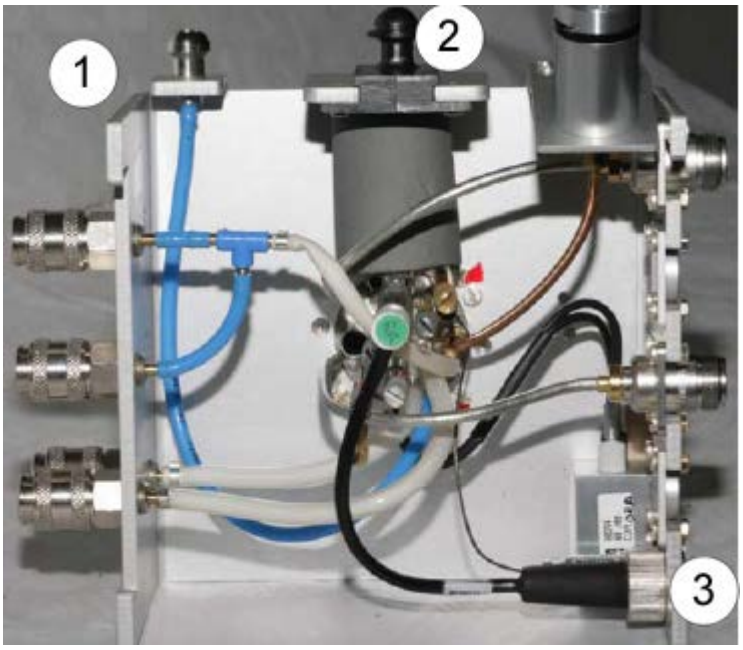


Figure 3.24: SB Probe MAS DVT Connections

1	Frame flush	3	Heater and TC connections
2	VT gas in		

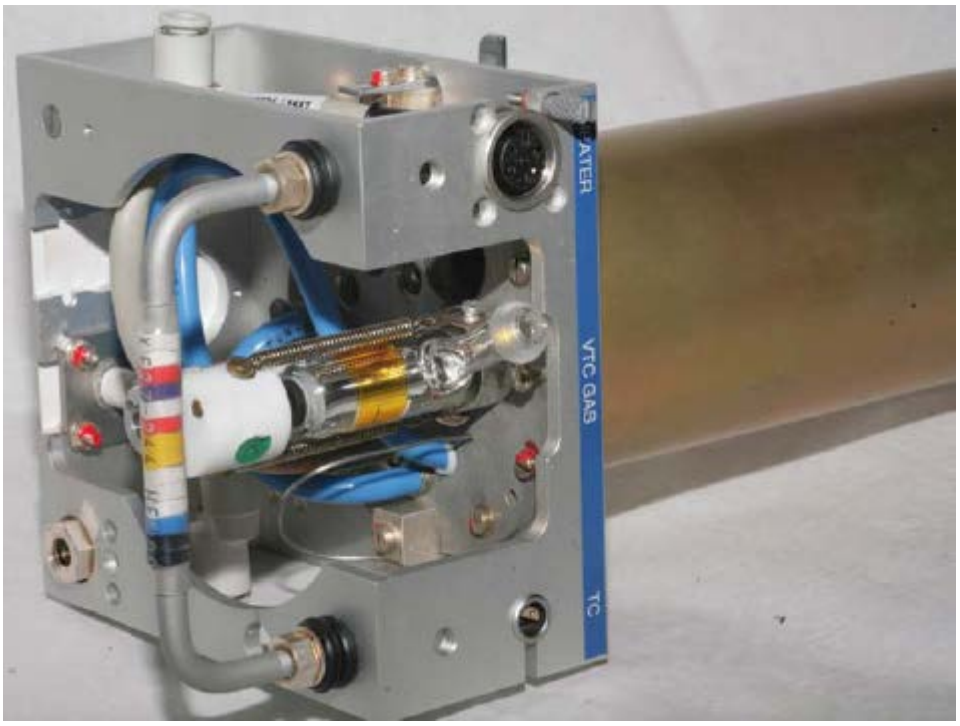


Figure 3.25: WB Wideline or PE Probes

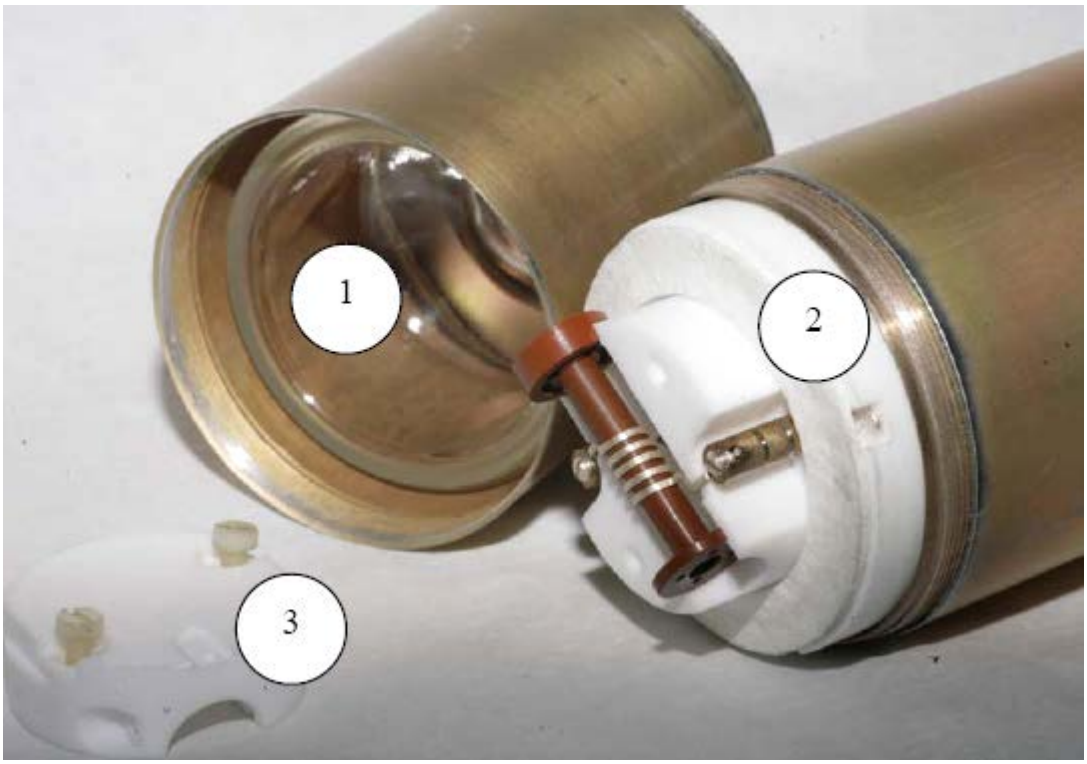


Figure 3.26: WB Wideline or PE Probe Connections

1	Bell shaped glass dewar around sample chamber	3	Cover for coil/sample compartment, fixed with plastic or metal screws
2	Insulating and sealing Al_2O_3 - felt ring		

Figure 3.27: Low Temperature Heat Exchanger for VTN Probes (old style)

In the figure above is a low temperature heat exchanger for VTN probes (old style):

- 1 turn exchanger loop for SB probes;
- 2 turn loop for WB probes;
- 4 turn loop for DVT probes.

Figure 3.28: Low Temperature Heat Exchanger for DVT Probes

The low temperature heat exchanger for DVT probes in the figure above uses the exchanger coil with 6 turns, the larger one may be used for high resolution probes. To use the spring loaded connection device shown in the close-up:

1. Compress the spring.
2. Move the hollow part over the ball joint.
3. Release the spring.

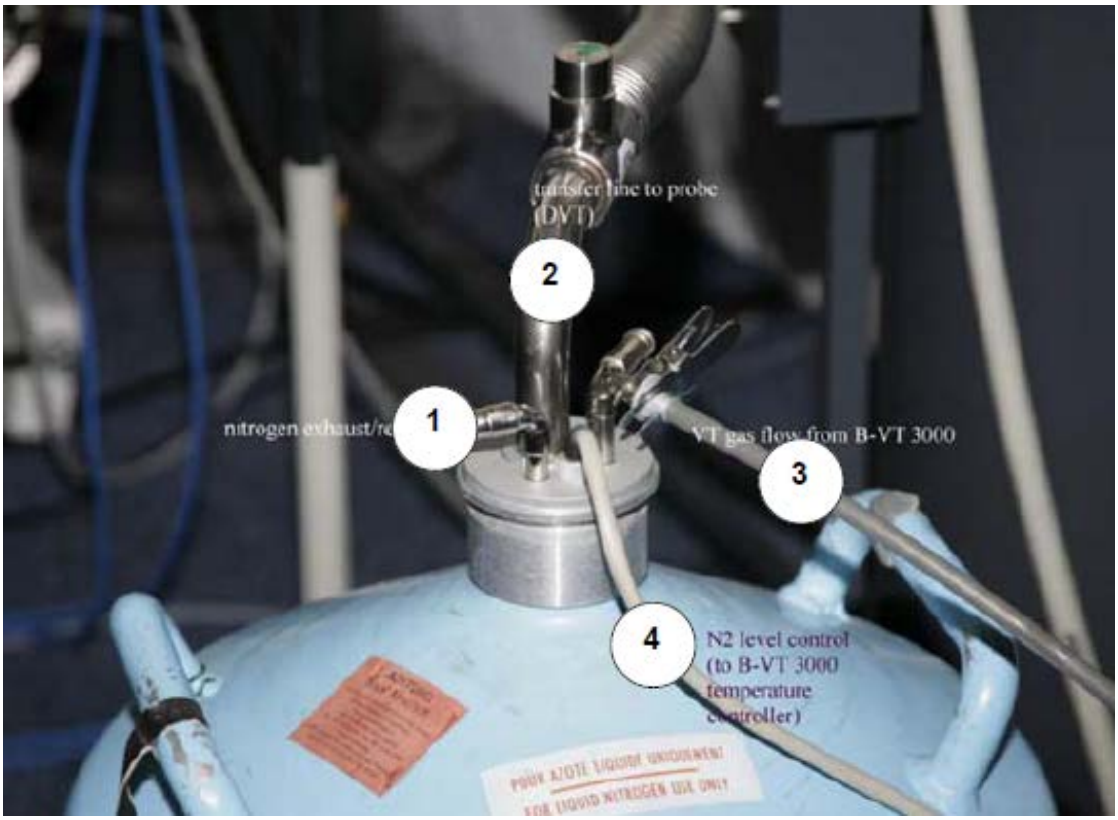


Figure 3.29: Low Temperature Liquid N2 Dewar with DVT Probe/Heat Exchanger

1	Nitrogen exhaust/refill	3	VT Gas flow from B-VT 3000
2	Transfer line to probe (DVT)	4	N2 level control (to B-VT 3000 temperature controller)

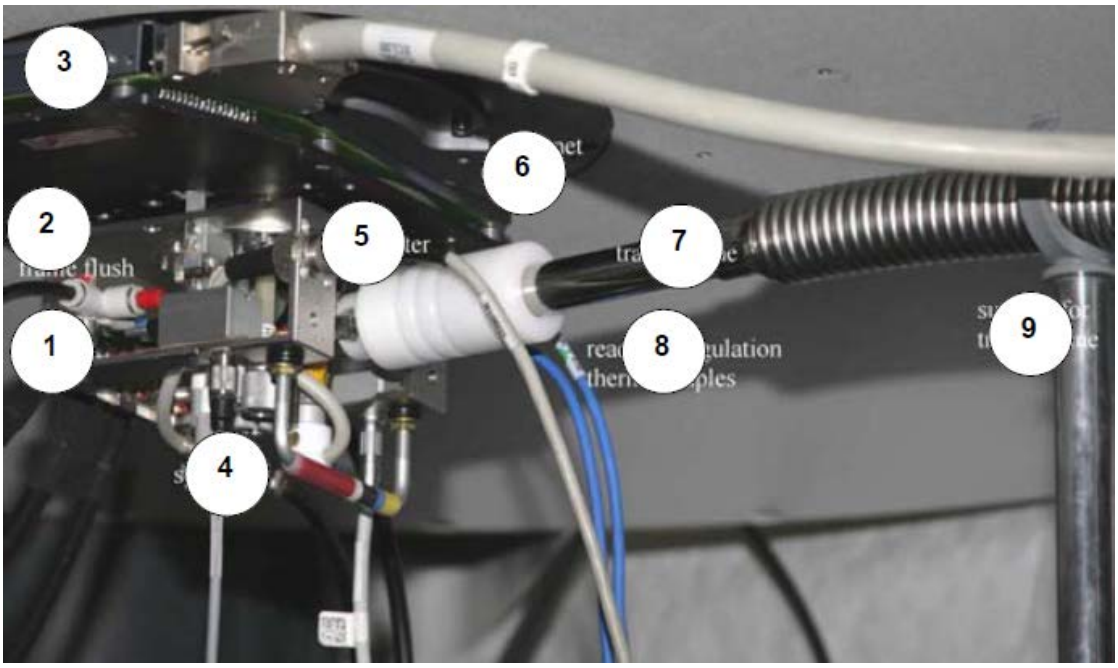


Figure 3.30: Bottom view of Low Temperature DVT Probe/Heat Exchanger

1	Gas in	6	Magnet bore
2	Frame flush	7	Transfer line
3	Shims	8	Read and regulation thermocouples
4	Spin rate	9	Support for transfer line
5	Heater		

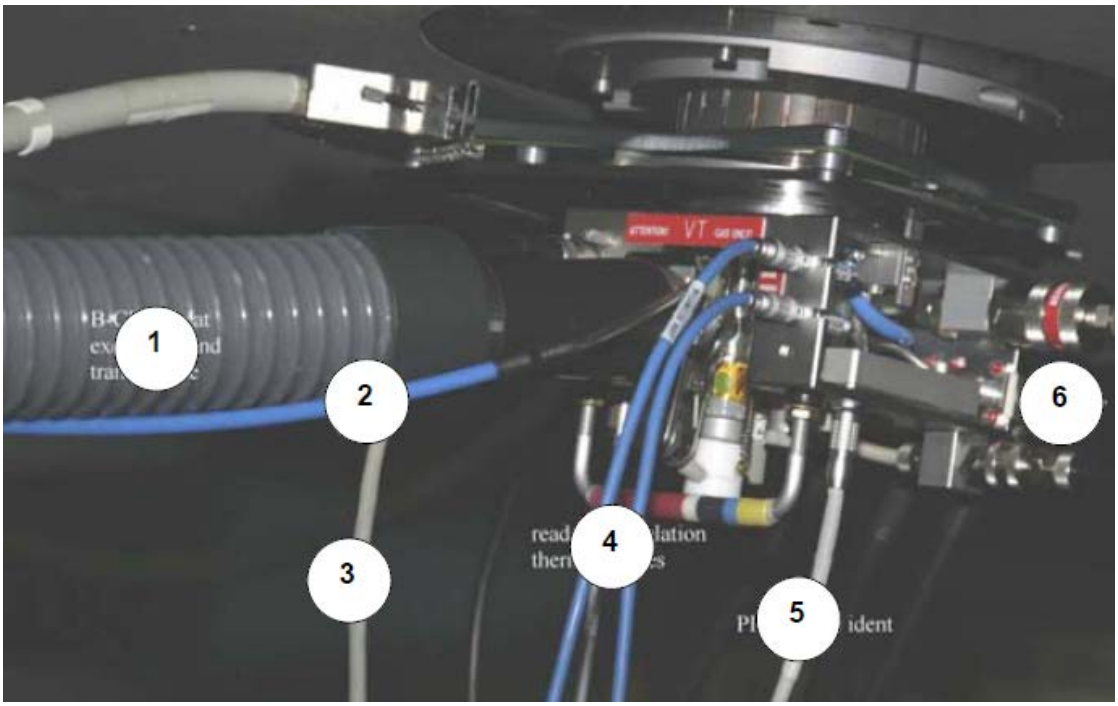


Figure 3.31: Low Temperature Setup with B-CU X (or B-CU 05)

1	B-CU X heat exchanger and transfer line	4	Read and regulation thermocouples
2	Bypass	5	PICS probe identification
3	Heater	6	Bearing, drive

In the figure above is the low temperature setup with a B-CU X (or B-CU 05) for DVT probes only, shown from the probe/magnet side.

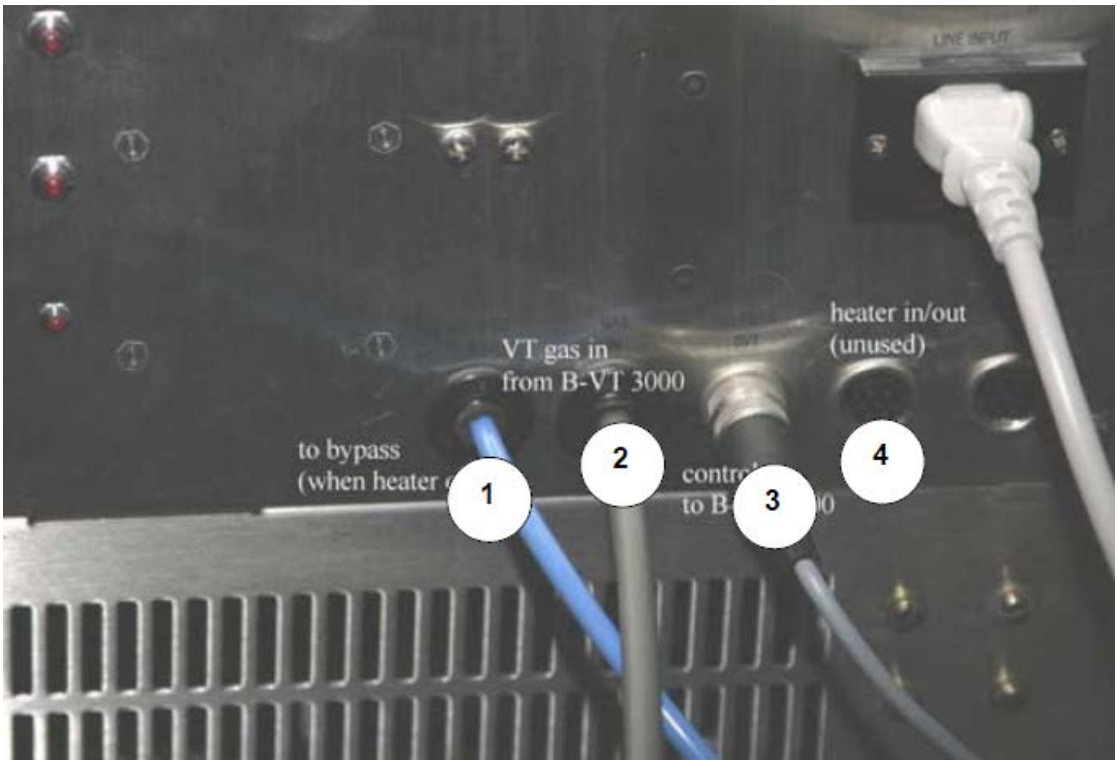


Figure 3.32: Low temperature setup with B-CU X

1	To bypass (when heater is off)	3	Control, to B-VT 3000
2	VT gas in, from B-VT 3000	4	Heater in/out (unused)

3.7 Probe Setup, Operations, Probe Modifiers

3.7.1 Setting the Frequency Range of a Wideline (single frequency) Probe

In a single frequency design, there are more degrees of freedom in tuning the circuit. The frequency range is set by a suitable NMR coil. Fine tuning is done by a variable capacitance (1) and a fixed capacitance inside an exchangeable tuning insert (3). The purpose of this setup is to adapt the probe to the desired task as much as possible. A wideline probe has to cope with a large range of frequencies and line widths, and must provide the shortest possible pulses and highest possible sensitivity at the shortest possible dead time. These requirements cannot be met with one standard setup. Principles of setting up such a probe are:

1. Select an NMR coil (2) with highest inductance that can still be tuned to the required frequency. Choose the smallest coil diameter permitted by your sample, reduce the sample diameter if appropriate
2. Select the symmetrization insert such that the desired frequency is close to the upper end of the available tuning range
3. Select the Q value of the insert according to the expected line widths (higher Q for line widths up to 100 kHz). Please note that multiturn coils, especially multifilament coils, have an intrinsically low Q.

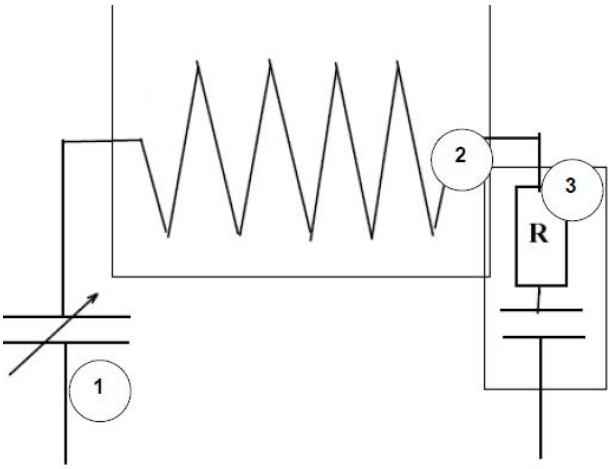


Figure 3.33: RF Setup of a Wideline Single Frequency Probe

1	Tuning capacitance	3	Exchangeable symmetrization and Q-reduction
2	Exchangeable NMR coil		

3.7.2 Shifting the Probe Tuning Range

Most probes cover a fairly wide frequency range. Changing the frequency range of a probe requires either a change in the inductance or capacitance of the circuit. The inductance of a circuit is hard to change unless a coil is mechanically lengthened or shortened. Most probes are tuned over a certain range by variation of a capacitance. The frequency range is then determined by the minimum and maximum capacitance that can be set. In order to make the inductance as high as possible (since the signal from the oscillating magnetization is detected in the inductive part), one usually selects a capacitance with very small minimum capacitance, which again means a too low maximum capacitance. So in most probes additional tuning components must be inserted (removed) to achieve the full tuning range. The highest signal to noise is always reached with maximum inductance and minimum capacitance, i.e. at the high end of the tuning frequency achieved with maximum inductance.

In a wideline probe the NMR coil is easily replaced. So with a few coils of different inductance one can extend the tuning range determined by the tuning capacitor inside the probe. This is difficult in a MAS probe for two reasons:

1. The coil must be carefully aligned such that it does not touch the spinning rotor.
2. There are two frequency ranges to be set- the X-tuning range and the proton tuning frequency.

In such a probe, changing the coil would throw off the proton tuning totally, so a coil change is not possible.

Extending the tuning range of a CP/MAS probe can be done in the following ways:

1. Switch the proton transmission line between λ/4 (low range) and λ/2 mode (high range). The proton transmission line is also part of the X-circuit and is higher (lower) in capacitance in λ/4 (λ/2) mode (only 400 MHz and up).
2. Add a parallel capacitance to the X-tuning capacitance, which makes the capacitance bigger (tunes to lower frequency). This is normally done to shift the tuning range to or below 15N.
3. Add a capacitance in series to another capacitance. This reduces the total capacitance and shifts to higher frequency. A capacitance in series to the λ/4 line will reduce its total capacitance and shift the X tuning to higher frequency.
4. Add a parallel coil to the NMR coil. This reduces the total inductance and shifts to higher frequency, however at the cost of filling factor. The bigger the parallel inductance, the smaller the high frequency shift and the loss.

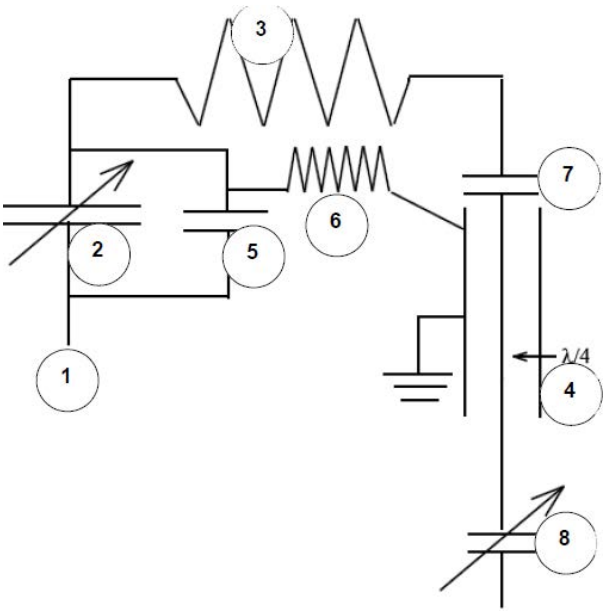


Figure 3.34: Possible Modifiers for Probe Tuning Ranges (400 MHz and up only)

1	X-frequency in	5	Parallel capacitance
2	X-tuning	6	Parallel inductance
3	NMR Coil	7	Serial capacitance
4	lambda/4 switch	8	H-tuning

The figure above illustrates modifiers for probe tuning ranges for 400 MHz and up only. In 300 MHz and lower probes only a λ/4 line can be used, because a λ/2 would be too long, but here λ/4 can be tuned over the full range.

All these modifications may be available for WB probes. In SB probes, they are usually built in (if necessary) and operated by a switch.

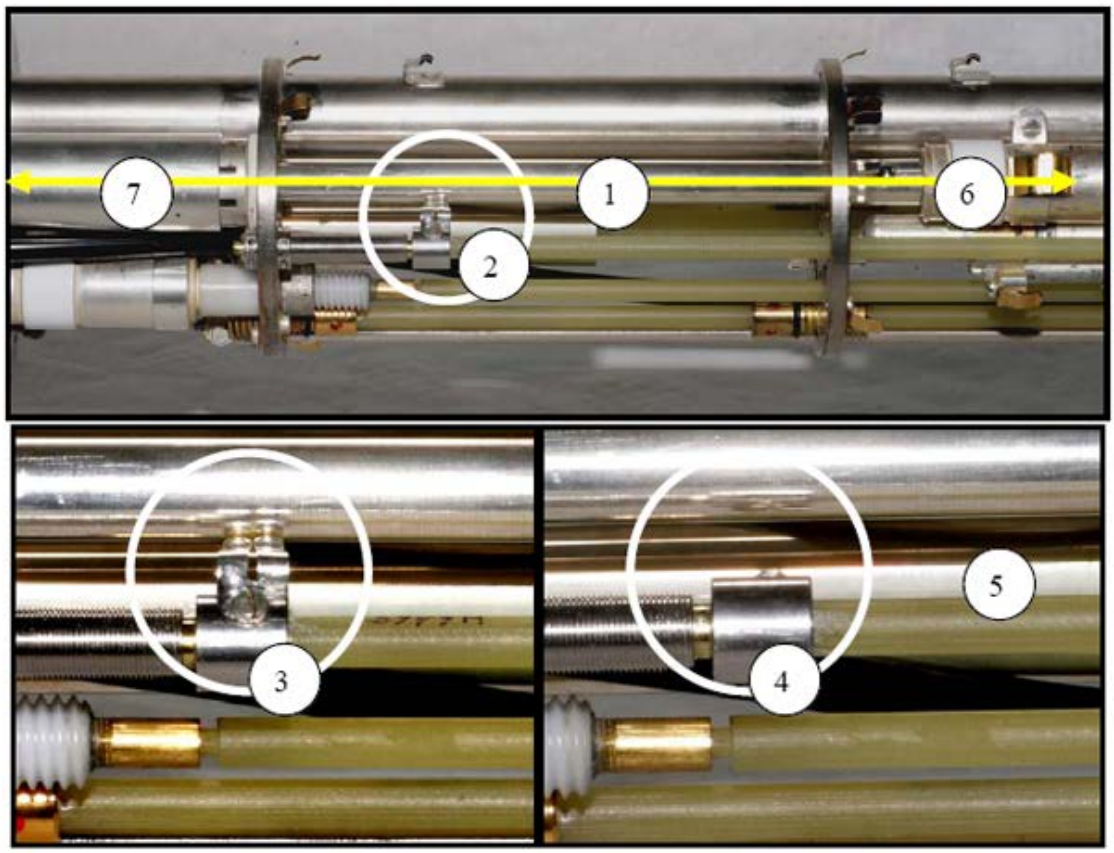


Figure 3.35: λ/4 (low range) and λ/2 Mode (high range), 400 MHz Probe

1	λ-line inner conductor.	4	Switch open rotating clockwise.
2	Rotating switch at λ/4 position.	5	Switch operating rod.
3	Switch closed rotating counterclockwise (seen from probe lower end). Contact springs (grounded) touch the λ-line at the λ/4 position.	6	Tuning capacitor at the end of the λ-line inner conductor: Fine tunes the effective length and there- fore resonating frequency of the λ-line. This tunes the proton channel frequency (“tune”).

At 400 MHz, the wavelength is large, so the λ/4 point is below the closed section (7) of the λ-line. At higher frequencies, the λ/4 point may fall within the closed section 7.

The proton channel (decoupling channel) is usually tuned via a so-called λ-line (transmission line). This is just a coaxial cable or a coaxial conductor with an arrangement of an outer conductor (a tube) and an inner conductor (a rod). The relative diameters and distances and also the dielectric in between (usually air in WB probes) determine the impedance of the transmission line. Since such a line is also an inductance as a capacitance, it is a resonating circuit. If the length of the transmission line equals λ/4 or λ/2 of the RF-wave, it is a λ/4 or λ/2 line. Since the upper end of the transmission line (inner conductor) is connected to the coil, high voltage is required there. This means that the λ/4 point has low voltage but high current, whereas the λ/2 point is at high voltage and low current. A short between inner and outer conductor at the λ/4 position enforces a low voltage/high current and fixes a certain resonance frequency. Some probes are tuned for the proton resonance frequency, with a tunable capacitor at the end of the λ/2-line (400 MHz and up), which changes the effective length of the λ/2-line. Some probes (400 MHz and below) are tuned by shifting the position of the λ/4 short to a higher (higher frequency, shorter length) or lower (lower frequency, longer length) position.

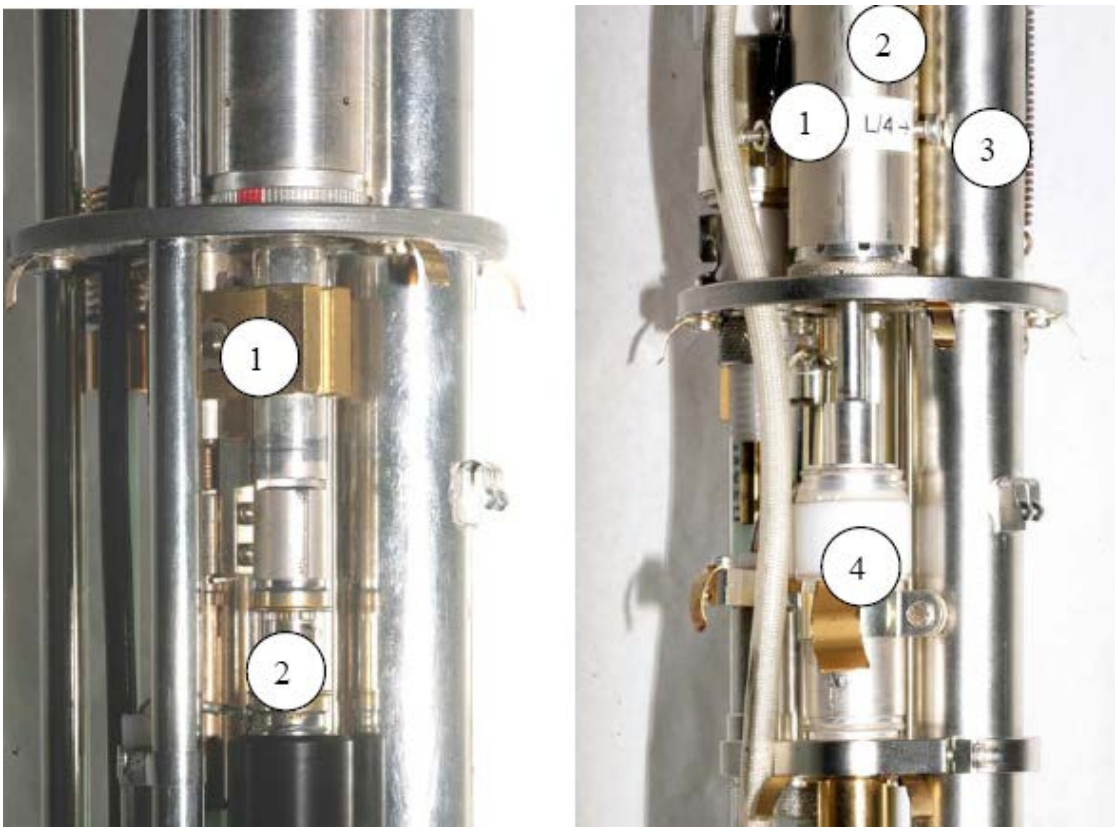


Figure 3.36: A λ/4 only probe (left) and a λ/4 - λ/2 probe (right)

1	Brass block to ground	3	Screw is used to set λ/4
2	Transmission line outer conductor	4	Proton tuning capacitor

On the left side of the figure above is a λ/4 only probe (200, 300 MHz, 400 low range only probes). The transmission line is only λ/4, proton tuning is done by moving the brass block to ground (1). Proton matching is done with capacitor 2.

On the right side of the figure above is a 600 MHz λ/4 - λ/2 probe. Due to the high proton frequency the λ/4-length shortens, the λ/4 point (1) moves inside the transmission line outer conductor (2). The screw (3) is used to set λ/4 (screw in) and λ/2 –mode (screw out). (4) Proton tuning capacitor.

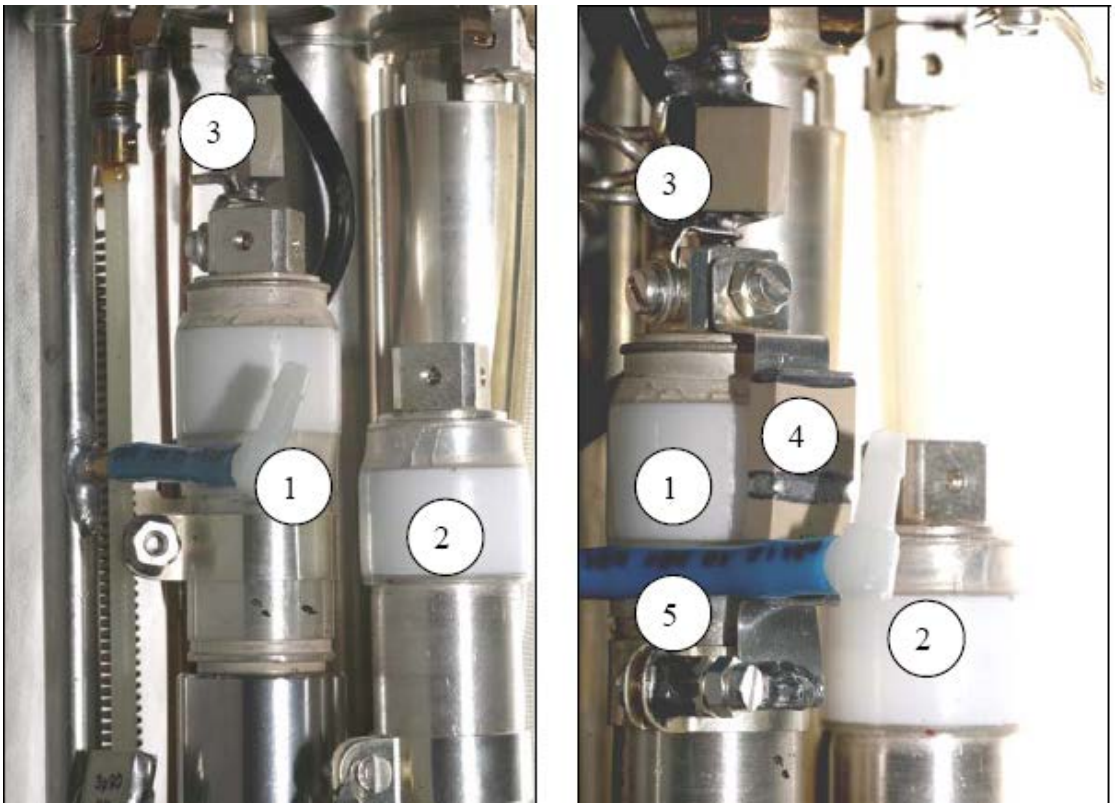


Figure 3.37: Without/with Parallel Capacitance to Shift the Tuning Range to Lower Frequency

1	X-tuning variable capacitance	4	Low range extension with parallel capacitance
2	Y-tuning capacitance	5	Frame flush air outlet (to protect sensitive capacitors)
3	Proton reject filter		

Adding a parallel capacitance does not decrease the efficiency of a circuit. However, a certain circuit has the highest possible efficiency if it is tuned with maximum inductance and minimum capacitance. Maximum inductance can usually not be achieved due to spacial restrictions in the stator (MAS probes) or due to losses in Q if the coil becomes too large (increase in resistance). Furthermore, capacitances are tunable, inductances usually are not, so a wide tuning range can only be achieved via exchangeable inductances and/or capacitance with a wide tuning range. Capacitances may look different than the one shown in the picture. Frequently, two larger capacitances are used in series to form a smaller capacitance withstanding higher voltage. A parallel capacitance lowers and narrows the tuning range.

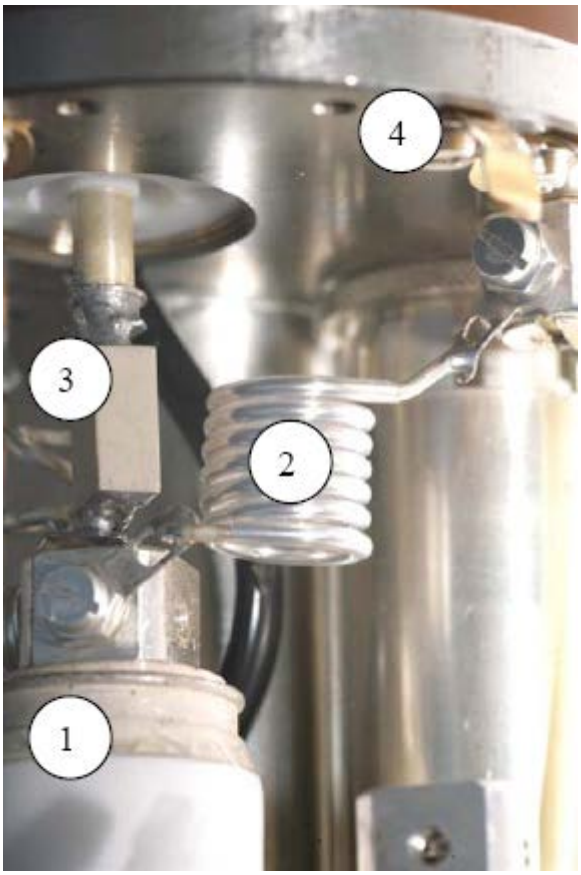


Figure 3.38: Parallel Coil to Shift the Tuning Range to Higher Frequency

1	X-Tuning cap	3	Proton reject filter („trap“)
2	Parallel coil	4	Probe ground

This additional coil is electrically connected in parallel to the detection coil. Since it is connected behind the proton trap, the proton channel is not influenced, but the X and Y channel will be affected, because only part of the inductance is now filled with sample. A parallel coil therefore reduces the RF efficiency quite substantially. The losses increase as the inductance (size and number of turns) of the parallel coil decreases. A coil of the same inductance as the detection coil will cost 50% in S/N and pulse voltage. Usually, these coils introduce about 30% loss.

3.7.3 Adding a Frequency Channel to a Probe (WB probes only)

Probes are produced as single channel, double (channel), triple or quadruple probes (1, 2, or 3 RF connectors on the probe). It is not possible to modify a probe produced as a double channel probe into a triple probe, but a triple probe may be used as triple or double probe. As multiple tuning will reduce the RF performance of a probe on the other channels (if they are part of the same RF network), it is better to remove an unused RF channel, if this is possible.

The usual case is triple probes or quadruple probes. A triple probe can be tuned to ^{1}H, X and Y (where X is the higher frequency, Y is the lower frequency, both in the X-nuclei range). A triple probe X/F/H only has 2 RF connectors, because ^{1}H and ^{19}F are tuned to the same RF connector simultaneously. Likewise, a quadruple probe is an X/Y/F/H probe, with the X channel tuned to X and Y, and the proton channel tuned to ^{19}F and ^{1}H, so there are only 3 RF connectors. Check the figures in section *RF-Filters in the RF Pathway [▶ 19]* for the connection of such probes.

So the double tuning of the X channel into an X and a Y channel is an optional operation available for WB probes. SB probes are always fixed multi channel probes, with no option to insert/remove an RF channel.

Such probes have 2 complete X-tuning circuits, almost identical in construction. Activating the second channel means: Insert a filter which is tuned to reject at the exact X-frequency (exactly to ^{13}C frequency for a $^{13}C/^{15}N$-^{2}H probe). The X channel frequency is fixed to the specified frequency (in this case, ^{13}C), while the Y frequency has a broader tuning range (in this case ^{15}N-^{2}H). With different filter-inserts, that same probe can be modified to different frequency combinations, with the following restrictions:

1. A frequency outside the tuning range of the probe in double mode cannot be reached in triple mode (for instance, if the probe in double mode does not tune to ^{31}P, a triple insert for ^{31}P-^{13}C cannot be provided (without a frequency range shift shown in the first figure in this section). As the triple tuning insert will act as a load to the whole circuit, the frequency range will shift a bit to lower frequencies.
2. Both frequencies (X and Y) should be within the same basic probe tuning range high or low (λ/2 (high range) or (λ/4 (low range)) mode.

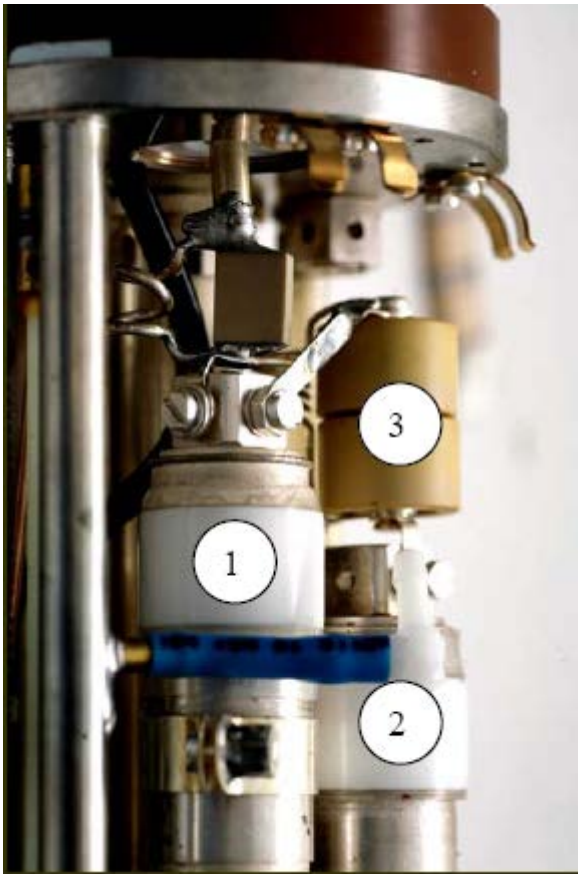


Figure 3.39: Mounting a Triple Insert into a Triple Probe

1	X-Tuning capacitor	3	X-Y trap (stops X-frequency into Y-channel, but not Y into X-channel)
2	Y-Tuning capacitor		

These inserts can be made to optimize the X or the Y channel. They should be mounted in exactly the position as indicated on the information sheet included with the probe.

Reversing the trap may mess up tuning! Observe probe manual instructions! For low range (below ^{13}C) nuclei, the probe must be in λ/4 mode. Combinations of low/high range nuclei are difficult/impossible and always lossy!

3.8 Mounting the Probe in the Magnet/Shim Stack

Usually, the service engineer installing the magnet and the rack has considered the local restrictions, and placed the system such that all operations which may be required for the proper probe installation are conveniently possible. This refers to the ease of access of the magnet bore from below and above, to mount the probe and the sample insert devices, and also to the possibility of placing VT equipment (liquid N2 dewar, B-CU 05 or B-CU X) in a convenient location that allows access without restricting standard operations.

Depending on the type of probe, VT control gas enters the probe from the side or from behind. So it must be possible to attach the heat exchanger transfer line from the appropriate side. Furthermore, the weight of the transfer line must be relieved from the probe dewar ball joint, so there should be appropriate fixation points for the transfer line. It is important that the transfer line enters the ball joint as straight as possible, as a ball joint will cut off the flow when strongly tilted to an angle.

It should also be possible to reach the probe tuning elements (since they need to be operated frequently) as easily as possible from the operators chair. Also, when tuning the probe, the video screen and the preamplifier display should be easily visible.

Figure 3.40: Example of a 600 WB NMR Instrument Site

In the example of a 600 WB NMR instrument site in the figure above provides easy access to the probe from either side.

3.9 EDASP Display: Software Controlled Routing

The menu edasp shows all relevant RF routing and allows the routing which is under software control to be changed. The following restrictions apply:

- Connections between the transmitter and preamplifier cannot be changed (the command **edasp setpreamp** with NMR Super user permissions does that).
- Channel F1 is the detection channel by default (which is no limitation).
- Detection must usually be routed via the same SGU as the F1 pulses are, since that SGU will supply the phase coherent reference signal. In some cases, pulsing and detection may use different SGU's, but provisions must be made to add the signal up coherently (using exactly the same frequency on both channels is usually coherent).
- Routing between SGU's and transmitters can only be selected if **cf** (the configuration routine) has found a hardware connection that supports this routing.
- Most routine applications will route correctly if the **default** button is pressed in the short menu version (receiver routing not shown).
- If an incorrect or potentially dangerous routing is selected, an error message or a warning will pop up. Error messages will not allow selection of this routing.
- On AVIII instruments (with SGU/2), one SGU can produce two pulse trains (within the same NCO frequency setting range, 5 MHz).

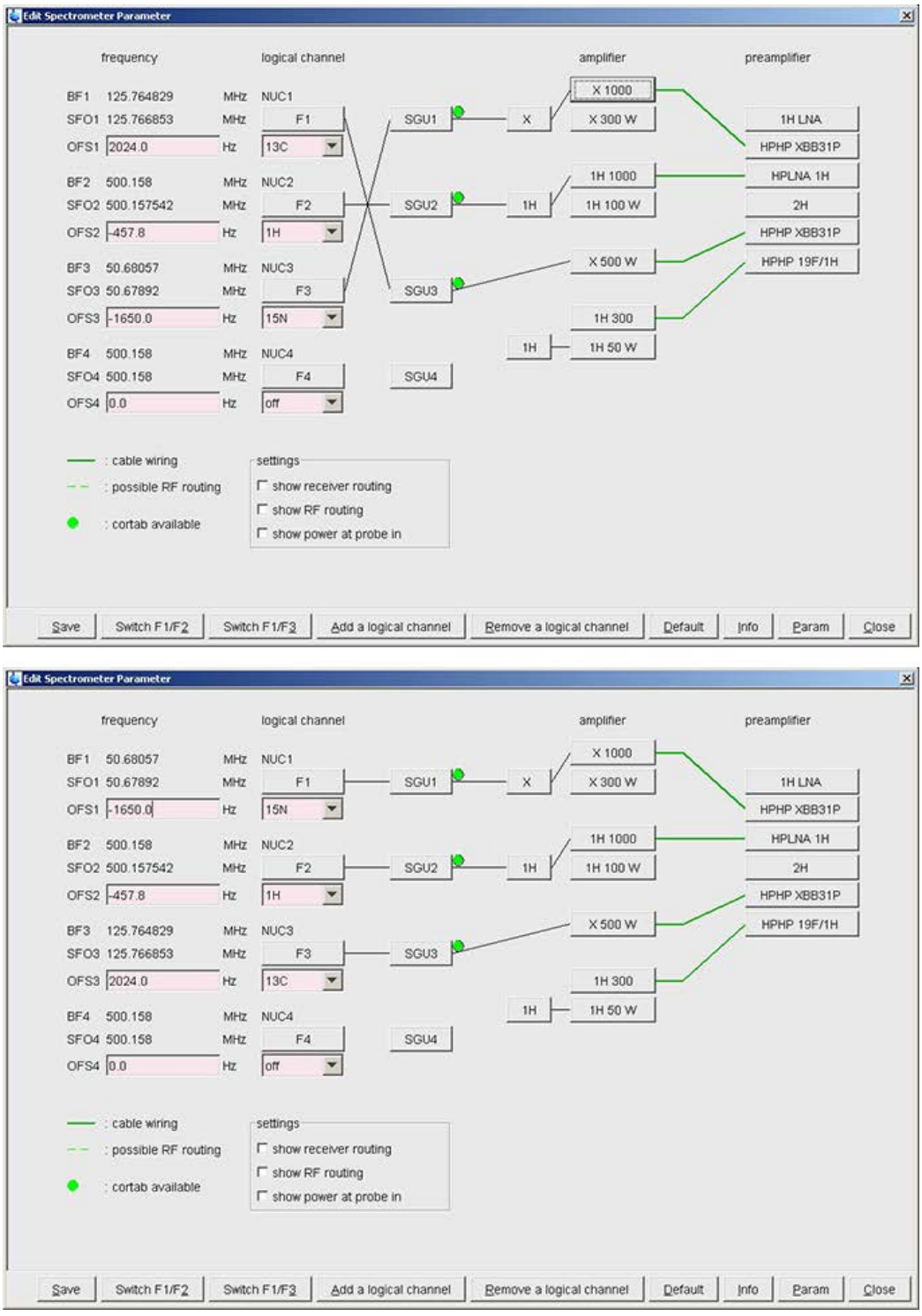


Figure 3.41: Short Display, Pulse Routing Only for C/N/H DCP or REDOR Experiment, observing ^{13}C (above) and ^{15}N (below).

In the figures above is a “short” display, pulse routing only, for a C/N/H DCP or REDOR experiment, observing C and observing N (without any hardware change!). Green dots indicate CORTAB linearization. ^{13}C routed via 500W transmitter since ^{13}C requires less power than ^{15}N.

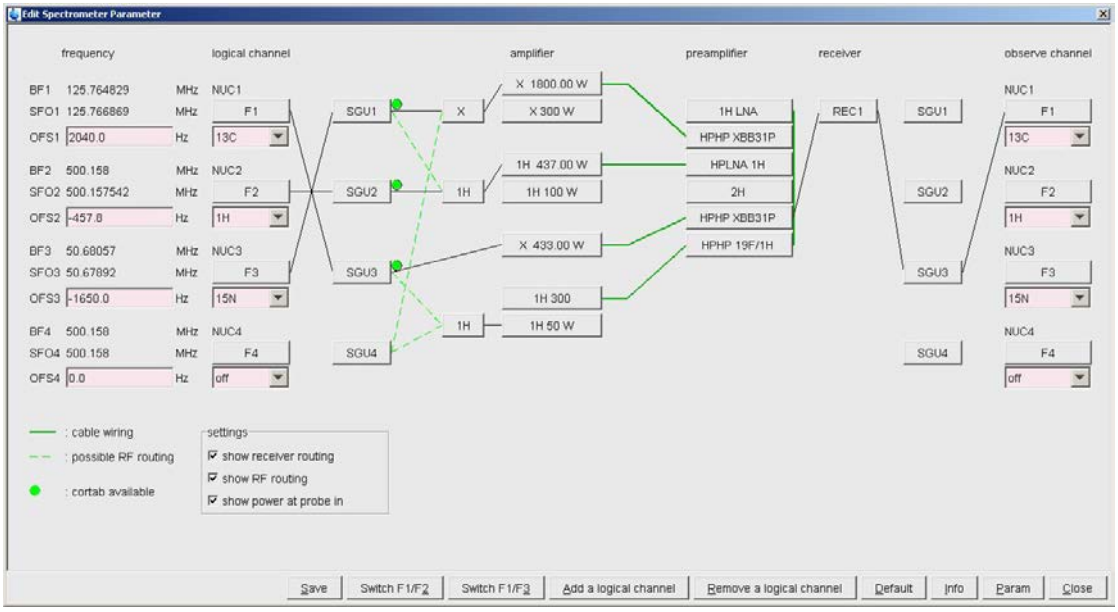


Figure 3.42: Long Display, Pulse and Receiver Routing

The figure above is a “long” display, with pulse and receiver routing:

- Green dots: CORTAB done for this path, transmitter linearized.
- Dotted green lines: Possible (hardwired) routing is shown (show RF routing).
- Receiver routing: SGU3 used for transmit and receive (show receiver routing).
- Power indication: Maximum possible power output (as measured, show power at probe in).

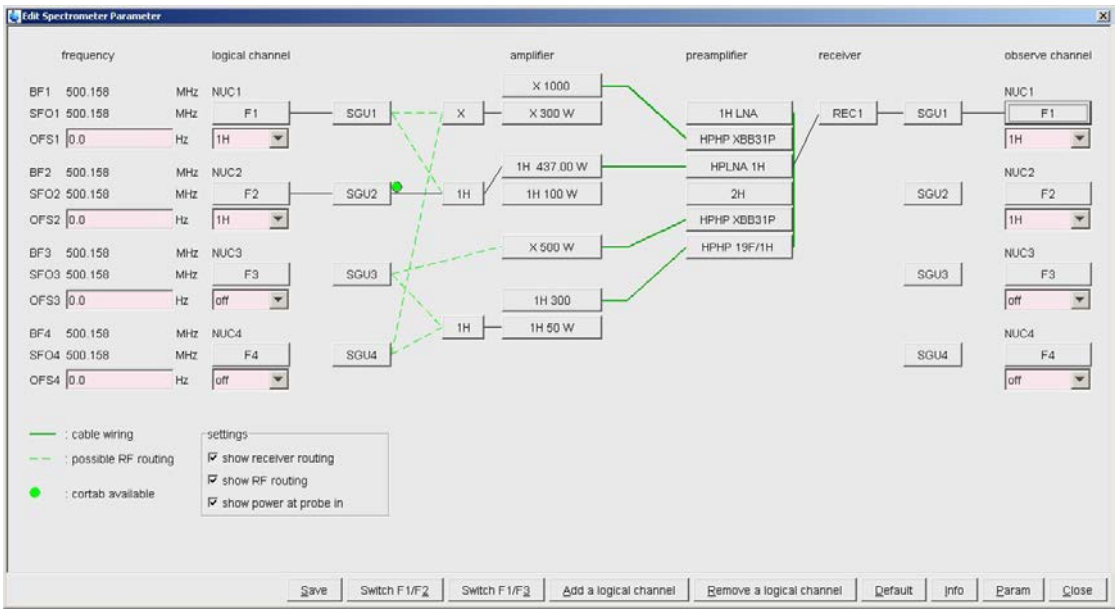


Figure 3.43: Pulse on F2, Observe on F1 - Routing

In the figure above the SGU2 is used for pulsing and the SGU1 is used to receive.

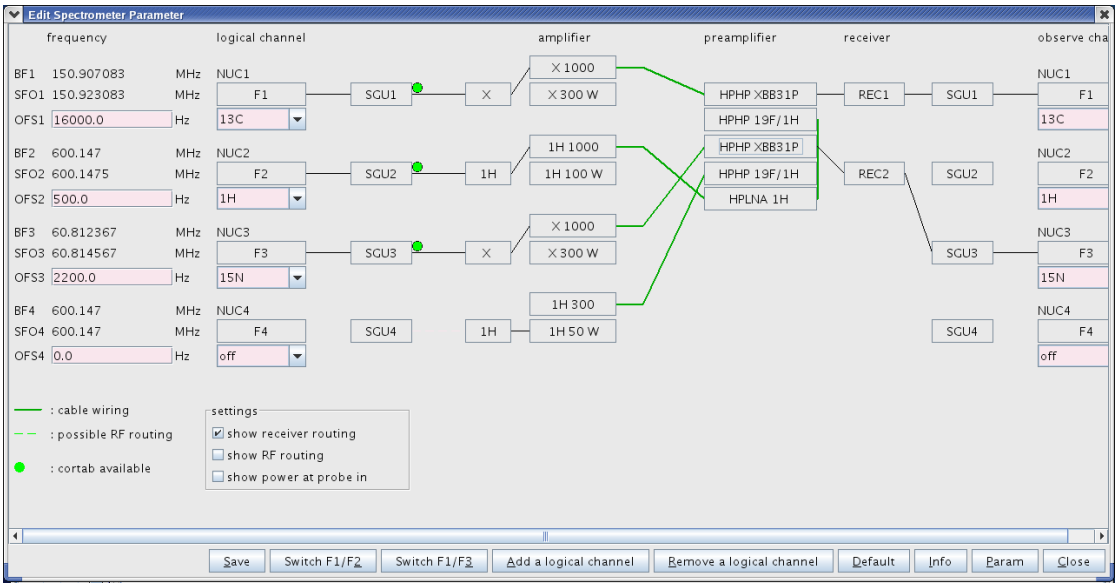


Figure 3.44: The edasp Display for a System with two Receiver Channels

In the figure above the edasp display for a system with two receiver channels, set for observe on ^{13}C and ^{15}N, while decoupling on protons. The same SGU is used for pulse and detection on both receiver channels.

4 Basic Setup Procedures

This chapter contains information and examples on how to set up basic solid-state NMR SSNMR) experiments. We'll begin with the settings for the RF-routing of the spectrometer, some basic setup procedures for MAS probes and how to measure their (radio frequency) RF-efficiency and RF-performance. Accurate measurement of the pulse lengths and the associated RF-power levels is essential for solid-state NMR experiments. In SSNMR, RF-field amplitudes are often expressed as spin nutation frequencies instead of 90° pulse widths. Spin nutation frequency nrf and 90° pulse width are related through the reciprocal of the 360° pulse duration $4t_{p90}$ such that:

$n_{rf} = 1/(4t_{p90})$= RF-field in Hz (with t_{p90} in µsec)

Setting up the magic angle, shimming a CPMAS probe, setting up cross polarization and measuring probe sensitivity for 13C will also be explained. This is part of probe setup and performance assessment during installation. However, regularly scheduled performance measurements should be part of the hardware, probes and spectrometer maintenance. Therefore, these checks should be performed periodically.

The checks also need to be performed if an essential piece of hardware has been exchanged. In the following, we describe all steps which are necessary to assess performance of a CPMAS probe, along with all necessary settings. Detailed information about TopSpin software commands is available in the help section within the appropriate chapter.

Setting up a CPMAS probe from scratch requires the following steps:

1. Mount the probe in the magnet and connect the RF connectors of the probe to the appropriate preamps.
2. Connect the spinning gas connectors and the spin rate monitor cable.
3. Insert a spinner with finely ground KBr and spin at 5 kHz.

It is assumed that these operations are known. If not, please refer to the following sources:

- Probe manual.
- MAS-II pneumatic unit manual in TopSpin/help.
- SBMAS manual in TopSpin/help.

This chapter will include:

4.1 General Remarks

Despite the fact that most spectra taken on a CP/MAS probe look like liquids spectra, the conditions under which they are taken must account for the presence of strong interactions. This basically means that:

- Fast spinning, and
- High power pulses are applied.

Fast spinning requires a high precision mechanical system to allow spinning near the speed of sound. This requires careful operation of the spinning devices. Please read the probe manual carefully!

High power decoupling in solids requires 20-fold RF fields compared to liquids spectroscopy, since we are dealing with >20 kHz dipolar couplings rather than maximum 200 Hz J-couplings! This means that RF voltages near the breakthrough limit must be applied, and that currents of far more than 20 A occur.

It is therefore essential that:

- Power levels for pulses must be carefully considered before they are applied. Always start at very moderate power levels with an unknown probe, find the associated RF field or pulse length and then work your way towards specified values. The same applies for pulse lengths, especially decoupling periods, since the power dissipation inside the probe is proportional to pulse power and duration. Always observe the limits for duty cycle and maximum pulse power. Please refer to the probe specifications for more information. Never set acquisition times longer than required!
- Spinners and turbine must be kept extremely clean. Any dirt, especially oil, sweat from fingers, water will decrease the breakthrough voltage dramatically. Make sure the spinner is always clean (wipe before inserting, touch the drive cap only) and the spinning gas supply is carefully checked to provide oil-free and dry (dew point below 0 °C) spinning gas. Compressors and dryers must be checked and maintained on a regular basis. Any dirt inside the turbine will eventually cause expensive repairs.

The following setup steps only need to be performed during installation or after a probe repair. The test spectrum on glycine should be repeated in regular intervals to assure probe performance.

4.2 Setting the Magic Angle on KBr

For the following steps, generate new data sets with appropriate names using the **edc** command to record all individual setup steps.

4.2.1 RF-Routing

The spectrometer usually has 2 or more RF generation units (SGU's), transmitters and preamplifiers. In order to connect the appropriate SGU to the appropriate transmitter and the transmitter to the associated preamp where the probe channels are connected, there are several routing possibilities. In order to minimize errors in hardware connections, the routing is under software control where possible. When cable connections need to be done manually, the software does not allow a change. These connections are made during instrument installation.

Enter the **edasp** command (or click the **Edit** button in the nucleus section in the acquisition parameter window eda) in order to get the spectrometer router display.

Alternatively, click on the routing icon in **eda**.

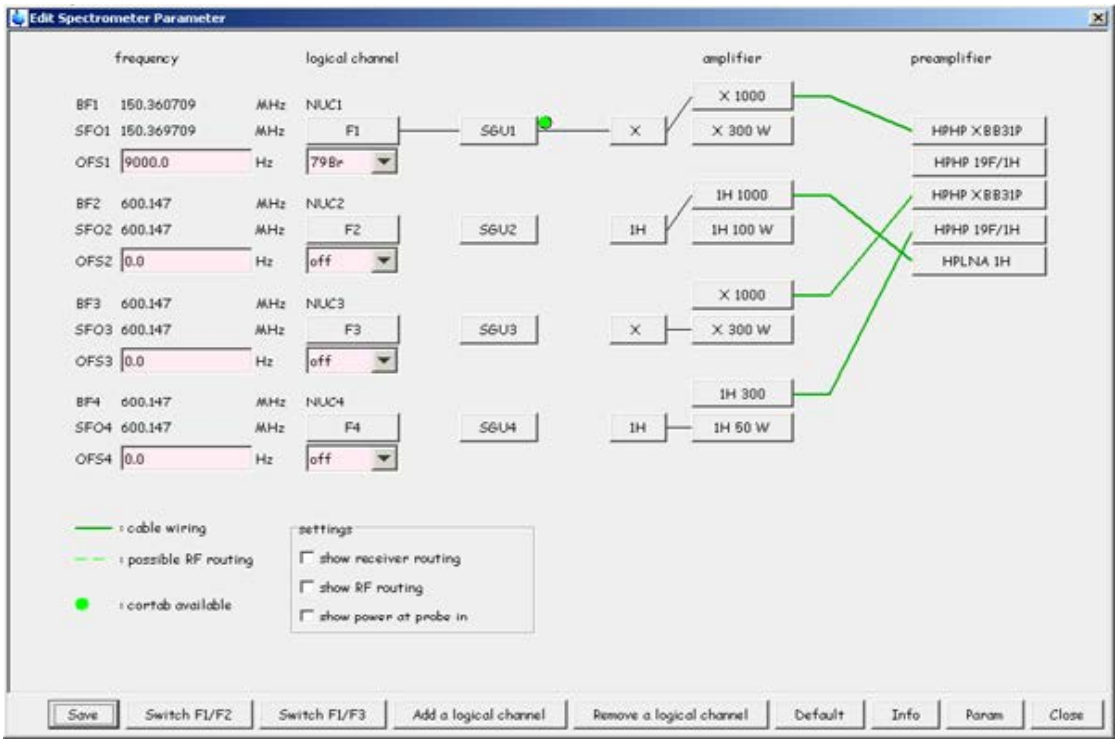


Figure 4.1: Routing for a Simple One Channel Experiment

The figure above shows the routing for a simple one channel NMR experiment using the 1000 W output from the high power amplifier.

In this menu, 4 RF channels are available. These 4 RF channels can be set up for 4 different frequencies. The two left most columns labelled frequency and logical channel define the precise irradiation frequency by setting the nucleus and the offset **O1** from the basic frequency. In this example, we want to set up for pulsing/observe on the nucleus ^{79}Br. Selecting ^{79}Br for channel F1 defines the basic frequency **BF1** of ^{79}Br (in this case on a 600 MHz spectrometer, 150.360709) and the adjustable offset (9000 Hz in this case). Both values are added to show the actual frequency setting, **SFO1**. The frequency setting is taken from a nucleus table which is calculated for the respective magnetic field B_0. The index 1 in **O1**, **BF1**, and **SFO1** refers to the RF channel 1 (which is also found in the pulse program where the pulse is defined as p1:f1. Note that this index does not exist for the following columns which represent the hardware components (SGU1-4 refers to the slot position in the AQS-rack). The lines connecting the (software) channel F1 to the actual frequency generation and amplification hardware can be drawn ad libitum as long as the required hardware connections are present. The connections between transmitter and preamp cannot be routed arbitrarily, because every transmitter output is hardwired to a preamplifier, so the lines are drawn in green. Please note that the nucleus in channel F1 is always the observe nucleus.

To set up for ^{79}Br observation, click on **Default** and the correct routing will be shown.

The green dot between SGU1 and amplifier 1 indicates that for this nucleus in this connection, the transmitter has been calibrated for amplitude and phase linearity (CORTAB).

For example, if you select a nucleus where this has not been calibrated and the green dot is not visible, the same power level setting in dB will produce >6 dB more power (> 4-fold power) which may destroy your probe. Calibrate power levels in such a case starting with 10 dB less power (higher pl(n)-value) to prevent destruction of your probe!

The connections between the SGU(n) and transmitter (n) can be altered by clicking on either side of the connecting line (removes the connection), and clicking again on both units you want to connect (route). In case of high power transmitters, you have two power stages which you can select by clicking on the desired stage. High power stages require the parameter **powmod** to be set to “high”. Selecting a path which is not fully routed will generate an error message.

To leave the display, click on **Save**. Make sure your probe X-channel is connected to the selected preamplifier (the sequence of preamplifiers in **edasp** represents the physical position of the preamplifier in the stack). If the preamp is a high power type, make sure the correct matching box is inserted into the preamp (for 500-800 MHz systems, it would be labelled for the frequency range 120-205 MHz). Connections are shown in the figure below:

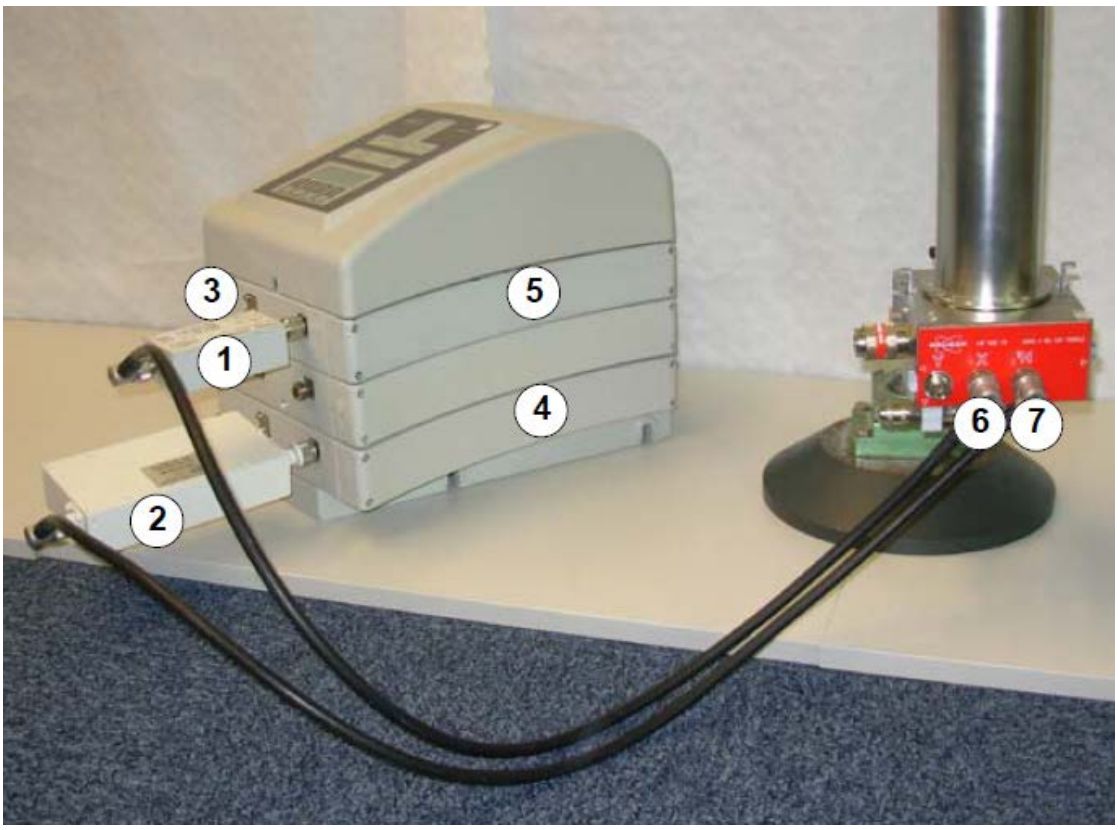


Figure 4.2: Probe Connections to the Preamplifier

1	X Low Pass Filter	5	13C Matching Box
2	Proton Bandpass Filter	6	X Probe Connector
3	X-BB Preamplifier	7	1H Probe Connector
4	1H HP Preamplifier		

The figure above shows the probe connections to the preamplifier with appropriate filters, placed on a table for better illustration.

4.2.2 Setting Acquisition Parameters

Create a new data set for the experiment by typing **edc** in the command line.

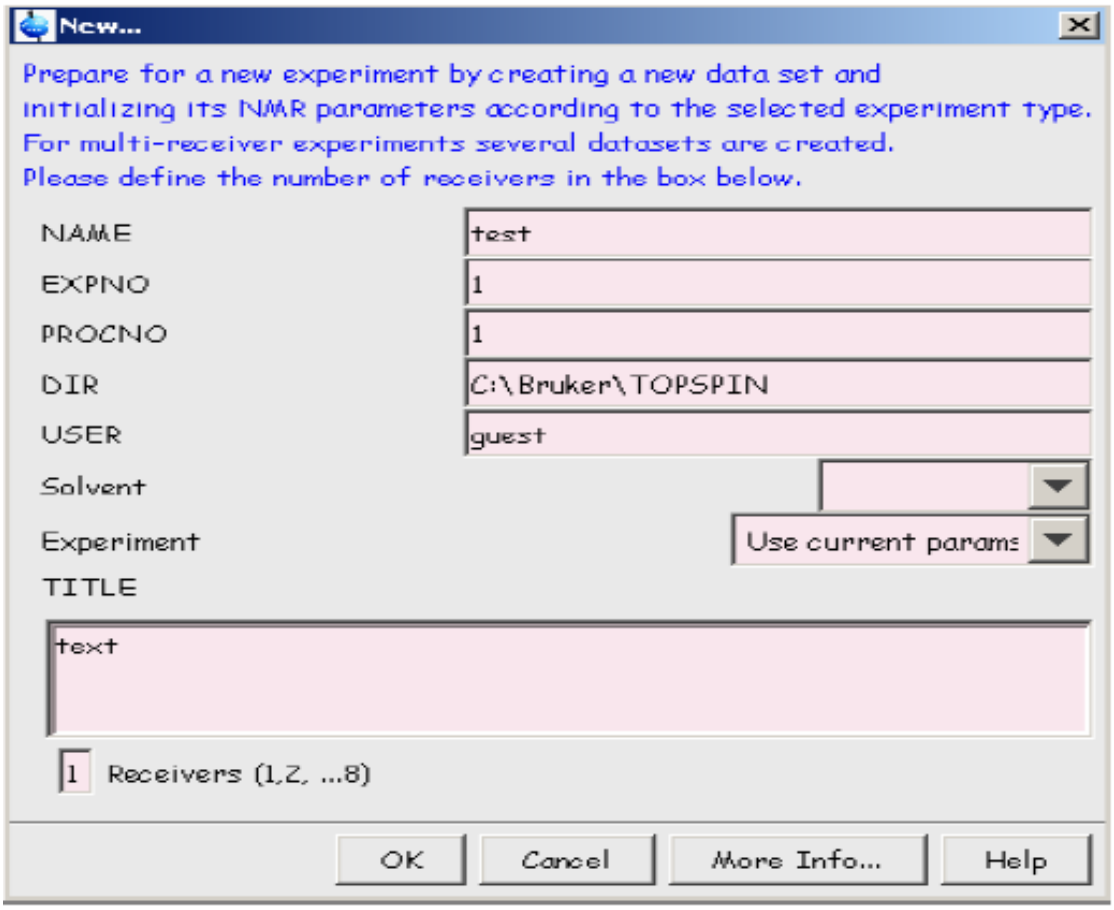


Figure 4.3: Pop-up Window for a New Experiment

Spin the KBr sample moderately (~5, 2.5 mm: 10 kHz). In order to set up the experiment, type **ased** in the command line to open the table with parameters used for this experiment.

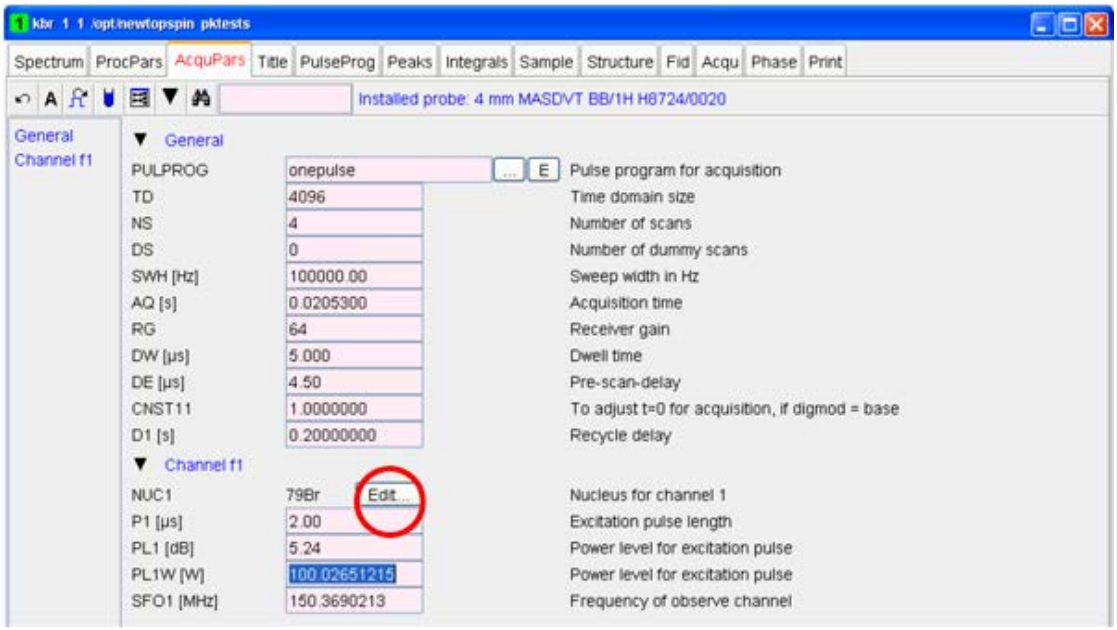


Figure 4.4: The ased Table with Acquisition Parameters for the KBr Experiment

Check the RF routing by clicking on the **Edit** button, or by typing **edasp** in the command line and check **powmod** by clicking the default button (as described above). The RF routing for this experiment is shown in the first figure above. Next set p1 = 2 µs, ns = 8 or 16. The power level at which p1 is executed is pl1. As high power transmitters are used it is important to be aware of the pulse power that is applied. With TopSpin 2.0 and later, **ased** shows **pl1** and **pl1w**, if the transmitter has been linearized (green dot in **edasp**) and the transmitter power has been measured. Set the power to about 100W.

For a non-linearized transmitter, pl1 should be set to 10 in case of a 1000W transmitter and to 4 (5) for a 300W or 500W transmitter. You can also check durations and power levels in a graphical display by clicking the experiment button in the Pulprog window as it is shown in the figure below.

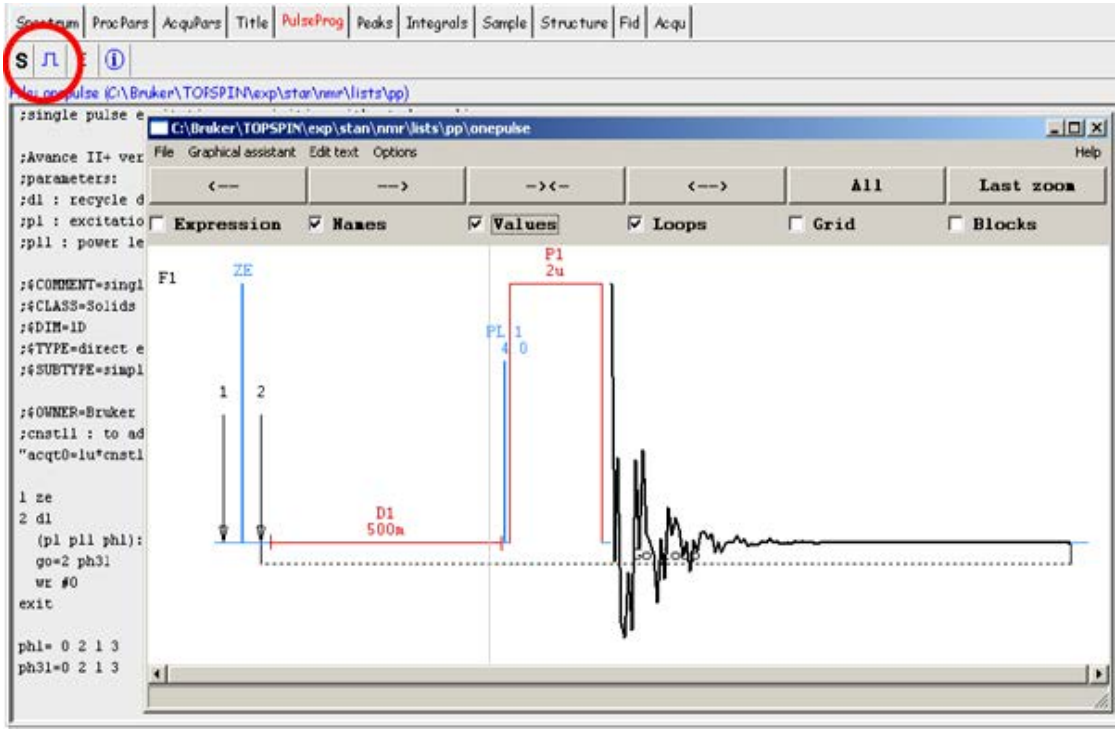


Figure 4.5: Graphical Pulse Program Display

In the figure above, the experiment button for opening the graphical display is marked with a red circle.

Match and tune the probe for this sample using the command **wobb**. This will start a frequency sweep over the range of SFO1+/-WBSW/2. The swept frequency will only be absorbed by the probe at the frequency to which it is tuned.

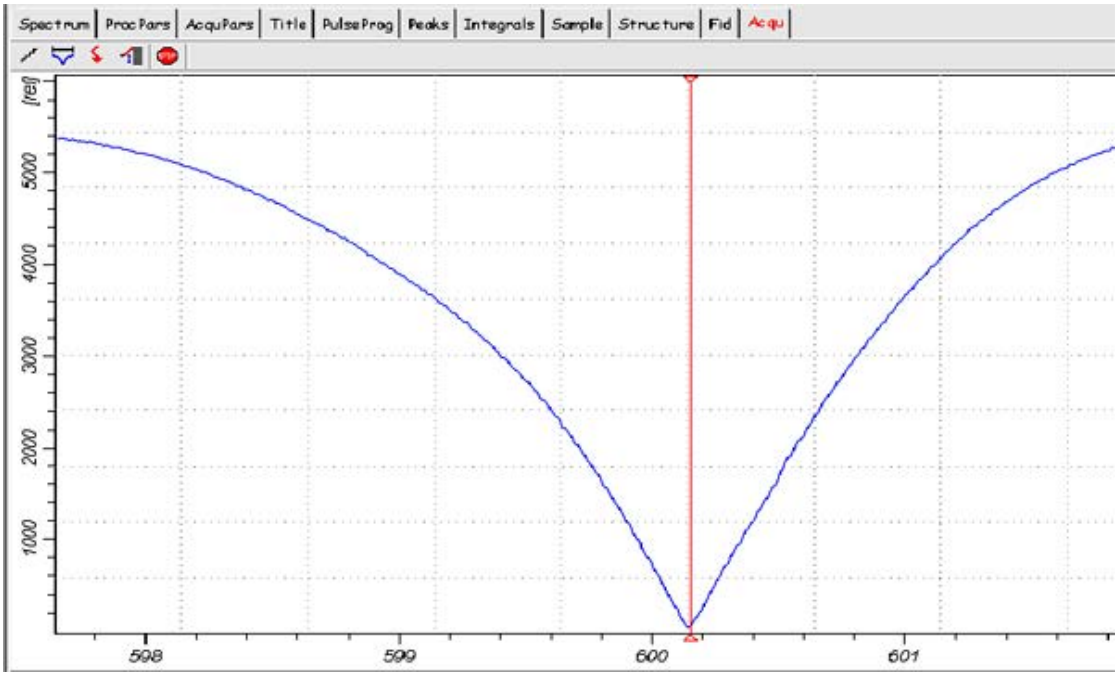


Figure 4.6: Display Example of a Well-tuned Probe

At frequencies, where the probe is not matched to 50 Ohms, the curve will lift off the zero line. If tuned to a frequency within SFO1+/-WBSW/2, but ≠ SFO1, the probe response will be off center.

Note: Fake resonances may appear which do not shift with probe tuning. It is always a good idea to keep track which nucleus was tuned last so it is clear what direction to tune to. Usually, turning the tuning knob counter clockwise (looking from below) will shift to a higher tuning frequency.

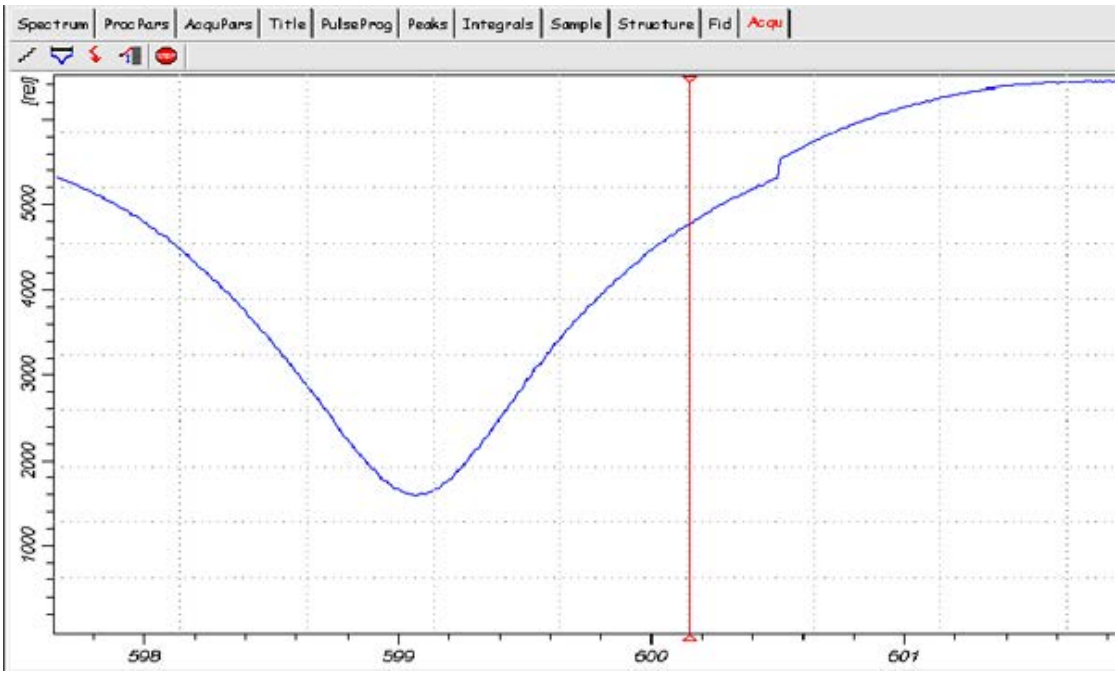


Figure 4.7: Display Example of an Off-Matched and Off-Tuned Probe

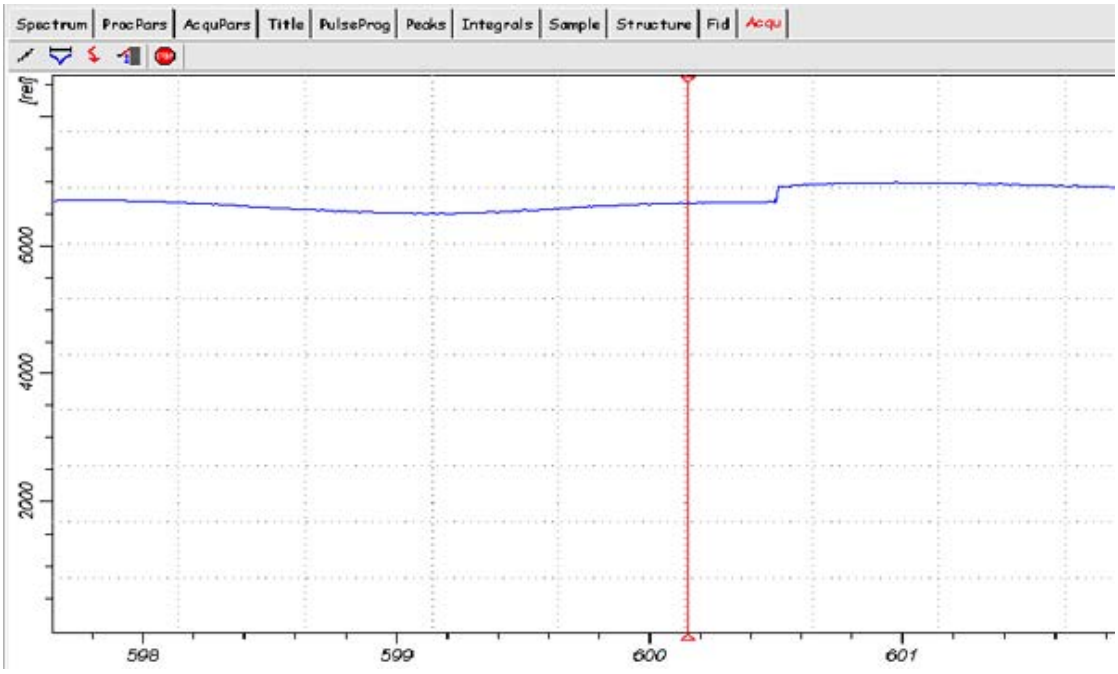


Figure 4.8: Display Example Where Probe is Tuned to a Different Frequency

The figure above is an example of where the probe is either tuned to a completely different frequency outside this window, or the probe is not connected to the selected preamp. Check **edasp** for a correct routing, then check for a correct matching box frequency range. Increase WBSW to 50 or 100 and try to find the probe resonance position.

When the probe is well tuned, start an acquisition by typing **zg** in the command line or clicking on the black triangle (upper left side) in the acquisition display. Do a Fourier transformation and a phase correction by typing **ft** and phase correct.

Set offset **O1** to the value obtained for the center peak and start **xau angle**. This will allow you to view the Fourier transformed spectrum or the FID after ns scans. The magic angle is adjusted best when the spikes on the FID or the spinning sidebands in the spectrum display have maximum size. This is most easily seen with the carrier exactly on resonance and the un-shuffled FID display mode. The spinning sidebands should have maximum intensity, the rotational echoes on the FID should extend out to at least 8 msec in the FID display.

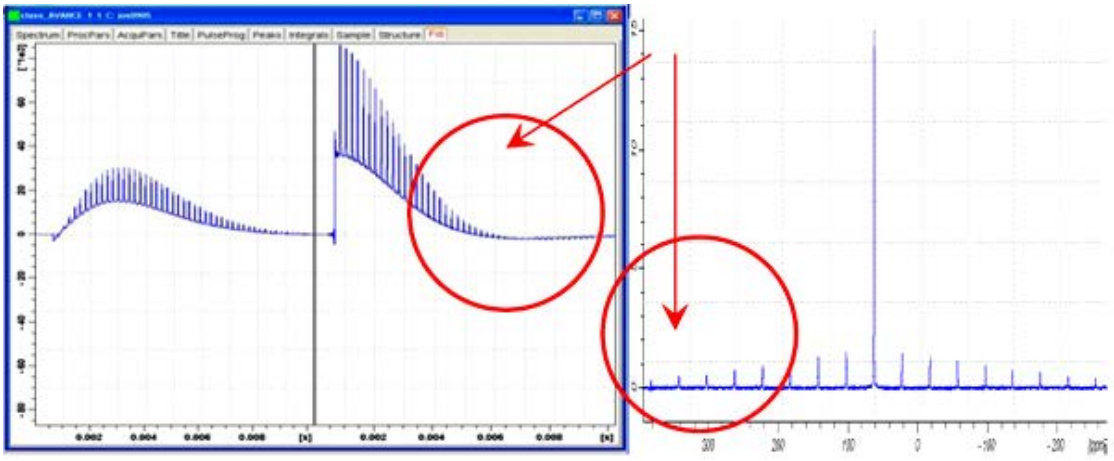


Figure 4.9: FID and Spectrum of the 79Br Signal of KBr used to Adjust the Magic Angle

4.3 Calibrating 1H Pulses on Adamantane

Spin the KBr sample down and change to a spinner filled with adamantane. Spin at 5-10 kHz. Generate a new data set from the KBr data set by typing new. Set the instrument routing for ^{13}C observe and ^{1}H decoupling, as shown in the following figure:

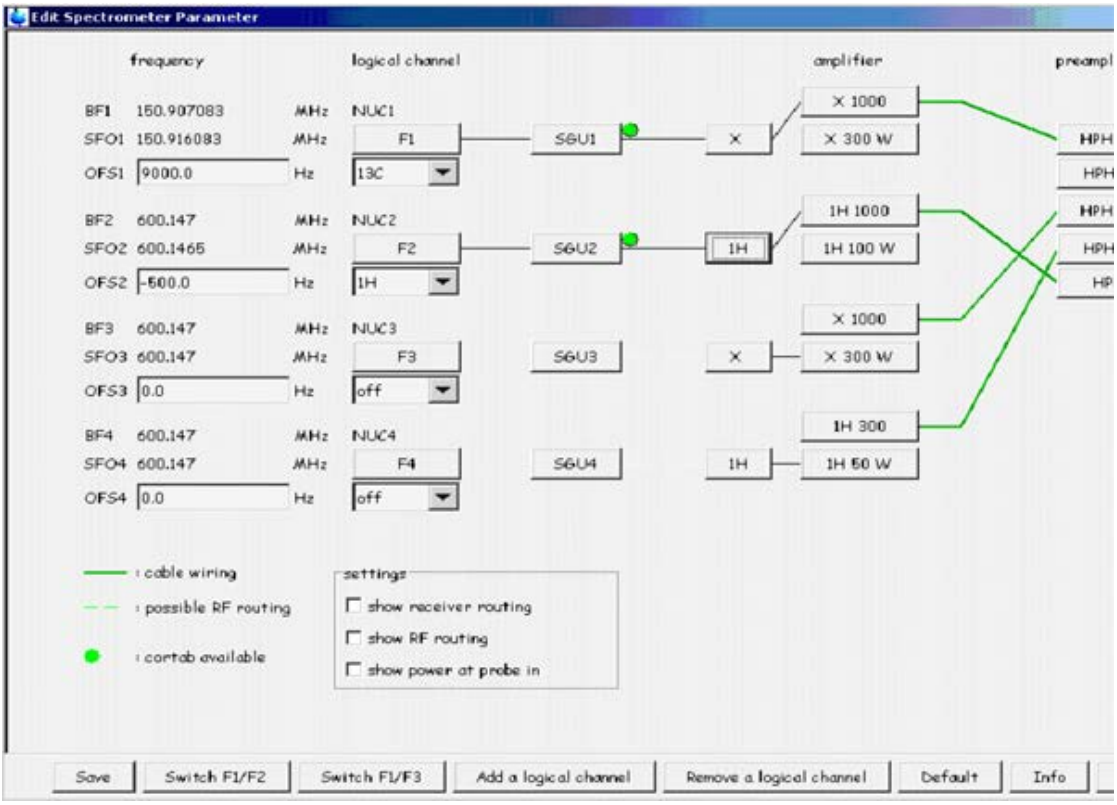


Figure 4.10: Routing for a Double Resonance Experiment using High Power Stage for H and X-nucleus

The figure above shows the routing for a double resonance experiment, e.g. a ^{13}C experiment with ^{1}H decoupling. For high power transmitters, the parameter **powmod** must be set to high. To check which power mode is selected, click the **Default** button, and change to **powmod high** in the command line if necessary. Note: the routing is only effective if the parameter powmod = high.

To change to proton observe, click **SwitchF1/F2**.

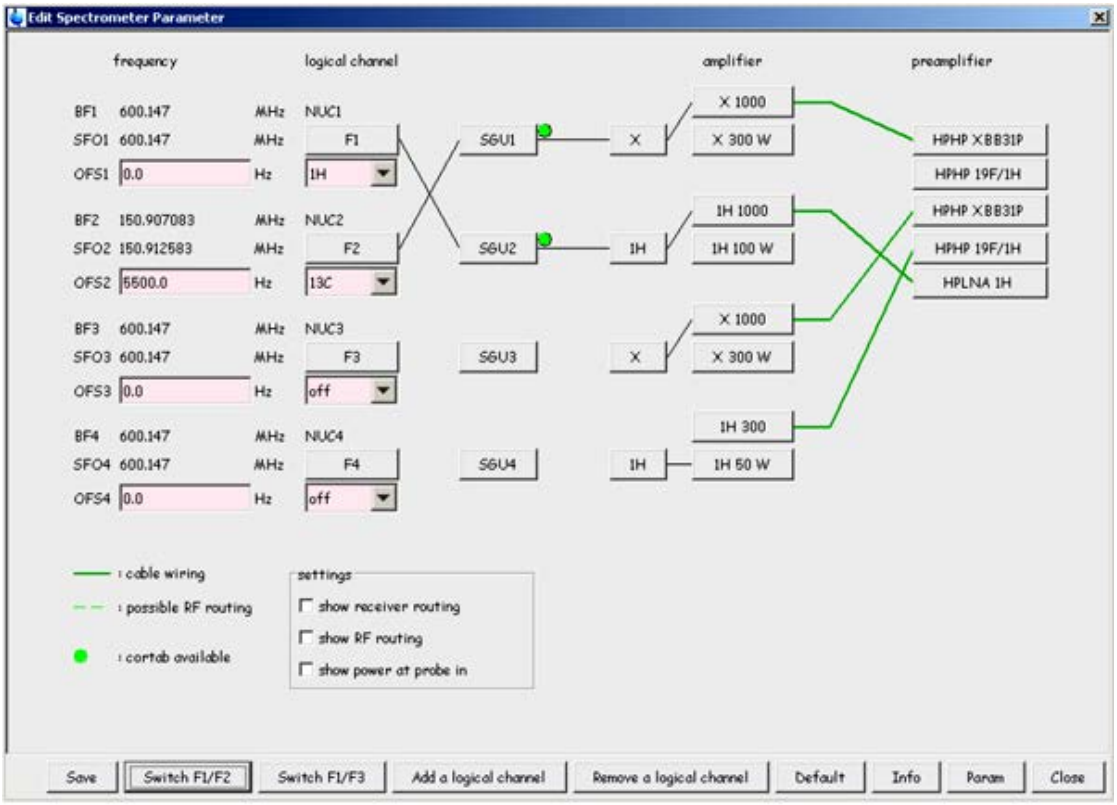


Figure 4.11: Routing for a Double Resonance Experiment, Changed for Proton Observation

In the figure above channel F2 need not be used.

The settings for F1 and F2 are interchanged. Change **rg** to 8-16 and **d1** to 4 sec. Set **pl1w** to 50W, or to 10 dB (high power proton transmitter), 7 dB (500W proton transmitter), 5dB (300W proton transmitter) or -4 (100W transmitter), if the green dot does not appear in the 1H channel in **edasp**. Connect the probe proton channel to the proton preamp. A proton band pass filter must be inserted between the preamp and probe. Tune the proton channel of the probe using the command **wobb high**. This means that the highest frequency is tuned first. Stop and type **wobb** again. Adjust the tuning of the X channel to ^{13}C. Alternatively, you can switch to the lower frequency channel within **wobb high** by clicking on the frequency table symbol in the wobb display or by pressing the second touch button on the preamp cover module twice. Then acquire ns =2 scans on the protons of adamantane.

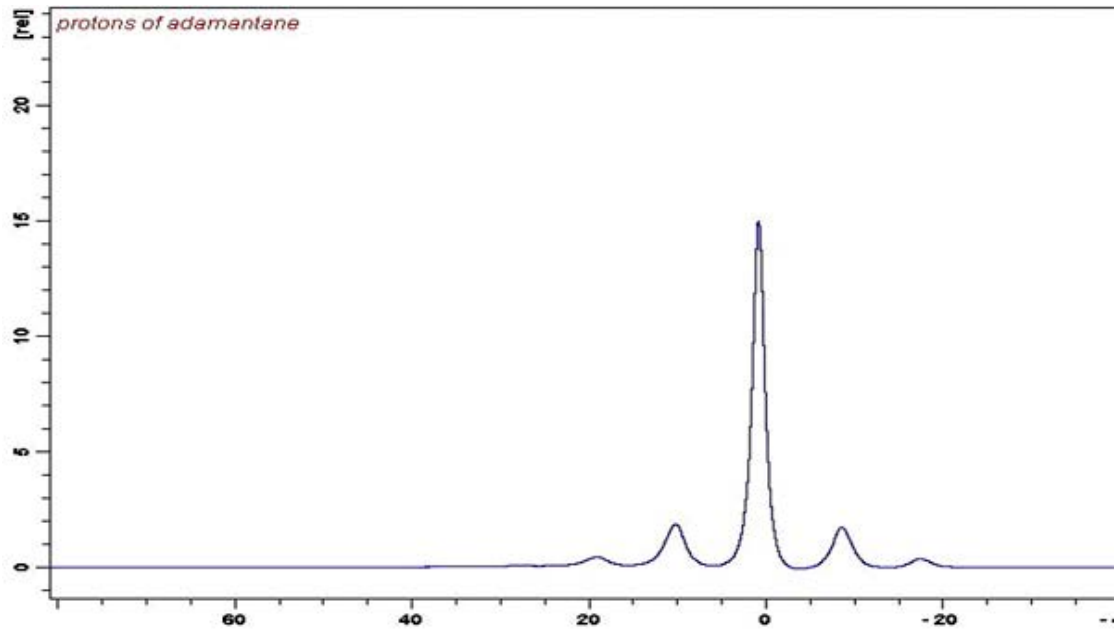


Figure 4.12: Proton Spectrum of Adamantane at Moderate Spin Speed

Set the carrier frequency O1 on top of the biggest peak using the encircled button in TopSpin.

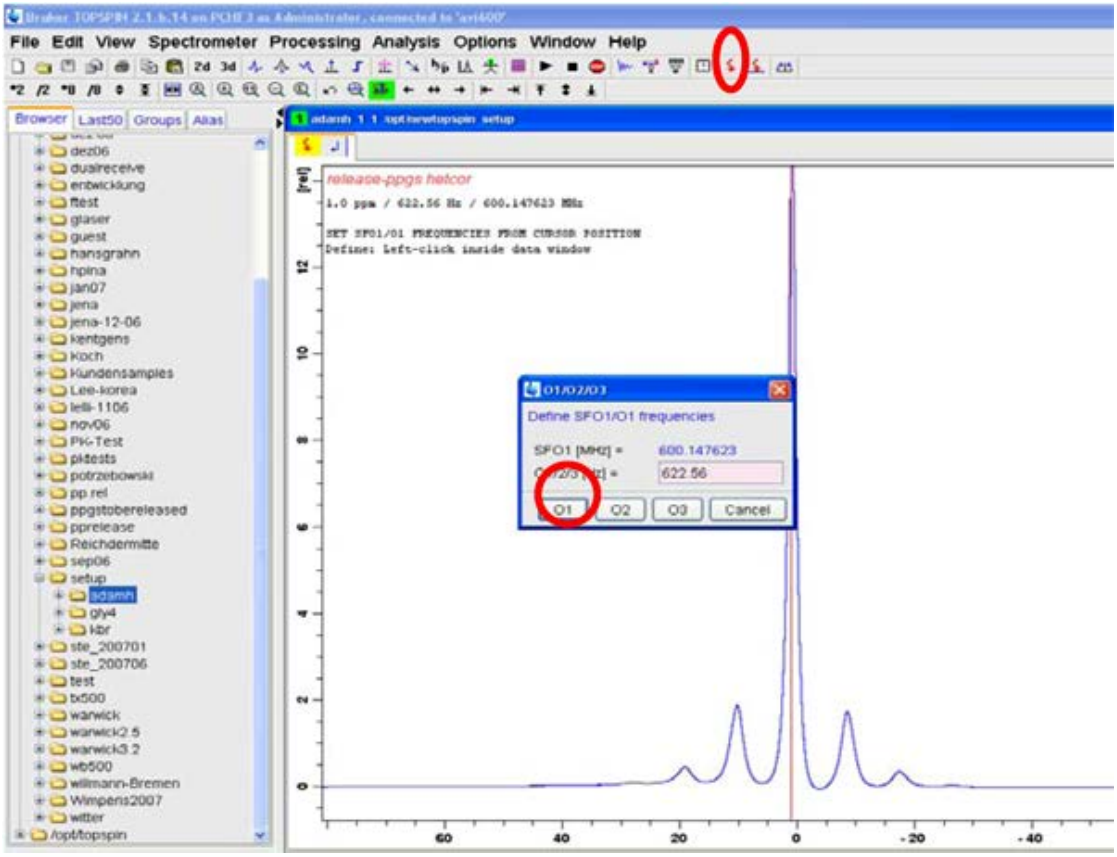


Figure 4.13: Setting the Carrier on Resonance

Click on the marked red arrow to set the observe frequency, set the position of the cursor line, and left click on the **O1** button. Acquire another spectrum, ft and phase.

Expand the spectrum around the adamantane proton signal including the spinning sidebands by clicking on the left margin of the region of interest and pulling the mouse to the right margin of the region of interest as shown in the following figure:

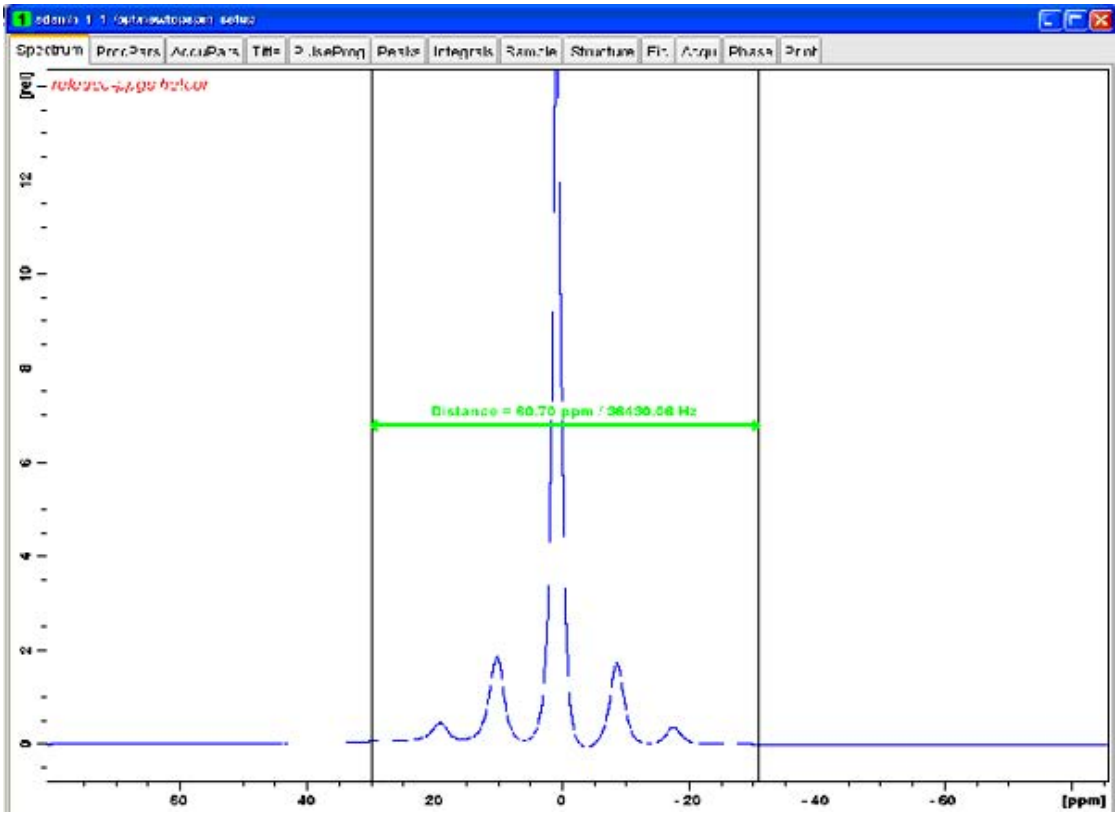


Figure 4.14: Expanding the Region of Interest

Right click in the spectrum window. When the Save Display Region to menu pops up, select Parameters **F1/2** and **OK** or type **dpl** in the command line.

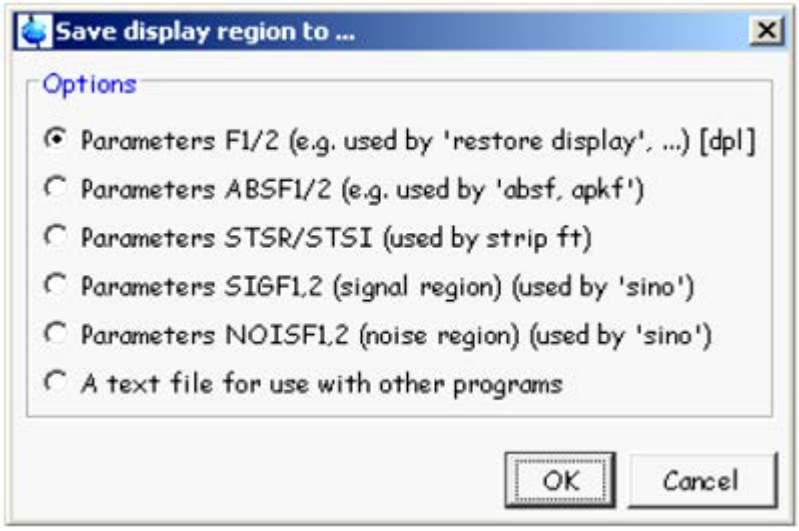


Figure 4.15: Save Display Region to Menu

Start parameter optimization by typing **popt** in the command line. The popt window will appear.

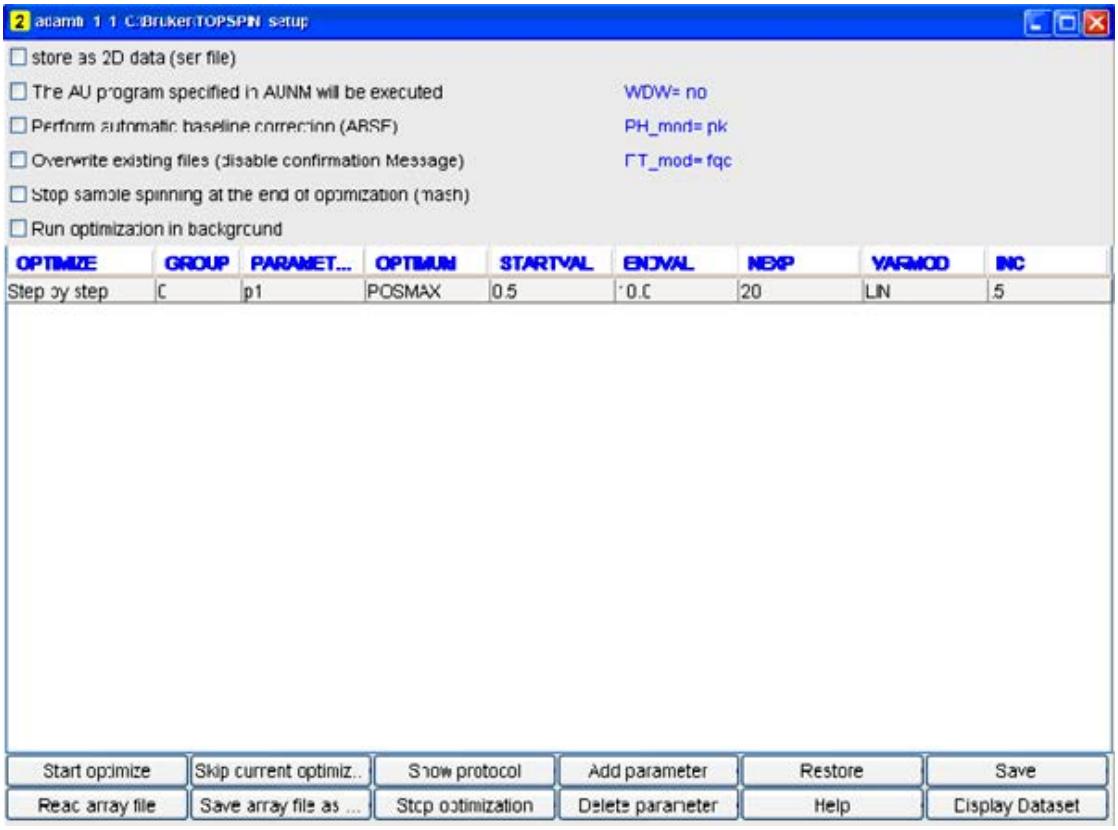


Figure 4.16: The popt Window

Use optimize **Step by step**, parameter **p1** to optimize parameter **p1**, optimum posmax to find the highest signal intensity (90 degree pulse) for the given value of pl1w or pl1,and varmod **LIN** to use linear increments for optimization. The value for group is not used for optimizing only one parameter and the number of experiments NEXP is set automatically when clicking on the **Save** button. Then save the table by clicking on the **Save** button and click on **Start optimize** to start the optimization procedure.

The parameter value obtained by the program is written into the parameter set of the actual experiment at the end of the optimization.

In order to stop the execution of popt use the **skip** or **stop optimization** buttons. Skip optimization will evaluate the obtained data as if popt had finished regularly and writes the parameter into the parameter set. Stop optimization will stop without evaluation of the data. You can also type **kill** in the command line and click on the bar with **poptau.exe** to stop optimization. This will work like stop optimization.

Popt will generate a data set, where the selected expansion part of the spectrum is concatenated for all different parameter values (in this case, for p1). It will have a procno around 999. To achieve this, processing parameters are changed appropriately. Fourier transforming a normal FID in such a window will generate an incorrect spectrum window.

Therefore:

Never start an acquisition in such a window, first read in the procno where popt was started using the **rep n** command where n is the source procno.

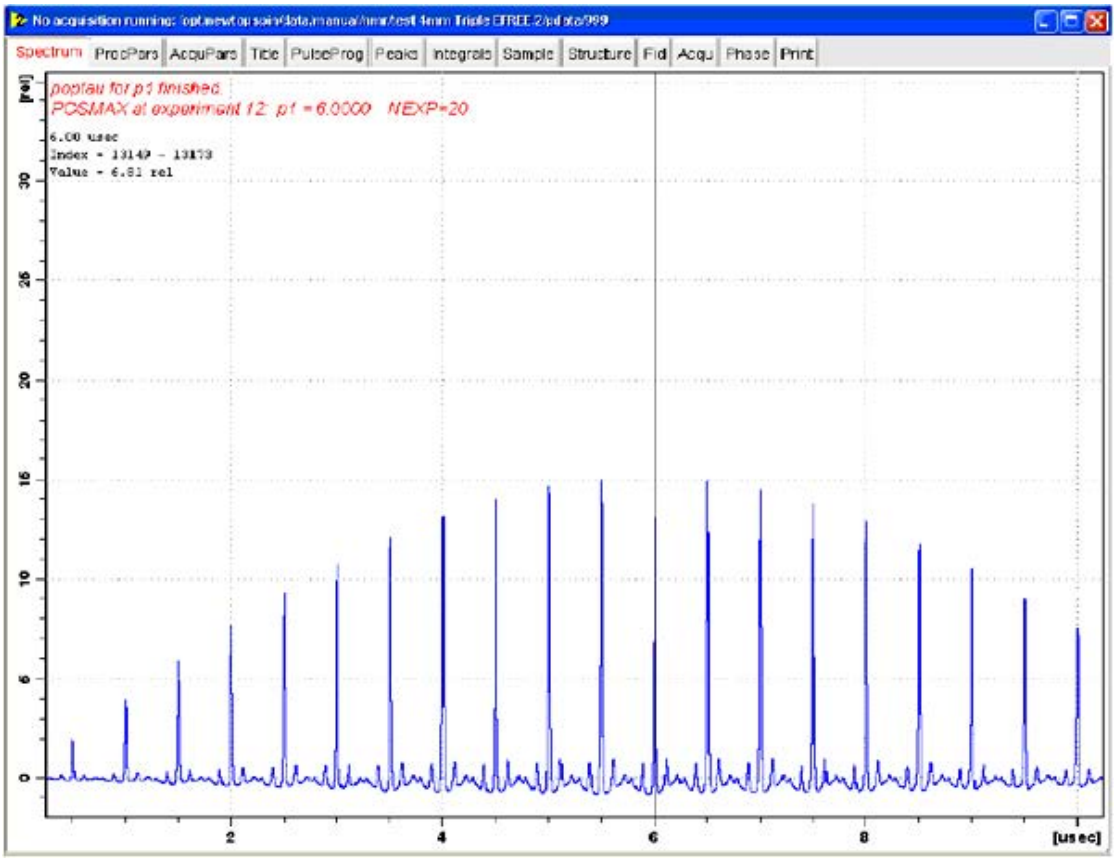


Figure 4.17: The popt Display after Proton p1 Optimization

The figure above shows the popt display after proton p1 optimization, the biggest signal is obtained at 6 µsec in this case.

Once you have obtained a 90-degree pulse for a given power setting, you can calculate power levels for different RF fields using the AU program **calcpowlev**. Type **xau calcpowlev** into the command line and follow the instructions in the popup window.

Calculate the power level (in Watts or dB) required to achieve a 4.5 µsec proton 90 degree pulse. In this case, 6 µsec were obtained. The command **calcpowlev** calculates a power level 2.5 dB higher than used above to achieve 4.5 µsec pulse length. Set the new pl1 as calculated and check whether 2*4.5 µsec for p1 will give a close to zero signal. This is a safe power level for all probes for pulses up to 100 msec. length.

4.4 Calibrating 13C Pulses on Adamantane and Shimming the Probe

A high power decoupling experiment on ^{13}C of adamantane is used to measure ^{13}C pulse parameters.

NOTE: For experiments where long decoupling pulses on protons are executed, the proton preamplifier must be bypassed, i.e. the transmitter should be wired to the probe directly (via the proton bandpass filter) without going through the preamp if a high power proton preamplifier is not available. For HP-HPPR modules $^{1}H/^{19}F$ this is not absolutely necessary, but recommended. For HPLNA 1H modules it is not required to bypass. Note that when bypassing the preamp which attenuates by about 1 dB, the proton power levels should be corrected by adding 1 dB to the pl-values.

Type **edasp** in the command line. You should get a display like that in the previous section. Click on **Switch F1/F2** to set for ^{13}C observation with proton decoupling. Load the pulse program hpdec. Set cpdprg2=cw. Set pl12 to the power level that yields a 4.5 µsec proton pulse. Set pl1 such that in ased the power displayed is 200W for ^{13}C (7 mm probe), 150W (4 mm probe) or 80W (2.5 mm probe). If the green dot is not visible in edasp for the ^{13}C channel, set pl1 to 12 dB (1 kW transmitter), 9 dB (500W transmitter) or 7 dB (300W transmitter) for any probe. Make sure the proton channel is tuned (wobb high) and the carbon channel is also tuned (wobb). With **d1** = 4s, **rg** = 256, **swh** = 100000, **td** = 4k, **o2** set to be on resonance on the adamantane protons as found above, accumulate 4 scans. Set the carrier frequency between both adamantane ^{13}C peaks. Reduce the spectral width **swh** to 50 kHz, set **aq** =50 msec.

Acquire 2-4 scans and define the plot limits for the larger of the two peaks. Define the plot limits and determine the 90 degree carbon pulse **p1**, using **popt**. Recalculate **pl1** for a 4.5 µsec carbon pulse using **calcpowlev**.

Pulse continuously using **gs** and shim the z gradient for highest FID integral.

The gradient settings can be conveniently changed in the setsh display. The next two figures show the adamantane 13C FID without shims, with z-shim adjusted and the corresponding setsh displays.

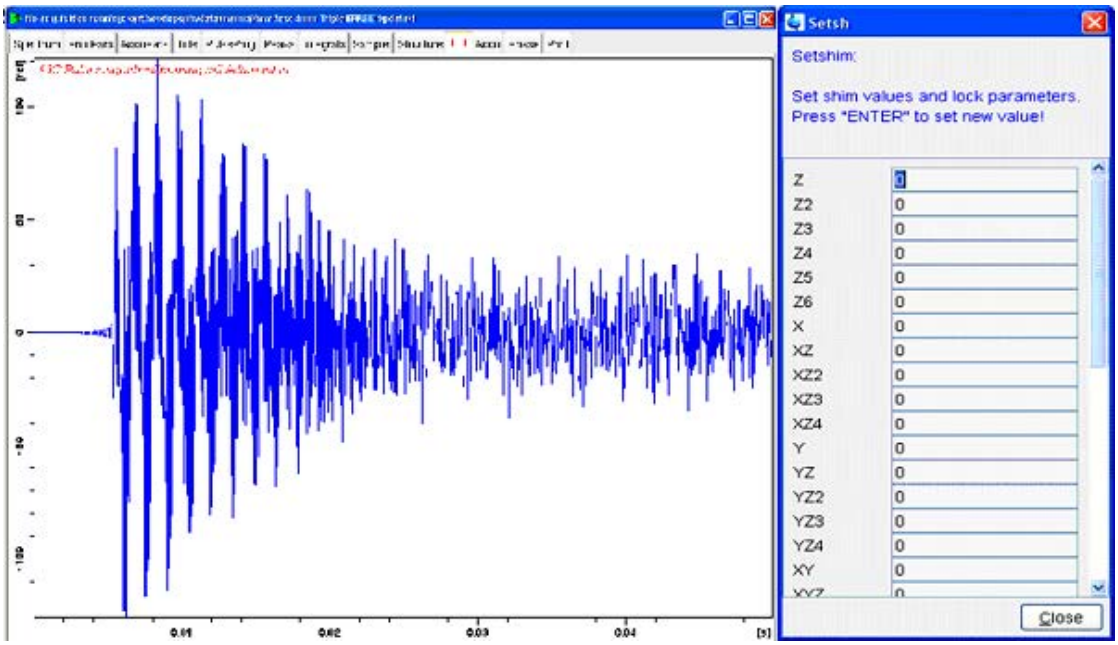


Figure 4.18: Adamantane 13C FID with 50 msec aq. setsh Display

The figure above is an Adamantane ^{13}C FID with 50 msec **aq**, with **setsh** display showing no shim values.

Note: Spinning removes part of the B0 in homogeneities. Probes which do not use susceptibility compensated coil wire can show much shorter T_2^* and require much more shimming effort (only with older probes up to 400 MHz proton frequency).

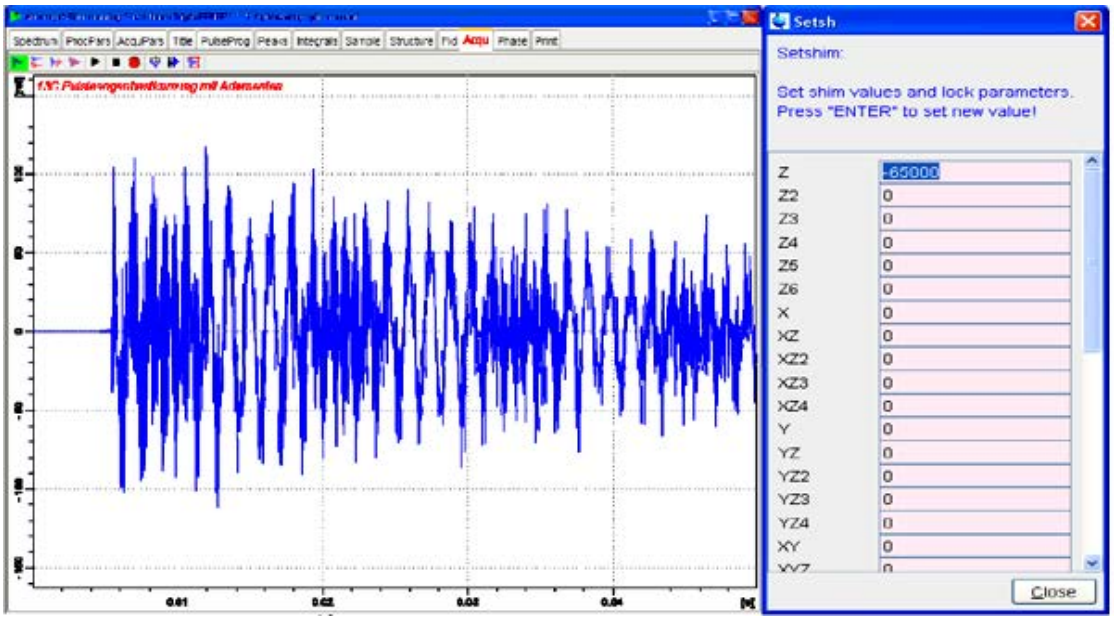


Figure 4.19: Adamantane 13C FID with 50 msec aq. setsh with Optimized Z-Shim Value

For optimum shims (rarely required) set the shims x, y, z^2, xy and x^2-y^2, as well as xz and yz. You may need to increase the acquisition time **aq** to see the effect of increasing resolution. Save the shims using the command **wsh** followed by a suitable name.

Before the shims are saved, it is recommended to reset the field value in the **bsmsdisp** menu to be exactly on resonance with your shift reference sample of choice (protons of adamantane or water in D_2O or silicon rubber).

This allows the command probe field (TopSpin version 3 and up) to set the field according to probe shims and magnet drift (see below).

Note: For long acquisition times (**aq** > 0.05 s) the decoupling power level pl12 must be set to +3 dB and d1 must be increased to 6s. To allow longer acquisition times than 50 msec, the **ZGoption –Dlacq** must be set in **ased**, if the pulse program contains the include file aq_prot.incl. Make sure the **–Dlacq** option is not left set for the following steps.

4.5 Calibrating Chemical Shifts on Adamantane

In TopSpin (as well as in the XWIN-NMR 3.5 release) the frequency list for NMR nuclei follows the IUPAC recommendations[1] for reference.

Set the ^{13}C low field signal of adamantane to 38.48 ppm. This will set the parameter **SR** which is used to calculate the chemical shift axis and the peak positions in the spectrum.

Note: All data sets generated from this data set will have the peak positions correctly calibrated, if the magnetic field B_0 is not changed. However, you must make sure that the magnetic field is always the same. It may change, if the magnet has a slight drift, or if different shim settings are loaded. Therefore the same shim file should be loaded and the field be set to the same value using the **bsmsdisp** command. If the magnet drift is noticeable, the calibration should be redone in suitable intervals and the field value recorded in the lab notebook.

[1] R.K. Harris, E.D. Becker, S.M. Cabral de Menezes, R. Goodfellow and P. Granger, NMR Nomenclature. Nuclear Spin Properties and conventions for Chemical Shifts, Pure Appl. Chem. Vol. 73, 1795-1818 (2001)

One can also use a spinner filled with H_2O to set the field position more precisely. Do not spin the sample and make sure the cap is well fitted. Set **o1p** to 4.85 ppm, set for proton observe (as described above for adamantane), and use **gs** for continuous pulsing and FID display. Change the field value in **bsmsdisp** until the FID is exactly on resonance. Then all spectra taken should be correctly referenced with **sr** = 0.

For all these experiments the field sweep must be off! When the BSMS unit is turned off and on again, the sweep will always be on. Running spectra with the sweep on will superimpose spectra at different fields! One can set the sweep amplitude to 0 in order to avoid such an accidental error condition.

4.6 Setting Up for Cross Polarization on Adamantane

Cross polarization is used to enhance the signals of X-nuclei like ^{13}C. The strong proton polarization is transferred (cross polarized) to the X-nuclei coupled to the protons via strong dipolar couplings. To achieve this, the protons and the X-nuclei must nutate at the same frequency. This frequency is the RF field applied to both nuclei at the same time (contact time). If this condition (Hartmann-Hahn-condition) is met, the transfer of proton magnetization to carbon is optimum. Since the proton signal of adamantane is resolved into spinning sidebands even at slow spin rates, this Hartmann-Hahn condition can be set to match for every proton spinning sideband. Using a ramp for the proton contact pulse, the Hartmann-Hahn match is swept over these possible match conditions and becomes insensitive to miss-sets and different spin rates.

Start from the data set used for observing ^{13}C under proton decoupling (1.4). Load the pulse program **cp** (in **eda** or typing **pulprog cp**). The pulse sequence is depicted in the following figure:

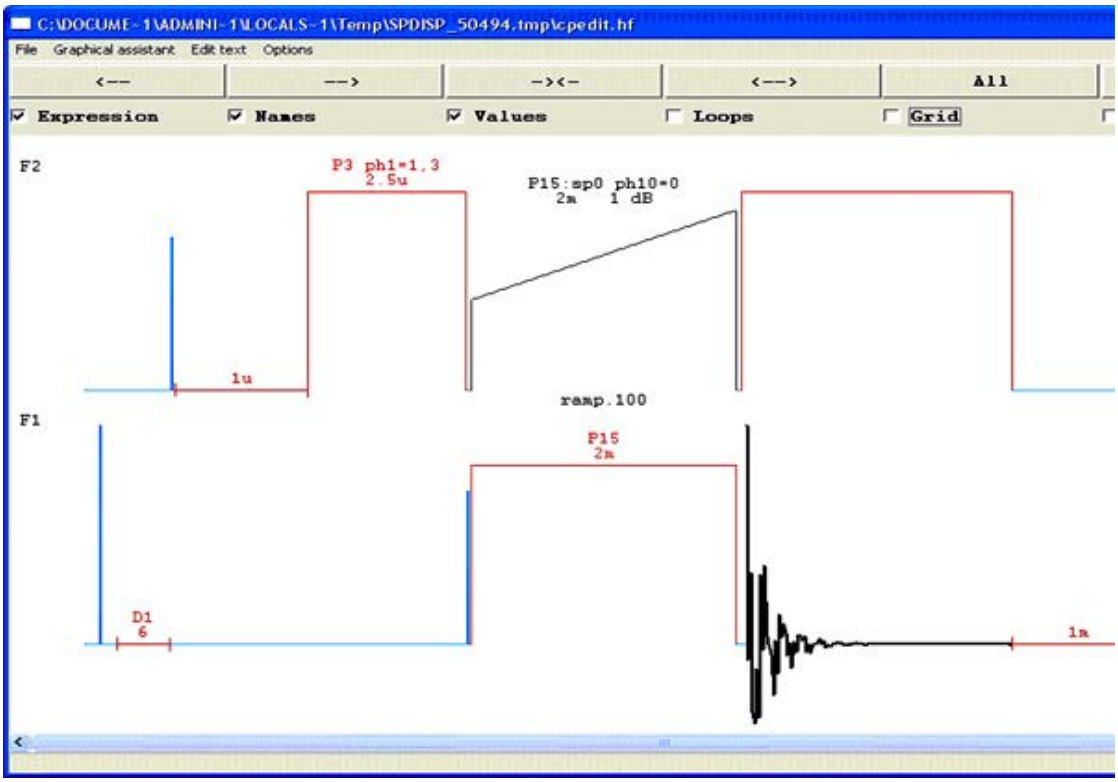


Figure 4.20: A cp Pulse Sequence

The following parameters are set:

- **pl12**, for the initial 90 degree pulse and the decoupling during acquisition: set for a 4.5 µsec proton 90 degree pulse (as previously determined).
- **pl1**, for the carbon contact pulse, set for a 4.5 µsec carbon pulse (as previously determined).
- **p3**, 4.5 µsec.
- **spnam0**, set to ramp.100 to sweep the proton contact RF field from 50 to 100%.
- **sp0**, set to pl12 -3 dB, to account for the lower average RF over the ramp.
- **p15**, 2-5 ms (after the value, specify m to make it milliseconds, otherwise it is taken as microseconds).
- **cpdprg2**, select **cw**.
- **o1**, set between both adamantane peaks.
- **o2**, set to be on resonance on adamantane protons.

Acquire 2 or 4 scans, then set plot limits for both peaks, and optimize p3 (+/- 2 µsec) and pl1 (+/- 2 dB) for best signal. The next figure shows a Hartmann-Hahn match optimization over 4 dB using a ramp contact pulse going from 50 to 100% amplitude.

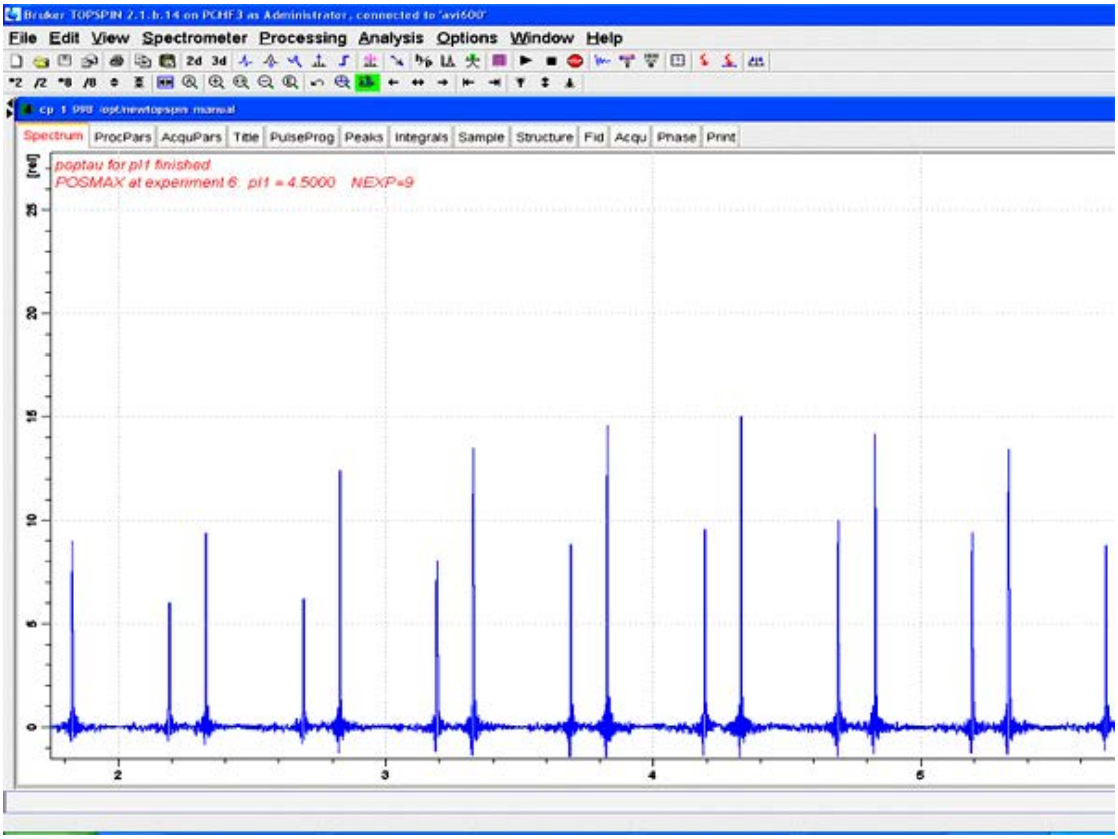


Figure 4.21: Hartmann-Hahn Optimization Profile

The Wiggles Profile besides the signals stem from truncation of the FID after 50 ms acquisition time.

To exemplify the existence of several HH-conditions on a spinning adamantane samples, another HH profile (see next figure) is shown where a square proton contact pulse is used. There are several maxima corresponding to matches on the sideband orders n+2, n+2, n+1, n+0, and n-1. The largest intensity is seen for n+/-1, the intensities are very sensitive to the RF-level which is varied in 0.2 dB steps. Using a ramp makes the experiment much more stable and more quantitative. Problems may arise if the proton T1ρ is short, since usually longer contact times must be employed. Therefore it makes sense to use a flatter ramp (70-100%) and optimize for the spin rate which is used.

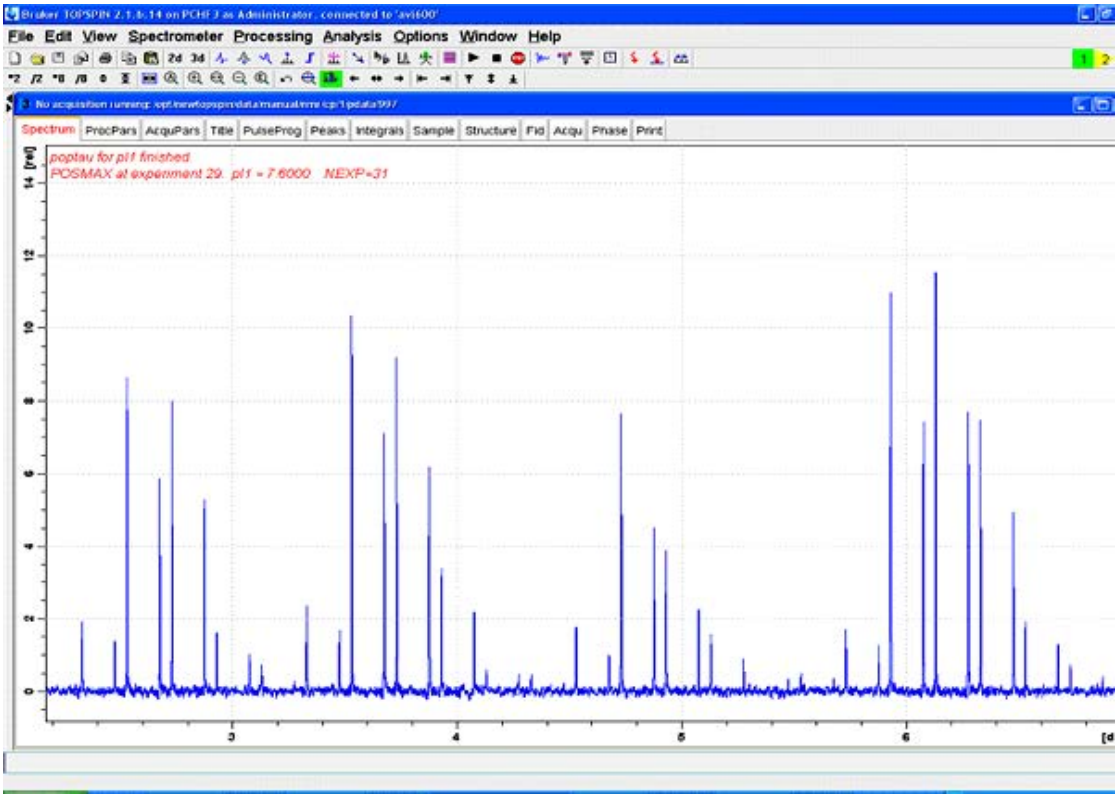


Figure 4.22: Hartmann-Hahn Optimization Profile Using a Square Proton Contact Pulse

The sideband order 0 at 4.8 dB gives a rather small intensity. A ramp sweeping over 3.5-6.5 dB would cover both most efficient HH conditions. Note that increasing the spin rate would shift all maxima except the one at 4.8 dB further out.

4.7 Cross Polarization Setup and Optimization for a Real Solid: Glycine

Adamantane is highly mobile even in the solid state. Therefore it behaves differently from a “hard” solid like glycine. For instance, it is not sensitive to incorrect decoupling adjustments, and is also not sensitive to incorrect settings of the magic angle. It is however extremely sensitive to incorrect HH adjustment. Glycine is therefore used for fine tuning of the decoupling parameters and signal-to-noise assessment. Start with the parameters found for adamantane, using a 50-100% ramp (ramp.100) and **p15**=2 ms for contact, **aq** = 20 msec. Change the sample from adamantane to glycine.

Since glycine may exist in two different crystal modifications with very different CP-parameters, and since packing of the spinner determines crucially the achievable S/N value, it is useful to prepare a reference spinner with pure α-glycine, finely powdered and densely packed. α-glycine is prepared by dissolution of glycine in distilled water and precipitation with acetone, quick filtering and careful drying in a desiccator. Drying is important because wet glycine may readily transform, especially when kept warm, into γ-glycine. α-glycine has two carbons with shifts of 176.03 and 43.5 ppm. γ-glycine shows resonances somewhat shifted to higher field, sharper lines, longer proton T_1 and shorter proton $T_{1\rho}$ which results in longer experiment time and less signal to noise.

Spin the glycine sample at 5 kHz (7 mm spinner), or 10 kHz (smaller spinners 4, 3.2 or 2.5), tune and match the probe.

The glycine CP/MAS ^{13}C-spectrum taken under the same conditions as adamantane previously, will look like in the following figure, far from optimum:

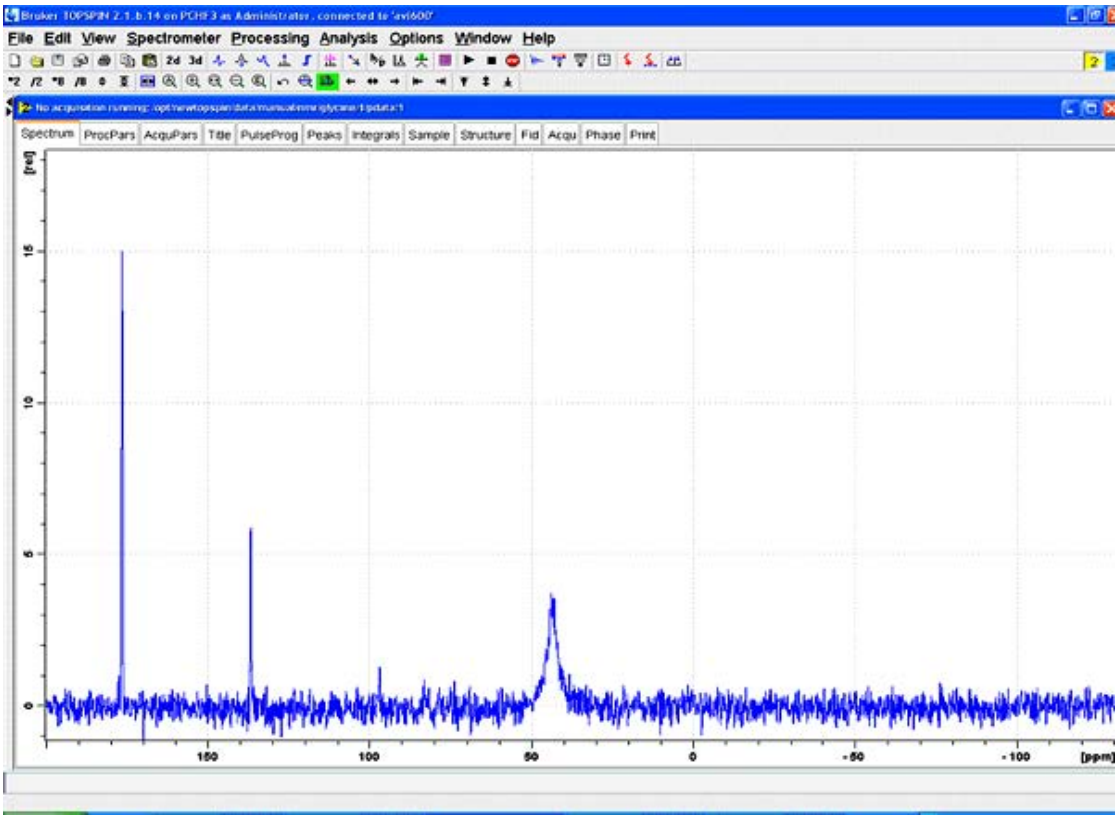


Figure 4.23: Display Showing α-Glycine Taken Under Adamantane Conditions, 4 scans

The figure above shows α-glycine taken under adamantane conditions, 4 scans: Incorrect carrier setting, α-carbon at 43 ppm insufficiently decoupled. Angle is set correctly, because carboxyl peak at 176.03 ppm shows a narrow lorentzian line shape. HH condition looks okay.

Reset the carrier, **o1p** should now be around 100 ppm, in the middle of most carbon spectra. Acquire a spectrum, set the plot limits for the peak at 43 ppm, and start **popt**, optimizing **o2** for maximum signal (+/- 2000 Hz around the current position) in steps of 500 Hz. The following result will be obtained:

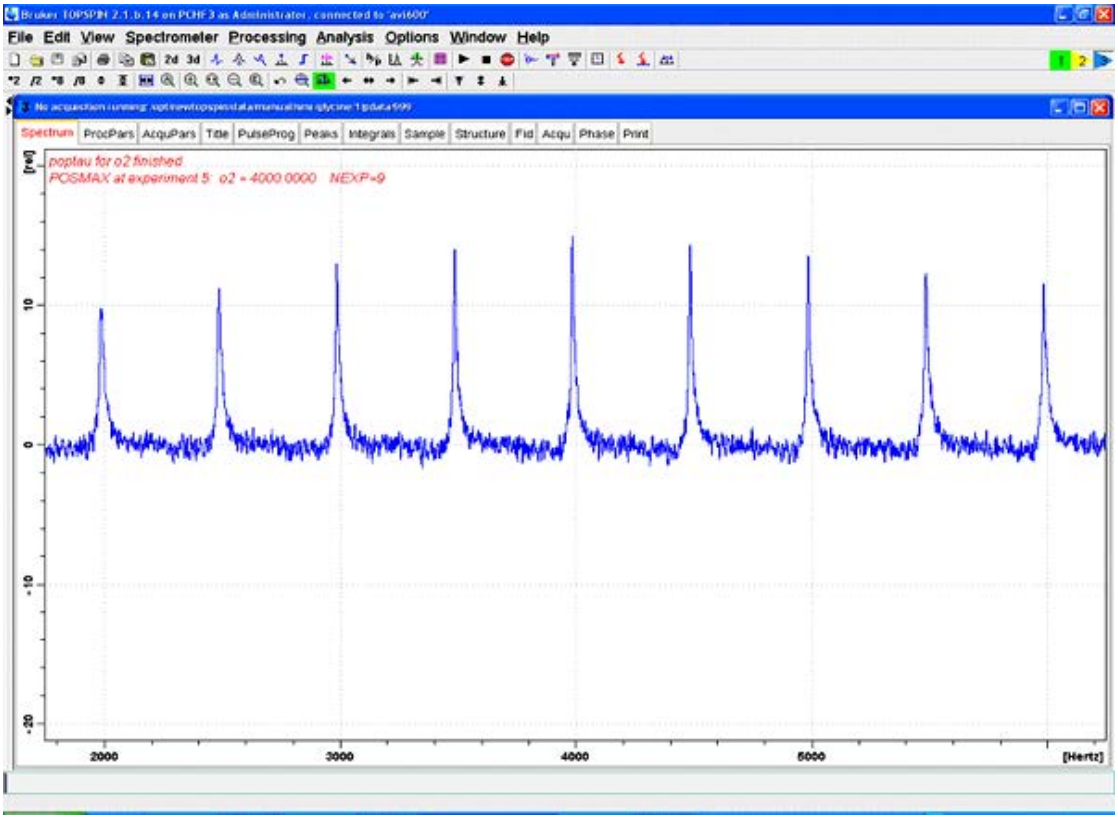


Figure 4.24: Optimization of the Decoupler Offset o2 at Moderate Power, Using cw Decoupling

Since the proton spectrum of glycine extends around 5 ppm, the optimum decoupler offset will be obtained at higher frequency than the adamantane proton peak (around 1.2 ppm). Decoupling is still inefficient, since cw decoupling is used which does not cover the whole proton shift range. Also decoupling power is too low with a proton pulse of 4.5 µsec. Glycine requires about 90 kHz of decoupling RF, corresponding to a 2.7 µsec proton 90 degree pulse. This can be obtained with probes of 4 mm spinner diameter and smaller (2.5, 3.2 mm). For a 7 mm probe, 3.5 (4 µsec) can be expected at proton frequencies below 500 (at 500) MHz. Use **calcpowlev** to calculate the required power level **pl12** and set **p3** to twice the expected proton pulse width. Check with 4 scans whether a close to zero signal is obtained. Compared to 4.5 µsec, a 2.7 µsec pulse requires about 4.5 dB more power (corresponding to almost 4 times more power!!!).

With **p3** properly set, a spectrum like the one below should be obtained, with about 93 kHz decoupling RF field.

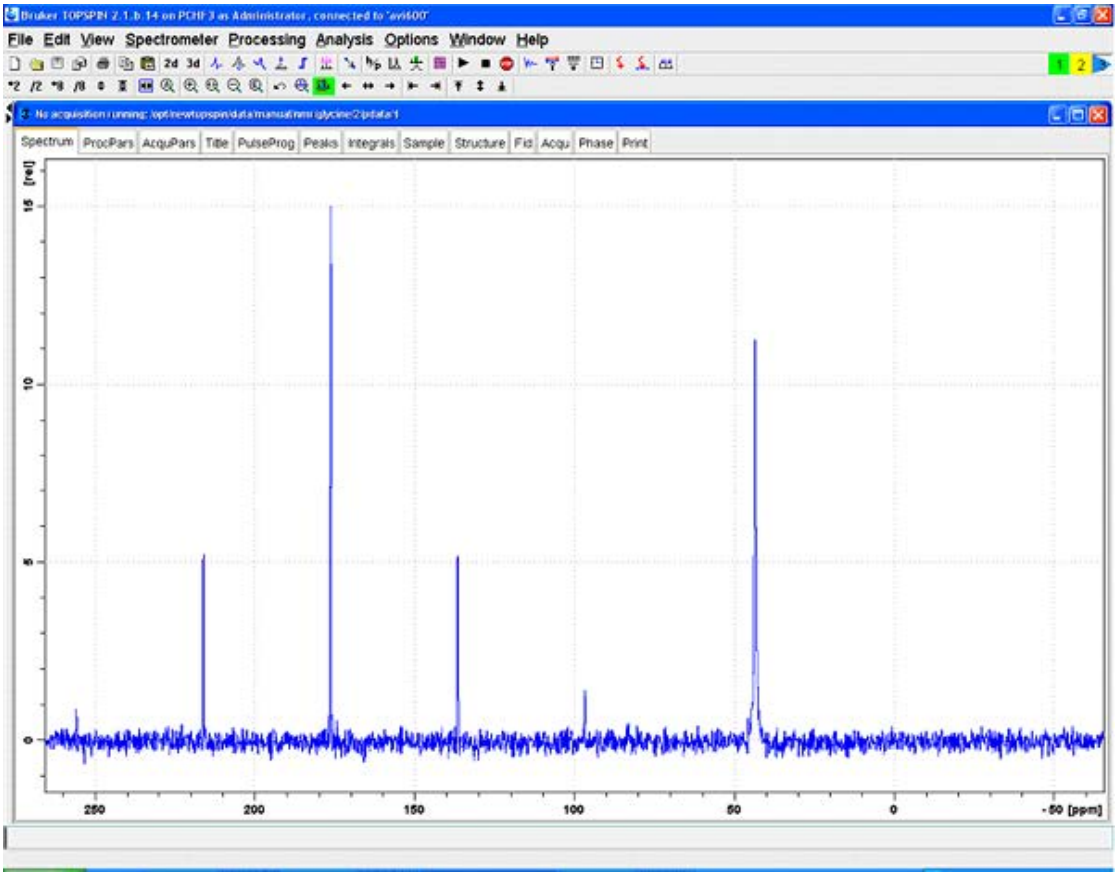


Figure 4.25: Glycine with cw Decoupling at 90 kHz RF Field

In the spectrum above, a lorentzian deconvolution (**Analysis** menu) shows a line width of 71 Hz for the peak at 43 ppm. The line width achievable under optimum decoupling conditions varies with the magnetic field. At fields below 9.4 Tesla (400 MHz) this line is substantially broadened by second order quadrupolar interaction to ^{14}N. At fields above 9.4 Tesla (500 MHz and higher), the residual line width is mostly determined by chemical shift dispersion and insufficient decoupling. Here less than 60 Hz (at 600 MHz) are expected. More efficient decoupling schemes must be applied especially at higher magnetic fields. A more efficient decoupling scheme is spinal64. Select **cpdprg2** = spinal64, set **pcpd2** to proton 180 degree pulse – 0.2 µsec for a start. A glycine spectrum as shown in *Figure 4.26 [▶ 73]* is obtained.

Parameter	Value	Comments
PULPROG	cp	cp.av for AV1 and 2
NUC1	13C	Nucleus on f1 channel
O1P	100 ppm	^{13}C offset
NUC2	1H	Nucleus on f2 channel.
D1	4 s	Recycle delay.
NS	4	Number of scans.
SWP	300 ppm	Spectral width for Glycine.
TD	2048	Number of acquired complex points.
CPDPRG2	SPINAL64	Decoupling scheme f2 channel (^{1}H).
SPNAM0	ramp.100 or ramp 70100.100	For ramped CP.
P15	2 ms	Contact pulse (f1 and f2 channel).
PL1		Set for 4-4.5 µsec P90.
SP0 (or pl2 AV1+2)		Set for 4-4.5 µsec P90 – 2 dB (optimize).
PL12		High power level f2 channel (^{1}H) excitation and decoupling.
P3		90°^{1}H pulse at PL12 (f2 channel).
PCPD2 or P31 (AV1+2)		SPINAL64 decoupling pulse.
O2P	2.5 - 3 ppm	^{1}H offset - optimize in 400 - 500 Hz steps for maximum signal of aliphatic peak.

Table 4.1: Summary of Acquisition Parameters for Glycine S/N Test

Note that the spectral window (**swp**) is set in ppm which makes the acquisition time dependant on the B_0-field at a given **td** of 2k. This is intentional and accounts for the line widths dependence on the B_0-field. The glycine lines show a broadening proportional to B0 due to chemical shift dispersion. To make S/N values more comparable, this accounts for shorter T_2 at higher field.

Parameter	Value	Comment
SI	2-4 k	Twofold or fourfold zero filling.
WDW	No	No apodization used for S/N measurement In this case.
PH_mod	Pk	Phase correction if needed.
BC_mod	Quad	DC offset correction on FID.
FT_mod	Fqc	

Table 4.2: Processing Parameters for the Glycine S/N-Test

Note: No line broadening is applied since the acquisition time is set appropriately.

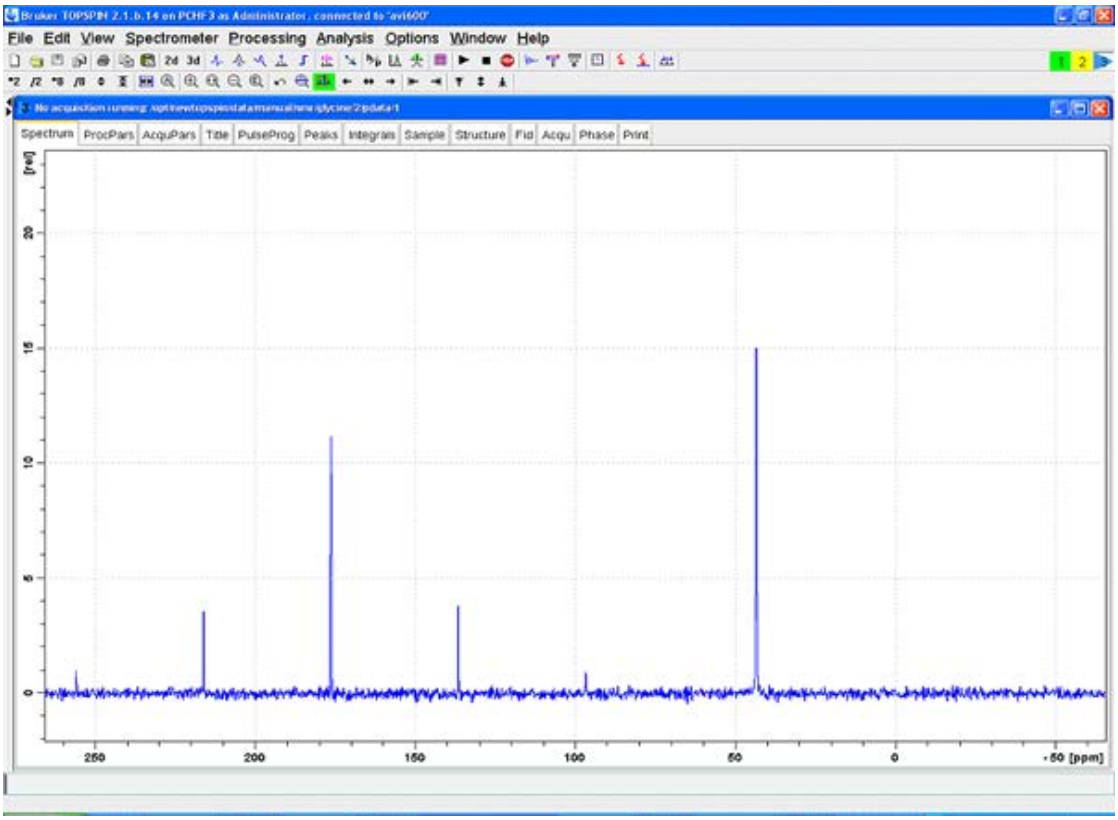


Figure 4.26: Glycine Spectrum with Spinal64 Decoupling at 93 kHz RF field

Here, the line width of the line at 43 ppm is fitted to be about 50 Hz: Correspondingly, the intensity is much higher. The **sinocal** routine calculates 80:1 S/N, using 250 to -50 ppm spectrum range, 50 to 40 signal range and 10 ppm noise range. On this triple probe, more than 100:1 is expected. What else needs to be optimized? Two more parameters are essential:

1. The power level at HH contact
2. The decoupling pulse **pcpd2**.

The spectrum in the previous figure was taken at contact power levels as set for adamantane. Furthermore, a 50% ramp was used, which has a rather low average RF level corresponding to about 5.7 µsec in this case (25% less than 4.5 µsec). This does not spin lock the protons well enough. So the power level of the contact needs to be increased. Set **spnam0** = ramp70100.100. Set **sp0** and **pl1** to about 2 dB less attenuation and check S/N again. Re-optimize the HH condition observing the peak at 176 ppm (which is less strongly coupled to protons and therefore exhibits a sharper HH matching condition) in steps of 0.3 dB. In this case, S/N improves to 100:1. Then optimize the decoupling pulse **pcpd2** in steps of 0.2 µsec, observing the peak at 43 ppm (which is more sensitive to decoupling mis-sets). Here, this led to another 10% improvement in S/N.

Good Laboratory Practice requires that evaluation measurements be taken in suitable periods. Store the optimized glycine spectrum together with the following important information:

1. Value of field setting.
2. Name of the shim file.
3. Name of the operator.
4. Probe setup (triple mode or double mode, high range or low range setting, WB probes only, name or part number of the probe).

5. Description of the sample (which reference rotor, weight of glycine and spinner).
6. Any additional comments, for instance the reading of the micrometer setting for the X tuning adjustment (not available on all probes).

Write this information into the title file so it is stored with the data set as well as all other acquisition and processing parameters. Recalling this data set and acquiring a new data set should give the same spectrum within +/- 10% of S/N.

4.8 Some Practical Hints for CPMAS Spectroscopy

Some general recommendations for reasonable RF-fields used in WB probes:

Probe	Nucleus	Decoupling power over 50 ms, 200 ms, 500 ms. Contact pulse up to 10 ms
2.5 mm CPMAS double resonance 35 kHz max sample rotation	^{1}H	115 kHz (2.2µs 90º pulse), 75 kHz, 40 kHz 71 kHz (3.5 µs) contact
2.5 mm CPMAS double resonance 35 kHz max sample rotation	^{13}C	83 kHz (3 µs 90º pulse) 71 kHz (3.5 µs)
3.2 mm CPMAS double resonance 24 kHz max sample rotation	^{1}H	110 kHz (2.3 µs), 60 kHz, 35 kHz 68 kHz (3.7 µs)
3.2 mm CPMAS double resonance 24 kHz max sample rotation	^{13}C	78 kHz (3.2 µs) 68 kHz (3.7 µs)
4 mm CPMAS double resonance probe (15 kHz max. sample rotation)	^{1}H	92.5 kHz (2.7us 90º), 50 kHz, 30 kHz 62 kHz (4 µs)
4 mm CPMAS double resonance probe (15 kHz max. sample rotation)	^{13}C	71 kHz (3.5 µs) 62 kHz (4 µs)
4 mm CPMAS triple resonance probe (15 kHz max. sample rotation)	^{13}C	66 kHz (3.8 µs) 50 kHz (5 µs)
7 mm CPMAS double resonance probe (7 kHz sample rotation)	^{1}H	70 kHz (3.6 µs 90º pulse), 35 kHz, 20 kHz 50 kHz (5 µs)
7 mm CPMAS double resonance probe (7 kHz sample rotation)	^{13}C	55 kHz (4.5 µs) 50 kHz (5 µs)

Table 4.3: Reasonable RF-fields for Max. 2% Duty Cycle

Note: Higher RF power levels should only be applied if necessary and within specifications. For special probes, maximum allowed RF fields may be lower. Check with your Bruker BioSpin applications support if in doubt.

In order to have quantitative information about the precision of your magic angle, one may measure the line width of the KBr central peak and compare it with the line width of the 5th spinning sideband. If the line widths compare within ± 8% then the MAS setting is acceptable. The line width comparison is conveniently achieved with the command **peakw**, expanding the display first around the center line, typing **peakw** and then repeating this with the 5th sideband to either side.

Most CP/MAS probes are tunable over a large range of X-frequencies. It can sometimes be fairly difficult to retune a probe to an arbitrary frequency within the tuning range. NEVER load a nucleus and blindly tune and match the probe, using a small wobble width (**wbsw**) of 10 MHz or less. Instead, either note the current tuning position of the probe in the lab notebook and start retuning to the new nucleus frequency from this frequency on, following the probe response over the whole frequency range using a large **wbsw** of 50 MHz. Alternately, check the micrometer setting of the X-tuning adjustment and conclude from that to which nucleus the probe is tuned. Make a list of micrometer settings for the most frequently measured nuclei.

Remember which way to turn the tuning knob to tune to higher and lower frequencies. On most probes, turning the adjustment counter clockwise tunes to higher frequency. Do not change the matching adjustment until you have found the current tuning position of the probe, otherwise you may lose the probe response totally. Do not tune without having the appropriate matching box fitted to the preamp. Fake resonances may appear due to filters between probe and preamp, because filters are also tuned circuits. Remove all filters before tuning over a wide range, and fine tune again (**wbsw** ≤ 10 MHz) when the probe is tuned close to the desired frequency.

Changing the proton tuning will affect the X-tuning, so always tune the proton channel first, then the X-channel.

Setting a probe from high range to low range mode (lambda/4 switch) will shift the X-tuning to a lower frequency by many MHz, the proton frequency will only change by a few MHz.

An empty probe may tune as much as 10 MHz higher on the proton channel compared to a probe with a spinner in.

When a probe has not been used over an extended period, humidity may collect inside the turbine, causing a few harmless arcs (RF-breakthrough) on the proton channel. If the arcing does persist and/or gets worse, have the probe checked. Usually this means that dirt has accumulated inside the turbine or on the RF coil. Cleaning should be done by a trained person only.

Regular probe performance checks include:

- Checking the magic angle setting (KBr)
- Checking the shims (Adamantane)
- Checking S/N performance on glycine

These checks must also be performed after a probe repair. Since a repair may result in a more efficient power conversion, start with slightly reduced power settings.

SB probes flip the stator vertical for sample eject. These probes require more effort to assure a correct angle setting.

Remember to always approach the magic angle setting from the same side!

To check the reproducibility of the magic angle setting, take a KBr spectrum, stop spinning, eject and reinsert the sample, take another spectrum into a new data set, compare in dual display mode.

If the second spectrum is worse, dial less than 1/8th of a turn *counterclockwise*.

Take another spectrum, compare again.

A **laboratory notebook** should be kept with the following entries (a suitable form for printout is supplied in the Appendix):

- Name of the shim file and field value for every probe.
- Value of the power level in dB and power in watt (if available) for proton decoupling (**pl12**, **pl12W**) and associated pulse lengths **p3**, **pcpd2**.
- Value of the proton contact power level in dB and watt (**sp0**, **sp0W**).
- Value of the carbon contact power level (**pl1**, **pl1W**) and associated pulse length **p1**.
- The S/N value obtained on glycine, the SR value for shift calibration, and the line width on α-carbon in Hz.

4.9 Field Setting and Shift Calibration

Note: If spectra taken at different times and/or taken with different probes need to be compared, it is essential that the shift calibration is executed correctly. If the magnetic field was not the same, the spectra will have different values of Spectrum Reference (sr).

To reduce the requirement of readjustment and to make referencing more reliable, the field should be set to be the same for all spectra.

To keep the field the same, 2 parameters must be taken into account:

1. The drift of the magnetic field

2. The difference in shims and field between different probes.

The drift of the magnet should be measured in the following way:

Insert a sample of D2O and run a proton spectrum just like it was done on adamantane in the chapter *Calibrating 1H Pulses on Adamantane [▷ 57]*.

Without changing any parameter, rerun the spectrum every 10 minutes for a full day. This can be done with **popt**, storing the data as a 2D experiment, or use the pulse program **zg** in a 2D data set, with appropriate settings for d1 and td1 so that the drift of the magnet can be followed by the changes in peak positions of the D_2O protons. The pattern of changes in the peak position will reveal whether the changes are solely due to magnet drift, or whether there are additional disturbances to the magnet field. A magnet drift will always be constant towards lower frequency.

Note: A freshly charged magnet may also drift to higher frequency, but this will change, so it makes no sense to account for this initial drift). A nonlinear drift pattern indicates temperature changes, abrupt changes reveal magnetic disturbances (elevator, cars or trucks trains, or the keys in your pocket) only if your magnet is not shielded.

From these data, filter out the linear magnet drift. Determine the number of digits the field value in the **bsmsdisp** menu must be changed to set the field back to the exact same value as in the first spectrum. Recalculate this number to a drift time of exactly 1h/24h. Note these values as your magnet drift rate.

This drift rate will only reach a stable and constant value sometime after charging, so the drift rate measurement should be repeated until the drift value is constant. This may take some weeks to months (for high field magnets).

The magnet drift value will allow calculation of the time dependent component of the field. The probe dependent component can be established by shimming every probe and setting the field on the same sample, following the same procedure for every probe after shimming. When then the shim file is written, the current (calibrated) field position is written to disk with the shim file.

Executing the AU program (or command, TopSpin version 3.0 and later) probe field will now set the field to the appropriate value according to drift and time (taken from the date of the shim file, so the computer clock should be correct) and probe shims (from the current probe setting, so the appropriate probe must be selected in **edhead**). However: this will only work if all shim files contain the precisely determined field value for the same reference compound.

4.10 Literature

Shift referencing:

R.K Harris, E.D. Becker, S.M. Cabral de Menezes, R. Goodfellow, and P. Granger, NMR Nomenclature. Nuclear Spin Properties and conventions for Chemical shifts, Pure Appl. Chem. Vol. 73, 1795-1818 (2001)).

W.L. Earl, and D.L. VanderHart, Measurement of 13C Chemical Shifts in Solids, J. Magn. Res. 48, 35-54 (1982).

C.R. Morcombe, and K.W. Zilm, J. Magn. Reson. 162 p479-486 (2003)

IUPAC recommendation (Harris et al.):

http://sunsite.informatik.rwth-aachen.de/iupac/reports/provisional/abstract01/harris_310801.html

Cross polarization:

D. Michel, and F. Engelke, Cross-Polarization, Relaxation Times and Spin-Diffusion in Rotating Solids, NMR Basic Principles and Progress 32, 71-125 (1994).

G. Metz, X. Wu, and S.O. Smith, Ramped amplitude Cross Polarization in Magic-Angle-Spinning NMR, J. Magn. Reson. A 110, 219-227 (1994).

B.H. Meier, Cross Polarization under fast magic angle spinning: thermodynamical considerations, Chem. Phys. Lett. 188, 201-207 (1992).

K. Schmidt-Rohr, and H.W. Spiess, Multidimensional Solid-State NMR and Polymers, Academic Press (1994).

S. Hediger, B.H. Meier, R.R. Ernst, Adiabatic passage Hartmann-Hahn cross polarization in NMR under magic angle sample spinning, Chem. Phys.Lett 240, 449-456 (1995).

5 Decoupling Techniques

Line shapes in solids are often broadened by dipolar couplings between the spins. If the coupled spins are of the same kind, it is called homonuclear dipolar coupling. Heteronuclear dipolar couplings exist between nuclei of different kinds. While most dipolar couplings between X-range nuclei can be removed by magic angle spinning, couplings between ^{1}H, ^{19}F and X nuclei cannot easily and efficiently be removed by spinning. Decoupling of homonuclear and heteronuclear interactions can be obtained by different forms of RF irradiation with or without sample spinning. It is possible to suppress homonuclear couplings without suppressing heteronuclear couplings. Most frequently, the nucleus ^{1}H must be decoupled when X-nuclei like ^{13}C or ^{15}N are observed, since it is abundant and broadens the line shapes of coupled X-nuclei strongly.

5.1 Heteronuclear Decoupling

5.1.1 CW Decoupling

CW decoupling simply means irradiating the decoupled spins (usually protons) with RF of constant amplitude and phase. The decoupling program is called `cw` or `cw13` and it uses `pl12` or `pl13`, respectively. The decoupling programs select the power level and `pl12` does not need to be specified in the pulse program, if it is not used elsewhere. In the decoupling program there is also a statement setting the RF carrier frequency, according to the parameter `cnst21`, which is zero (on resonance) by default. In order to optimize decoupling, one uses the highest permitted RF field (e.g. 100 kHz for 4 mm probes) and optimizes the carrier frequency `o2` or `o2p` using **popt**.

The cw –decoupling program is written as follows:

```
0.5µ pl=pl12 ; reset power level to default decoupling power level
1 100up:0 fq=cnst21 ; reset decoupling carrier frequency to o2+cnst21
jump to 1 ; repeat until decoupler is switched off by do in the
; main ppg
```

CW decoupling suffers from the fact that protons have different chemical shifts, so irradiating at a single frequency does not decouple all protons evenly. At higher magnetic fields this becomes more evident, since the separation due to the magnetic field increases. CW decoupling requires fairly high decoupling power to be efficient.

5.1.2 TPPM Decoupling

TPPM decoupling surpasses the traditional CW decoupling. The decoupling programs tppm15 and tppm20 use a 15 and 20 degree phase shift between the two pulses, respectively. Both operate at power level `pl12`. The cpd program tppm13 uses 15 degree phase shift, as tppm15, but operates at power level `pl13`.

In order to optimize the decoupling one optimizes `pcpd2` (`AV3`, or `p31`, AV1+2) with **popt** and the carrier frequency, by varying `o2` or `o2p`. Strongly proton coupled 13C-resonances narrow substantially, especially at high magnetic fields (>300 MHz).

The figure below shows an arrayed optimization using `popt` for the TPPM phase tilt and `pcpd2` (available in TopSpin 2.0 and higher).

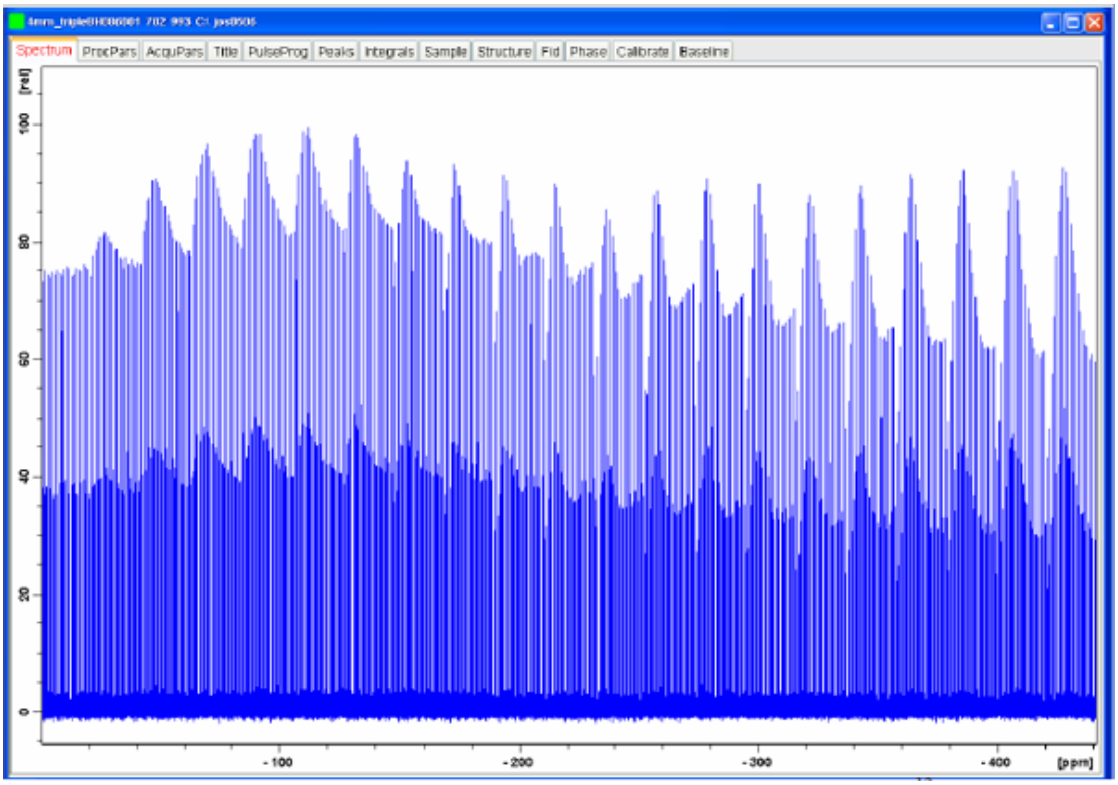


Figure 5.1: Optimization of TPPM Decoupling, on Glycine at Natural Abundance

The figure above shows optimization of TPPM decoupling, on glycine at natural abundance, ^{13}C CPMAS at 5 kHz spin rate. Each block represents a 2° degree increment of the phase toggle and the variation in each block stems from the incrementation of the pulse width in 0.2 µs increments. Optimum decoupling was found with a 4.5 µs pulse at a 16° phase toggle. It is obvious that more than one near optimum combinations of phase toggle and pulse length exist.

Reference:

A.E. Bennett, C.M. Rienstra, M. Auger, K.V. Lakshmi, and R.G. Griffin; Heteronuclear decoupling in rotating solids, J. Chem. Phys. 103 (16); 6951 – 6958 (1995).

5.1.3 SPINAL Decoupling

SPINAL provides adequate decoupling bandwidth even for high field (>400 MHz) instruments at an RF level of 80 kHz or higher. SPINAL 64 (64 phase permutations) outperforms TPPM and may be used as standard decoupling sequence. SPINAL 64 can be optimized in the same way as TPPM, by incrementing `pcpd2` (p31) (the phase shifts are fixed). The decoupling pulse is an approximate 180° pulse.

Reference:

B.M. Fung, A.K. Khitrin, K. Ermolaev, J. Magn. Reson. 142, 97-101 (2000).

5.1.4 Swept-Frequency-TPPM

This decoupling method combines TPPM and a frequency variation via the pulse length variation to achieve a wider decoupling bandwidth. The decoupling efficiency is better than TPPM (especially at high fields), and comparable to if not better than SPINAL-64. The corresponding cpd-program is called `swftppm`.

Reference:

R.S. Thakur, N. D. Kurur, and P. K. Madhu, J. Magn. Res. 193, 77 (2008).

5.1.5 XiX Decoupling

XiX decoupling requires high spinning speeds, but decouples at a moderate RF level. 180° proton pulses are used, synchronized to the rotor speed such that recoupling does not occur (pcpd2≠n/4*rotor periods). Usually, `pcpd2` is selected to be about 1/3 rotor period. The decoupler power level must be adjusted to produce a 180° pulse of (rotor period)/3.

Reference:

A. Detken, E. H. Hardy, M. Ernst, and B. H. Meier, Chem. Phys. Lett. 356, 298-304 (2002).

5.1.6 Pi-Pulse Decoupling

Pi−pulse decoupling is a decoupling program, for weaker nuclear interactions like J couplings or weak dipolar interactions, using rotor synchronized 180° pulses. π-pulse decoupling uses the xy-16 phase cycle for large bandwidth. Abundant protons cannot be sufficiently decoupled with this method, but it is very suitable for removing couplings to ^{31}P, which is hard to do with **cw** or **tppm**, since the chemical shift range is wide. Likewise, it can be used to decouple dilute spins or spins which are homonuclear decoupled by spinning (^{19}F).

Reference:

S.-F. Liu and K. Schmidt-Rohr, Macromolecules 34, 8416-8418 (2001).

5.2 Homonuclear Decoupling

Homonuclear decoupling refers to methods which decouple dipolar interactions between like spins. Those are only prominent between abundant spins like ^{1}H, ^{19}F and ^{31}P (and potentially some others). This interaction cannot easily be spun out in most cases and renders NMR-parameters like chemical shifts of the homonuclear coupled spins or heteronuclear couplings and J-couplings to other (X-)nuclei unobservable.

5.2.1 Multiple Pulse NMR: Observing Chemical Shifts of Homonuclear Coupled Nuclei

Multiple pulse NMR methods are covered in the chapters about CRAMPS of this manual collection. The principle of those methods (CRAMPS, if MAS is used to average CSA interactions simultaneously), is to set the magnetization of the spins into the magic angle, using a suitable pulse sequence. In this case, the dipolar couplings between those spins are suppressed. Short observation windows between pulses allow observation of the signal from the decoupled nuclei.

Reference:

S. Hafner and H.W. Spiess, Multiple-Pulse Line Narrowing under Fast Magic-Angle Spinning, J. Magn. Reson. A 121, 160-166 (1996) and references therein.

5.2.2 Multiple Pulse Decoupling

Multiple Pulse Decoupling: Observing dipolar couplings and j-couplings to homonuclear coupled nuclei.

Homonuclear couplings between abundant spins (usually protons) superimpose their heteronuclear dipolar couplings to X-spins and J-couplings to X-spins so these (distinct) couplings are not observable. homonuclear decoupling protons while observing X-spins makes these couplings observable. Any method used in multiple pulse NMR may be used to achieve this.

5.2.2.1 BR-24, MREV-8, BLEW-12

Used as heteronuclear decoupling methods, the window between pulses may be shortened or omitted (semi-windowless or windowless sequences). These sequences work well, but have rather long cycle times and are therefore not suitable for fast spinning samples. Or else they work in a similar fashion as the sequences covered in the following. BLEW-12 decoupling is supplied as a standard `cpd` program. It consists of a windowless sequence of 90° pulses with suitable phases. High RF levels for decoupling provide better resolution.

5.2.2.2 FSLG Decoupling

The Frequency Switched Lee Goldburg (FSLG) sequence may be used at spin rates up to 15 kHz. It is a homonuclear decoupling sequence which rotates the interaction Hamiltonian around an effective field, aligned at the magic angle (arctan $\sqrt{2}$) with respect to the Zeeman field in the rotating frame. The tilt is achieved by off resonance irradiation at the Lee Goldburg frequency fLG according to the Lee Goldburg condition.

The off-resonance condition depends on the RF field of irradiation, but not very sensitively however. The irradiation frequency jumps between +/- (RF field)/sqrt(2) for the duration of two 360°pulses on resonance (=293° pulses at the LG-frequency) with a 0/180° phase alternation. The include file <lgcalc.incl> calculates all values according to the RF field (set in Hz as cnst20) within the pulse program.

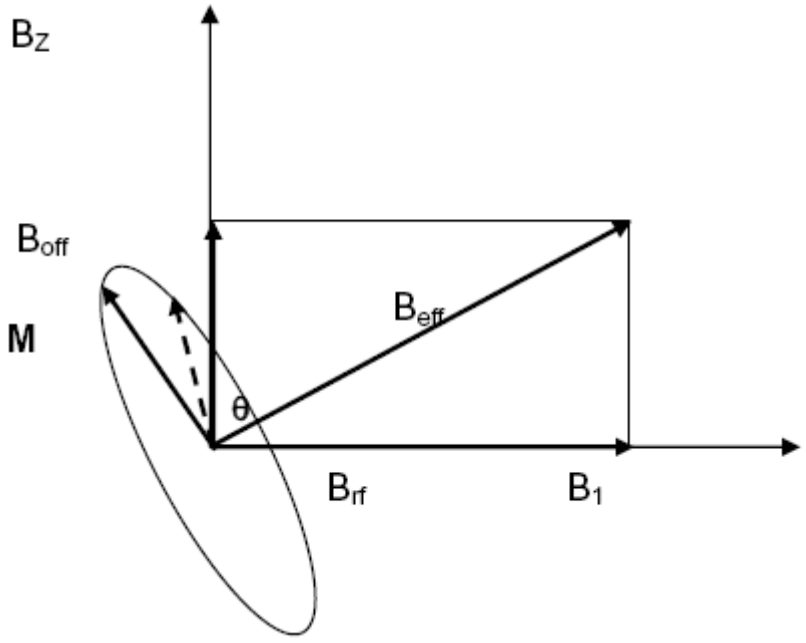


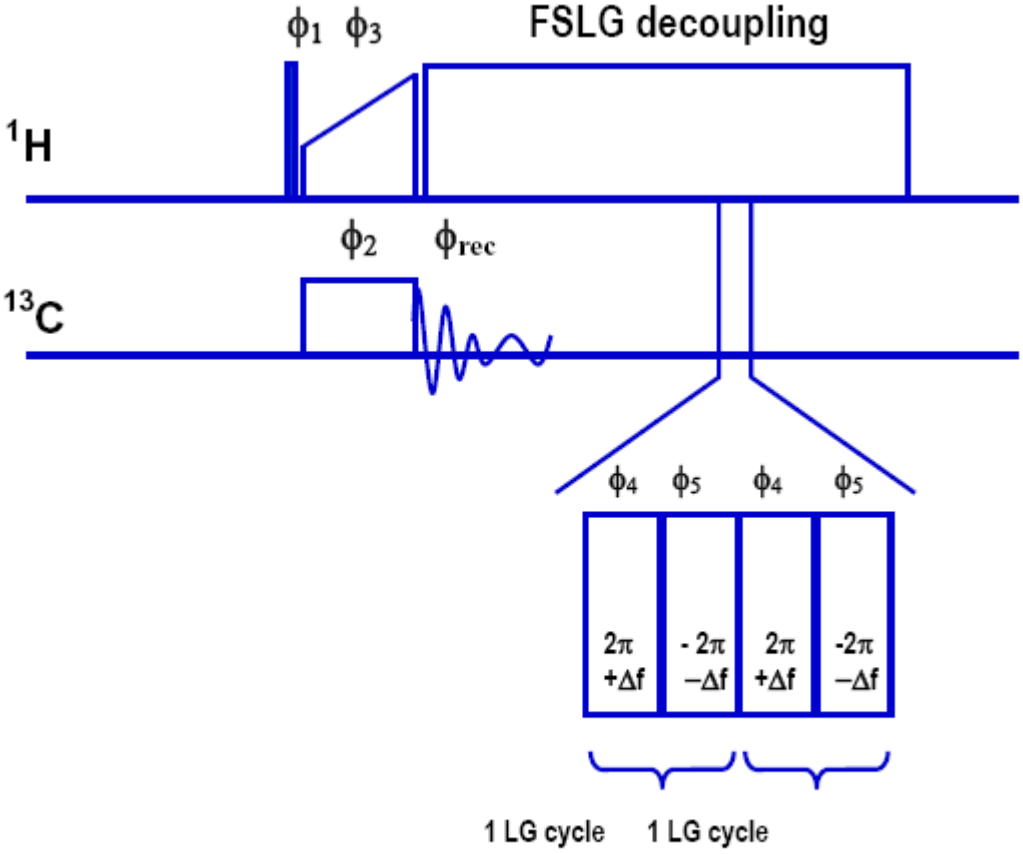
Figure 5.2: Geometry for the FSLG Condition

Note that Beff points along the 1 1 1 direction in the 3 dimensional space. Note the sign of Boff when calculating the actual direction of the effective field. A positive Boff and a B1 with phase 0 results in the effective field being in the positive quadrant along the magic angle in the X-Z plane of the rotating frame.

Two methods are available to achieve such a frequency switch experimentally. One method is simultaneous switching of frequencies and phases. The other method uses phase-modulation. Frequency, time and phase relate to each other as derivative of phase and time

as to get $^{2\pi}$ = .../. The relationship describes the rate at which a phase of the RF pulse must be changed in order to achieve a certain frequency offset. Vinogradov et al. describe this approach under the acronym PMLG (Phase Modulated Lee Goldburg).

Used in combination with `cp` signal generation, both methods allow observing proton-J-couplings to the observed X-nucleus. However, only samples with very narrow lines will produce well resolved J-couplings as shown below on adamantane. Harder solids require careful adjustment and fairly high power levels to show barely resolved couplings, since the line widths achieved are broader than what can be achieved with standard decoupling sequences like `tppm`. Since the heteronuclear X-H-coupling remains, there may be spinning sidebands from this coupling, in addition to CSA sidebands.



ϕ_1=1 3	ϕ_2 = 0 0 2 2 1 1 3 3
ϕ_3=0	ϕ_{rec}= 0 2 2 0 1 3 3 1
ϕ_4= 0	ϕ_5= 2

Figure 5.3: FSLG Decoupling Pulse Sequence Diagram

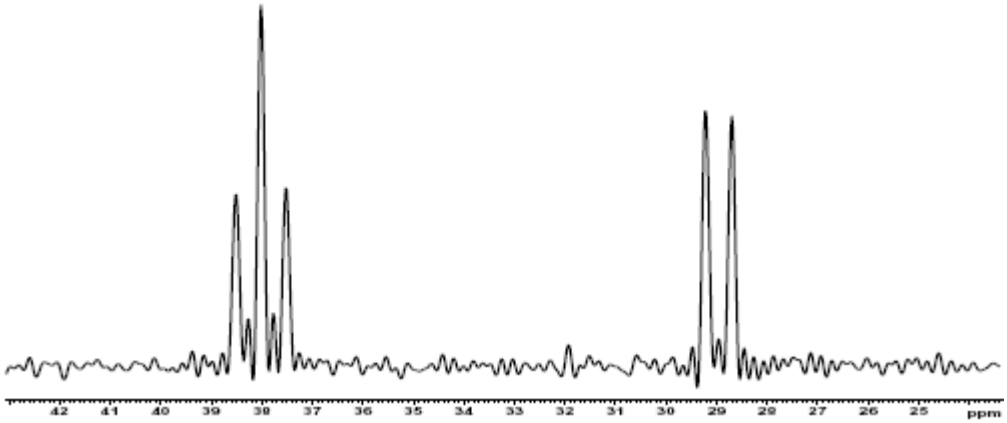


Figure 5.4: Adamantane, FSLG-decoupled, showing the (downscaled) C-H J-couplings

The figure above shows homonuclear proton decoupling on center packed adamantane sample rotating at 7 kHz, 100 kHz 1H decoupling field. Note that good B1-homogeneity is required. Use a CRAMPS spinner (12 µl sample volume in a 4mm spinner) for this experiment.

Setting up the experiment:

1. Use center packed adamantane in a CRAMPS rotor, unlabeled and a spinning rate of 10 kHz for adamantane.
2. Start from a data set with a well adjusted HH condition on adamantane.
3. Generate a new data set with **edc**.
4. Readjust the decoupling power pl12 and p3, pcpd2 for 70-100 kHz RF field, determine the precise RF field (preferably via a 360° proton pulse p3).
5. Load the pulse program `fqlg`. This uses frequency shifts with simultaneous phase shifts for FSLG-decoupling at pl13.
6. Set pl13 to achieve the same RF field as measured in step 4, set cnst20 to the value of RF field in Hz. The pulse program contains the include file <lgcalc.incl> which calculates the required frequency shifts to either side (shown in `ased` as `cnst22` and `cnst23`). Cnst24 provides an additional overall offset to compensate for phase glitch. With proper probe tuning and 50 Ohms, `cnst24` should be close to zero.
7. Set acquisition and processing parameters.

A spectrum like that in the previous figure should be obtained. If the splitting is worse, optimize with `pl13` and `cnst24`. Usually, somewhat less power than calculated is required.

The FSLG decoupling scheme is also implemented as cpd-program `cwlgs`. The include file lgcalc.incl is also required. With `cpdprg2` = cwlgs, the standard cp pulse program can be used. The ZGOPTN –Dlacq (ased) should be set in order to allow decoupling times >50 ms.

Parameter	Value	Comments
pulprog	fqlg	AV3, use fqlg.av for AV1+2.
d1	4 s	Recycle delay.
ns	4-16	Number of scans.
aq	80 ms	Acquisition time.
spnam0	ramp.100 or, ramp70100.100	For ramped CP.
pl12, p3	set for p3=90°	
sp0, pl1	set for cp	
p15	5-10m	
pl13	set for 70-100 kHz	Optimize for best resolution.
cnst20	70000-100000	Equals the applied RF field.
cnst24	0	To be optimized.
cnst21	0	Reset proton frequency to SFO2.

Table 5.1: Acquisition Parameters

Parameter	Value	Comment
SI	2*td	Adequate 4fold zero filling.
WDW	no	No apodization.
PH_mod	pk	Phase correction if needed.
BC_mod	quad	DC offset correction.

Table 5.2: Processing Parameters

As mentioned above, frequency shifts can also be generated by a phase gradient shape. A phase change of 360° per second corresponds to a frequency of 1 Hz, as can easily be visualized. The frequency shift which needs to be achieved is ±RF field/$\sqrt{2}$. Since the pulse duration must achieve a 2π rotation off resonance, corresponding to a 293° flip angle on resonance, it can easily be calculated that a phase change over 209° during a 293° flip angle pulse is required to achieve this.

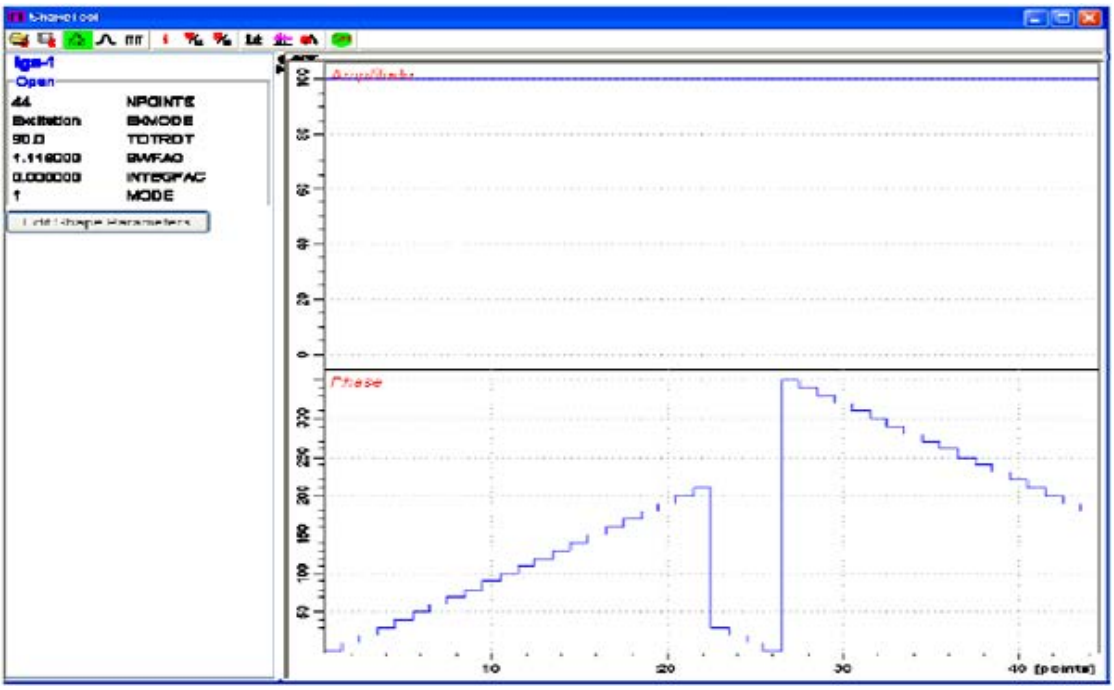


Figure 5.5: Shape with Phase Gradients

In the figure above: Shape with phase gradients for positive and negative offsets and corresponding phase change, stdisp –display of lgs-1 shape. Amplitude is 100% throughout.

Vinogradov et. al. have published shapes with much fewer steps and different phases. The pulse length for the shape does not depend on the number of steps, but only on the applied RF field. Using the include file lgcalc.incl, a pulse p5 is calculated from cnst20=RF field in Hz. The total shape pulse length must be 2*p5.

To use pmlg decoupling, save the pulse program fqlg under a different file name and change calculations and loop as follows:

```
define loop counter count ; calculate number of LG periods according
to aq
"count=aq/(2*p5)"
define pulse pmlg
"pmlg=2*p5"
"sp1=pl13" set shape power to pl13 for LG (TS2.1 only)
.
3 (pmlg:sp1 ph3): f2 for one full PMLG unit, as for lgs-1 shape
lo to 3 times count
```

References:

A. Bielecki, A.C. Kolbert, and M.H. Levitt, Frequency-Switched Pulse Sequences: Homonuclear Decoupling and Dilute Spin NMR in Solids, Chem. Phys. Lett. 155, 341-346 (1989).

A. Bielecki, A.C. Kolbert, H.J.M. deGroot, R.G. Griffin, and M.H. Levitt, Frequency-Switched Lee-Goldburg Sequences in Solids, Advances in Magnetic Resonance 14, 111-124 (1990).

E. Vinogradov, P.K. Madhu, and S. Vega, High-resolution proton solid-state NMR spectroscopy by phase modulated Lee-Goldburg experiments, Chem. Phys. Lett. 314, 443-450 (1999) and references cited therein.

5.2.2.3 DUMBO

DUMBO (Decoupling Uses Mind Boggling Optimization) is a phase modulation scheme where the phase modulation is described in terms of a Fourier series

$$\phi(t) = \sum_{n=0}^{+\infty} a_n \cos(n\omega_c t) + b_n \sin(n\omega_c t)$$

The shape can be created using the AU-program **DUMBO**. The DUMBO shape file in the release version of TOPSPIN is calculated for 32 µs pulses. To create your own DUMBO shape you can also use the au-program dumbo. See instructions in the header of the au-program for proper use.

The above pulse program to observe J-couplings with DUMBO decoupling would be written as follows:

```
define loop counter count ;calculate number of LG periods according
to aq
"count=aq/(p10)"
.
3 (p10:sp1 ph3):f2 ;p10 set by AU-program DUMBO (n*32 µsec)
lo to 3 times count
```

References:

D. Sakellariou, A. Lesage, P. Hodgkinson, and L. Emsley, Homonuclear dipolar decoupling in solid-state NMR using continuous phase modulation, Chem. Phys. Lett. 319, 253-260 (2000).

Lyndon Emsley's home page: *http://www.ens-lyon.fr/STIM/NMR/NMR.html*

5.3 Transverse Dephasing Optimized Spectroscopy

Decoupling optimized under refocused conditions:

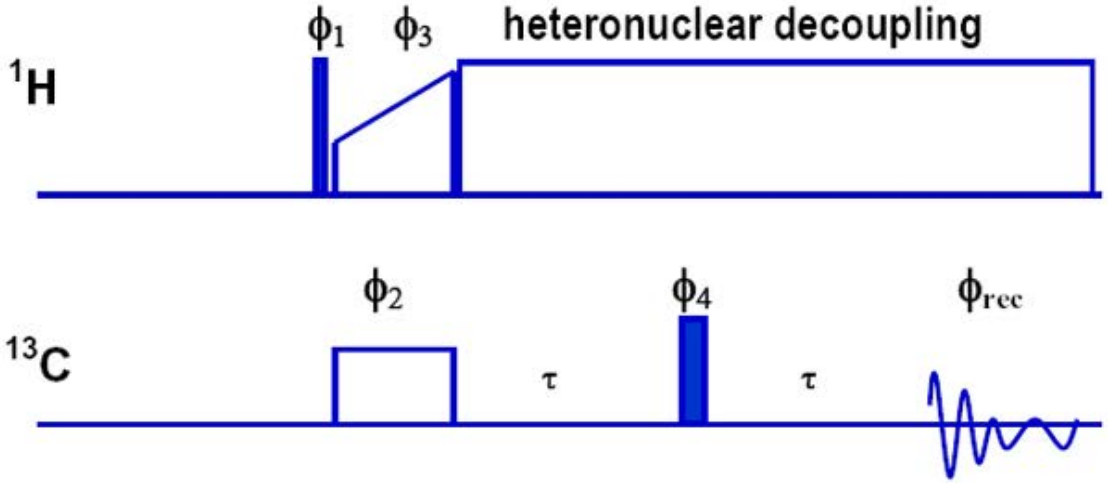


Figure 5.6: Pulse Program for Hahn Echo Sequence

Transverse Dephasing Optimized spectroscopy (G. De Paepe et al. 2003) uses a spin echo sequence for optimizing heteronuclear decoupling. The idea behind it is to simply remove of normally dominant J0 term (describing coherent residual line broadening effects) in the transverse relaxation rate R2 (A. Abragam chapter 8). With the normal CP experiment the observed line broadening (coherence decay time T*2) might be caused by other heterogeneous effects, such as distribution of chemical shifts or susceptibility effects and not reflect the true T'2 (coherence lifetime). The true T'2 achieved through good heteronuclear decoupling can then be observed with a hahn-echo experiment. Optimization is done by looking for the maximum signal amplitude of the decoupled resonances of interest. Be careful not to exceed the maximum decoupling time with high power decoupling.

Reference:

G. De Paepe, N. Giraud, A. Lesage, P. Hodgkinson, A. Böckmann, and L. Emsley, Transverse Dephasing Optimized Solid-State NMR Spectroscopy, JACS 125, 13938 – 13939 (2003).

6 Practical CP/MAS Spectroscopy on Spin 1/2 Nuclei

Once good setup parameters have been obtained to observe ^{13}C and get good S/N on glycine, it should be easy to also observe ^{13}C-CP/MAS spectra on other samples and on nuclei different from ^{13}C. Nevertheless, sometimes one comes across samples where it is difficult to observe ^{13}C. This chapter deals with strategies to optimize acquisition parameters for ^{13}C and other spin ½ nuclei.

6.1 Possible Difficulties

Usually, ^{13}C spectra are easily acquired. Several sample properties may however make observation difficult:

- Low concentration of ^{13}C in the sample.
- No or too few protons in the sample.
- Long proton T_1.
- Long $T_{I\text{-}S}$.
- Short proton $T_{1\rho}$.

If a nucleus different from ^{13}C should be observed, there are additional potential difficulties:

- Unknown chemical shift.
- Unknown Hartmann-Hahn-condition.
- Unknown relaxation properties (proton T_1, $T_{1\rho}$, T_I).

6.2 Possible Approaches for 13C Samples

- Collect as much information about the sample as possible. Do not accept samples for measurement with unknown composition. Request information about:
  - Possible hazards (upon a rotor explosion).
  - Concentration of the nucleus to be measured.
  - Structural information about the molecular environment of the nucleus of interest:
    - mobility (rigid environment: expect long T_1 and repetition delay),
    - proximity to protons (can one use cross polarization),
    - conductivity, dielectric loss (expect tuning and RF heating problems if sample is dielectrically lossy or even conductive).
- Collect information about the sample first by running an “easy” nucleus
  - Feasibility of cross polarization parameters is the required key information, because it decides the steps to follow.
  - If the sample information which you have collected shows that a ^{13}C CP/MAS experiment should be feasible (sample contains more than 20% protonated carbons), load a reference cross polarization data set (S/N test spectrum of glycine), spin the sample at the same spin rate, set contact time (**p15**) to 1ms, wait 1 min., do one scan. There should be a visible signal.
  - From there on, optimize the required repetition rate (**d1**), contact time (**p15**), number of scans (**ns**), spin rate (**masr**) and Hartmann-Hahn adjustment until the signal is optimum. In very few cases, the decoupler offset (**o2**) may require readjustment.

  - If no ^{13}C-signal is found, the reasons may be:
    - Incorrect setup (recheck reference sample).
    - Concentration lower than expected.
    - Unusual relaxation properties (long $T_{I\text{-}S}$, long proton T_1, short proton $T_{1\rho}$).
- Then the most important information about the sample (proton T1, proton T1ρ) can be obtained by looking at the protons in the sample. Set up for proton observation, set **swh** to 100000-500000, **rg** to 4 and **pulprog cpopt** (if not found in the library, copy the pulse program in the appendix), **p3** and **pl12** for p3=p90. Set **spnam0** = ramp.100, **sp0** = power level for HH, **p15** =100 µs. Do 1 scan and fourier transform/phase correct. Using popt, optimize **d1** for maximum signal.

  Note: CP/MAS probes usually have a substantial proton background signal. Do not be misled by this, it will not behave like a regular signal:
  - It will grow steadily with longer pulses
  - It will not show spinning sidebands
  - It will cancel when a background suppression pulse program like a ring is used with a full phase cycle.
- Knowing the required relaxation delay, the following step is to determine the cross polarization (contact time). On protons, we measure the time constant $T_{1\rho}$. Using popt in the previous setup, vary **p15** between 100 µsec and 10 ms (even 20 ms at reduced power, if a long $T_{I\text{-}S}$ is expected, as the distance between nucleus of interest is long or the mobility is high, leading to a small heteronuclear dipolar coupling between nucleus of interest and protons). This measurement will tell you how long the contact time p15 may be. A value of **p15** giving 50% of the initial proton signal amplitude will still give a 2-fold enhancement on ^{13}C. If the proton signal is below 50% at 1 ms spin lock time or even less, a full cp-enhancement cannot be expected.
- The minimum relaxation delay and the maximum contact time are now known. With these parameters used as d1 and p15, the measurement is just a matter of patience.

6.3 Possible Approaches for Non-13C Samples

If an arbitrary X-nucleus of spin ½ is under investigation (quadrupolar spins must be treated separately), the strategy follows the one described above, if the sample contains the protons bound to ^{13}C. In this case, running a ^{13}C CP/MAS spectrum allows setting and determining all proton parameters (recycle time, contact time) from the ^{13}C setup. To run the X-nucleus, cross polarized from protons, set the HH-condition from the known proton RF field, the spin rate, and the transmitter power at the NMR frequency of the X-nucleus such that the effective field at the X-frequency equals the effective field at proton frequency ±spin rate.

Example: Setting the HH-condition for ^{15}N from known parameters for ^{13}C-CP/MAS. The gyromagnetic ration of ^{15}N is lower by a factor of 2.5 compared to carbon (proton frequency: 400 MHz, ^{13}C-frequency: 100 MHz, ^{15N} frequency: 40 MHz). The probe efficiency is about the same for ^{13}C and ^{15}N (but not ^{1}H!), so one needs about 2.5 times higher RF voltage for the ^{15}N-contact pulse than for the ^{13}C-contact pulse, if the spin rate and the proton RF field are the same. This is equivalent to 2.5^2=6.25 times the power in watts. So if `ased` shows **pl1W** =150W for a well optimized ^{13}C-CP setup, ^{15}N will require 6.25*150 W= 938 W. This is far above specifications, so the same proton contact power level cannot be used, it needs to be lowered. The maximum allowed power for a contact pulse on ^{15}N is 500W. This means that the proton contact power should be lowered by approximately a factor of sqrt (938/500) ≈1.37. Precalculating power levels like this will get the parameters close enough to see a cp-signal on a good test sample, so further optimization is possible. See *Test Samples [▸ 11]* for suitable test samples.

The most efficient way of precalculating power levels for multi-nuclear spectroscopy is the following:

- Determine the power conversion factor for some nuclei of interest on a suitable test sample, from the low end to the high end of the probe tuning range. This means measuring a precise 360° pulse (make sure it is 360°, not 180° or 540°) and the associated power level. Make a table in your lab notebook as follows:

Nucleus Frequency	P90 (µs)	RF Field (KHz	Power (W or dB)	Remarks
^{1}H400.13	2.5	100	100	Low range
^{19}F/376.3				Not available
^{15}N/40.5	6.5	38.6	300	Probe in double mode
^{15}N/40.5	6.5	38.6	500	Probe in triple mode C/N
^{29}Si/79.5	6	41.7	300	Double mode low range
^{13}C/100.5	4	62.5	150	Double mode low range
^{13}C/100.5	5	50	200	Triple mode C/N
^{119}Sn/149.1	4	62.5	100	Double mode high range
^{31}P/161.9	3.5	71.4	150	Range switch up, double mode

Table 6.1: Power Conversion Table – 4 mm Triple Probe

- Once these values are measured, any HH condition can be calculated. Assuming you want to cross polarize ^{119}Sn, the sample spins at 12 kHz. The contact time is anticipated to be rather long, because ^{119}Sn atoms are large and far away from protons. So the power level for the contact should not be too high. Set the RF field to 50 kHz for the contact. When, for example, we decide to apply a ramp shape on the proton contact pulse, covering the ±1 spinning sidebands. This means that we need to apply a ramp from 38 to 62 kHz RF field, plus some safety margin, about 35 to 65 kHz RF field on the proton ramp. For ^{119}Sn we need to apply 50 kHz RF field. Since the RF field is proportional to the amplitude in a shape (RF voltage output is proportional to shape amplitude value), the shape power must range from 65 kHz to 35 kHz, from 100 to about 50% amplitude. Use **calcpowlev** to calculate the changes in dB to achieve the calculated RF fields (enter reference RF field to calculate required RF field instead of pulse lengths). In our case, the proton contact pulse power sp0 is calculated at + 3.74 dB (65 kHz compared to 100 kHz), the power level for ^{119}Sn is calculated at +1.94 dB (50 kHz compared to 62.5 kHz). Be sure to add the calculated number for a desired RF field lower than the reference field, subtract the number if the desired RF field is higher.
- If such a table is not available, but an oscilloscope is, one can measure the RF voltage for the X contact pulse of the known (^{13}C) HH condition, calculate the pp-voltage for the unknown HH condition from the NMR-frequencies of the two nuclei, and set this voltage for the unknown HH condition.

6.4 Hints, Tricks, Caveats for Multi-nuclear (CP-)MAS Spectroscopy

- Since T_1 relaxation tends to be slow in solids, direct observation of hetero-nuclei is usually time consuming, so CP is widely used because the proton T_1 is usually bearable. However, CP can only be used if the hetero-nucleus is coupled to protons (or whatever nucleus the magnetization is drained from). Whereas ^{13}C and ^{15}N usually bear directly bonded protons, this is not the case for many other spin ½ hetero-nuclei. So the magnetization must come from more remote substituents. Moreover they may also result because atomic radii increase as one goes to nuclei with higher atomic mass. In short: HH conditions may be very sharp, $T_{I\text{-}S}$ may be long, but proton $T_{1\rho}$ may still be short.
- Chemical shift ranges and chemical shift anisotropies increase with nuclei of higher order number and number of electrons in the outer shell. Therefore one may be confronted with two problems:
  - To find the signal somewhere within the possible chemical shift range.
  - To find the signal within a forest of spinning sidebands.
- Ease of setup therefore depends largely on the availability of a setup sample with decent T_1, efficient CP, and known chemical shift for referencing. The chapter *Test Samples [▷ 11]* lists some useful setup samples together with known parameters.

6.5 Setup for Standard Heteronuclear Samples 15N, 29SI, 31P

^{15}N on α-glycine:

- Calculate HH condition as described above.

Or else:

- Load α-glycine ^{13}C reference spectrum.
- Set observe nucleus N15 in **edasp**.
- Add 2 dB to **sp0 (spnam0=ramp.100)** .
- Subtract 2 dB from **pl1** (more is not required since the transmitter will usually put out 50% more power at ^{15}N frequency) .
- Set **p15** = 3 ms.
- Acquire 4-8 scans.
- Optimize HH condition, acquire reference spectrum with **aq**=25-35 ms.

^{29}Si on DSS

- Load α-glycine ^{13}C reference spectrum.
- Set observe nucleus to ^{29}Si in **edasp**.
- Add 2 dB to **sp0**.
- Acquire 4 scans with **aq** = 35 ms.
- Optimize HH condition, acquire reference spectrum.

^{31}P on ADP (ammonium dihydrogen phosphate $NH_4\ H_2\ PO_4$)

- Load α-glycine ^{13}C reference spectrum.
- Set observe nucleus to ^{31}P in **edasp**.
- Add 6 dB to **pl1**.
- Optimize HH condition, acquire 2 scans, reduce **rg** appropriately.

7 Basic CP-MAS Experiments

The following experiments can be run by calling a ^{13}C CP-MAS standard parameter, data set, or data, loading the appropriate pulse program and loading the pulse parameters obtained previously during the setup (see *Basic Setup Procedures [49]*). Some attention needs to be paid to special experimental parameters. Most of the parameters are explained in the header section of the pulse programs.

The CPPI experiment series in section *Spectral Editing Sequences: CPPI, CPPISPI and CPPIRCP [101]* requires measuring the HH match using a constant amplitude contact pulse. This can be accomplished using a rectangular shape **square.100**, or using the pulse program `cplg`.

7.1 Pulse Calibration with CP

The following figure is an example of a pulse calibration for ^{13}C pulses after cross polarization using a flip back pulse.

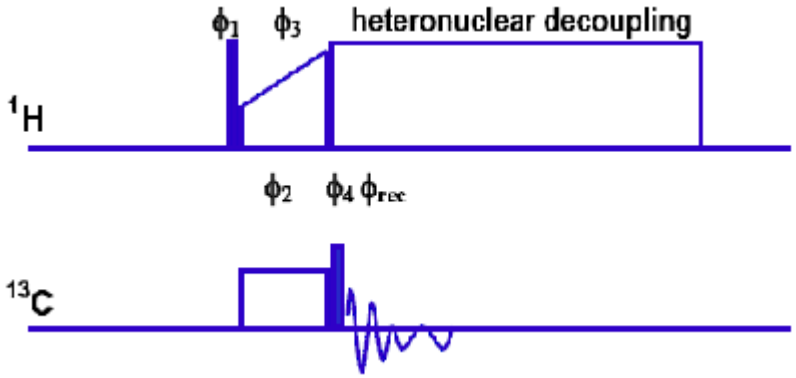


Figure 7.1: Pulse Program for CP with Flip-back Pulse

The experiment can be done directly after the CP-MAS setup procedure. Loading the pulse program `cp90` and setting `pl1=pl11` allows one to measure the X nucleus spin nutation frequency at the HH contact power. Of course, the experiment allows nutation frequencies to be measured at other power levels as well. The typical nutation pattern has a cosine form, so a 90 degree pulse gives null signal. Use glycine spinning at N kHz as before. When using POPT for such measurements the optimization type is “ZERO” so that the program looks for a zero crossing at the automatic data evaluation. To get nutation patterns without phase distortions, 90° pulses should always be executed close to the observed resonance. Larger offsets give different (shorter) p90 values and phase distortions for pulse lengths close to 180° and multiples thereof.

Parameter	Value	Comments
pulprog	cp90	AVIII, cp90.av for older instruments
nuc1	13C	Nucleus on f1 channel
nuc2	1H	Nucleus on f2 channel
sw	300 ppm	Spectral width for Glycine
o1p	45	Close to C-α
td	2048	Number of points sampled

Table 7.1: Acquisition Parameters

Fine adjustment of the π pulse on ^{13}C can also be done using the TOSS experiment, this will be discussed in the next section.

7.2 Total Sideband Suppression TOSS

The TOSS sequences permit complete suppression of spinning sidebands (SSB) in CP-MAS experiments. The TOSS sequence consists of the basic CP sequence plus a 2 rotor period sequence with four specially spaced 180º pulses. As is the case for all extra pulses on the X channel in CP-MAS experiments (with the exception of symmetry based sequences, see further below), these 180º pulses are set with `pl11`.

This experiment can be optimized for minimum spinning sideband intensity either by variation of the 180° pulse width or the associated power level **pl11**.

Two variations of the TOSS sequence exists, the default is TOSS A, which is appropriate for lower spinning speeds. TOSS B, for higher spinning speeds, is selected by setting ZGOPTNS to –Dtossb. The maximum spinning speed is either determined by common sense – if all sidebands are spun out, TOSS is not needed (low field instruments) – or by the shortest delay, which is d26 in both cases. For TOSS B, d26 = 0.0773s/cnst31-p2, with cnst31 the rotation rate in Hz and p2 the 180º pulse width in µs. For TOSS A, d26 = 0.0412s/cnst31-p2, so the maximum spinning rate is lower.

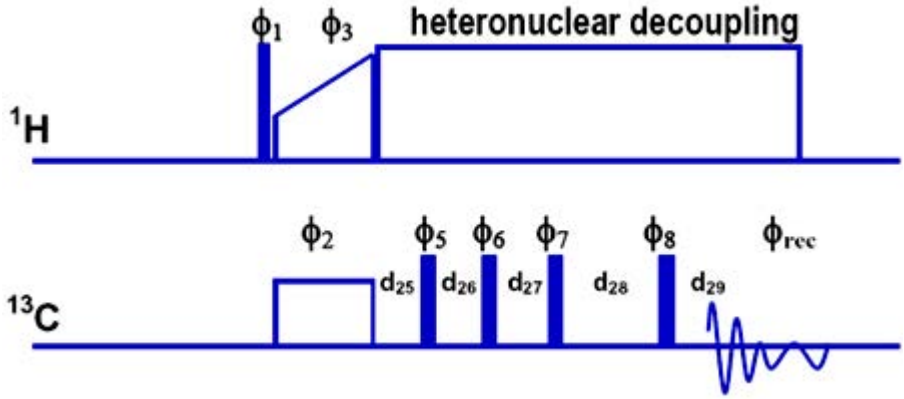


Figure 7.2: Pulse Program for CPTOSS

If the timing becomes a problem, alternative TOSS schemes need to be [2]. The SELTICS sequence is an alternative.

[2] foundO.N. Antzutkin, Sideband manipulation in magic-angle-spinning nuclear magnetic resonance; Progress in Nuclear Magnetic Resonance Spectroscopy 35 (1999) 203-266.

Set up the experiment using glycine or tyrosine-HCl at a moderate spinning speed. Get a good CP-MAS spectrum first then run a TOSS spectrum.

Parameter	Value	Comments
pulprog	cptoss cptoss243	
p2		180° pulse on X nucleus
pl11		Power level driving P2 on X-channel
cnst31		Spinning speed in Hz e.g. 5 kHz the entry would be 5000
zgoptns	-Dtossb	Tossb if needed because of high spinning speed or long p2

Table 7.2: Acquisition Parameters

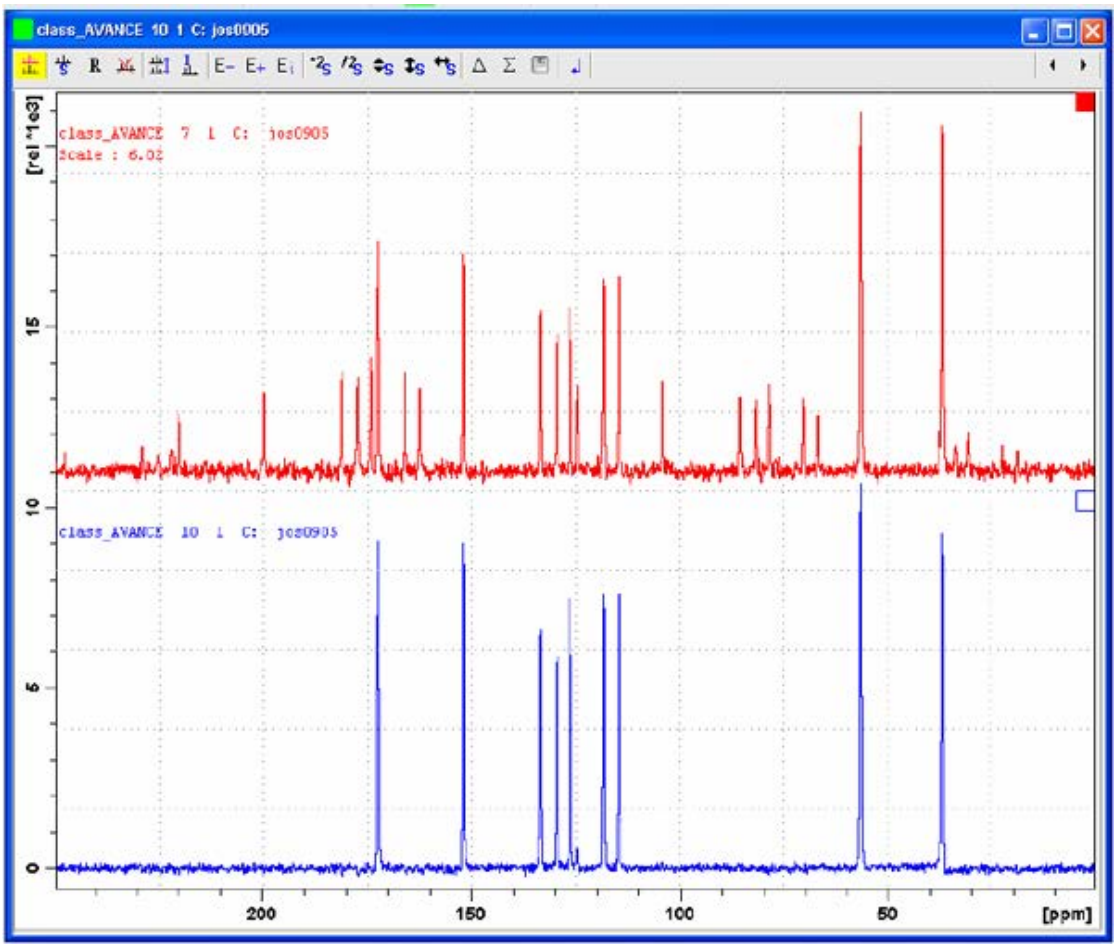


Figure 7.3: Comparison of a CPTOSS and CP-MAS Experiment

The figure above compares a CPTOSS experiment (lower spectrum) to a CP-MAS experiment (upper spectrum) on tyrosine HCl at 6 kHz sample rotation using a 4 mm CP-MAS double resonance probe at 500 MHz with 16 accumulated transients.

The sequence is not perfectly compensated for experimental artifacts and if perfect suppression of SSB is required, one can use a 5 pulse sequence with a long phase cycle, requiring a minimum of 243 transients for complete artifact suppression using the pulse program `cptoss243`, where the extension.av is added in case of the AV2 console. The following figure shows the advantage of the well compensated TOSS sequence with its 243 phase cycle steps over the above 4 pulse sequence. Besides better compensation, the `cptoss243` pulse sequence is also shorter and uses only 1, instead of 2 rotor cycles. This pulse program can be used with fairly high spinning speeds, up to about 12.5 kHz sample rotation, depending, of course, on the width of the employed π – pulses. The last figure in this section shows a comparison of the results obtained with the 4 pulse sequence with 256 scans.

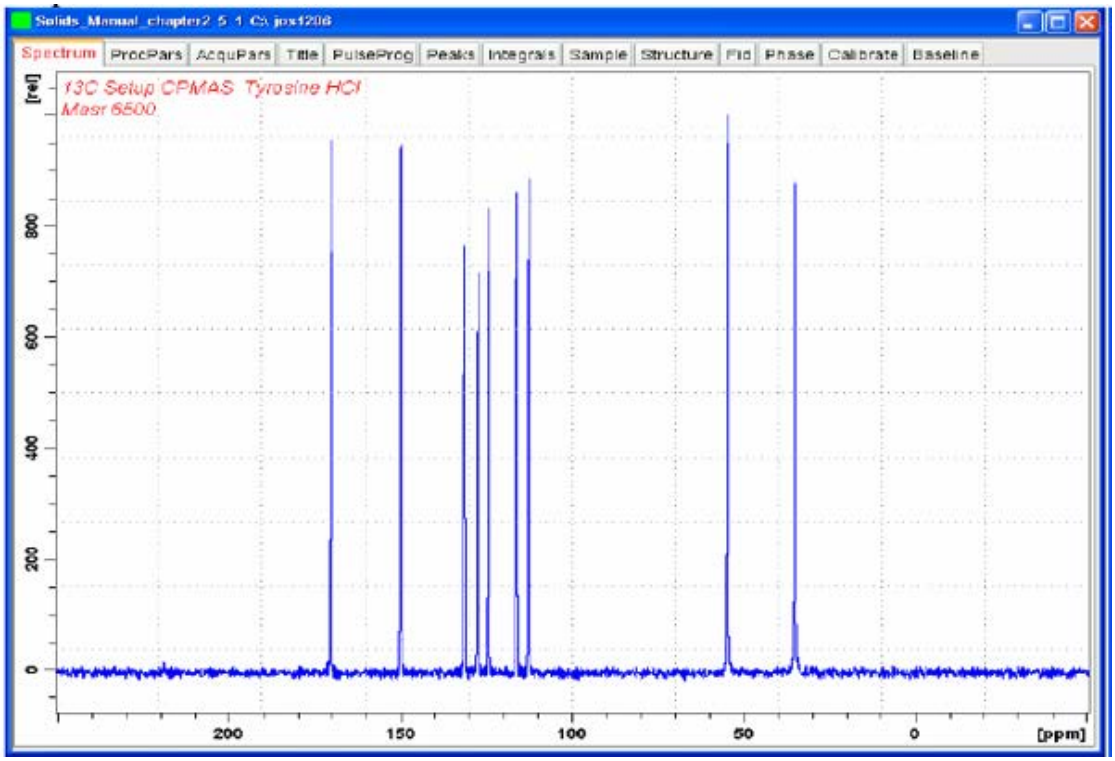


Figure 7.4: CPTOSS243 Experiment on Tyrosine HCl at 6.5 kHz

The figure above is a CPTOSS243 experiment on tyrosine HCl at 6.5 kHz sample rotation using a 4 mm CP-MAS triple resonance probe at 500 MHz with 243 accumulated transients. No spinning sideband residuals can be observed, with a noise level below 2% peak to peak compared to the highest peak intensity.

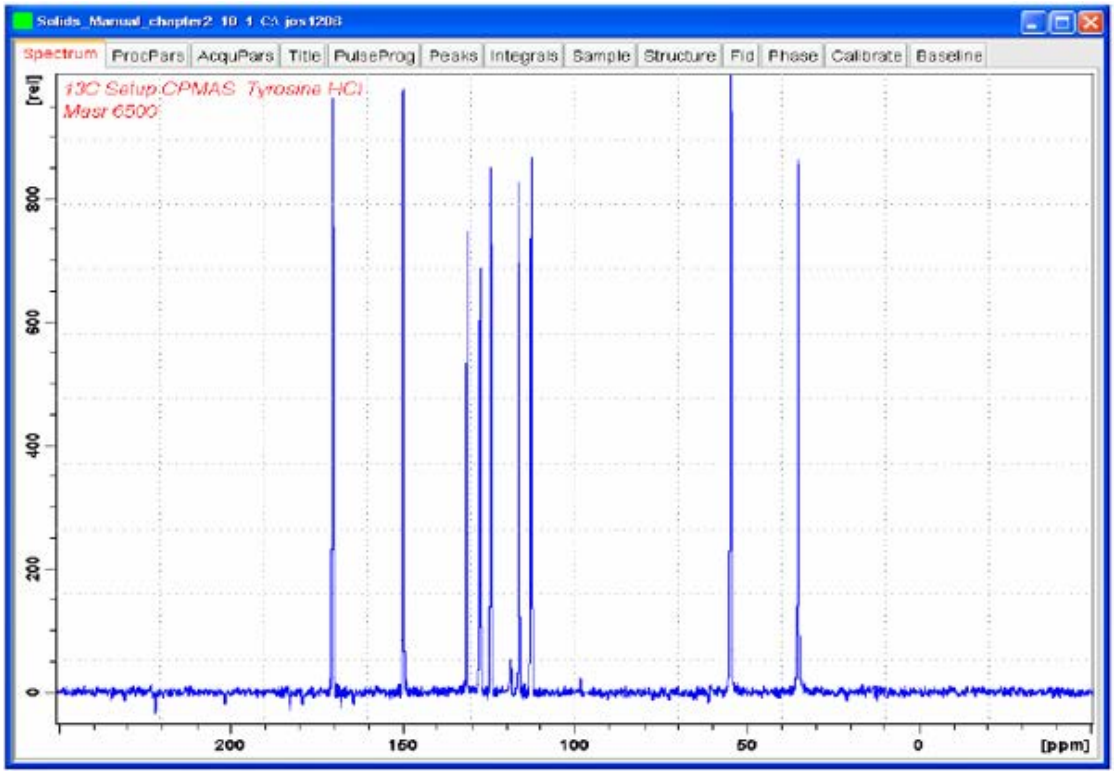


Figure 7.5: CPTOSS Experiment on Tyrosine HCl at 6.5 kHz

The figure above is a CPTOSS experiment on tyrosine HCl at 6.5 kHz sample rotation using a 4 mm CP-MAS triple resonance probe at 500 MHz with 256 accumulated transients. Spinning sideband residuals can be observed outside a noise level of approximately 2% peak to peak compared to the highest intensity. The residual sidebands have up to 5% intensity compared to the highest resonance.

7.3 SELTICS

Like the TOSS experiment, SELTICS (**S**ideband **EL**imination by **T**emporary **I**nterruption of the **C**hemical **S**hift) is an experiment for spinning sideband suppression. Pulses on the 13C channel are driven with **pl11** and pulse times are rotor synchronized. For optimum suppression, the shortest pulse (τ/24) of the sequences, where τr is the rotor period, should be a π/2 pulse or stronger. Choose **pl11** accordingly. Unlike TOSS, SELTICS is only 0.5 rotor periods long.

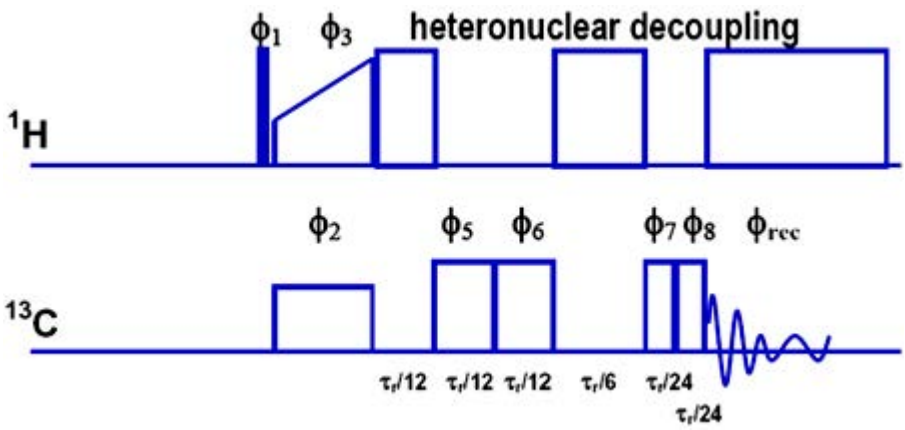


Figure 7.6: Pulse Program for SELTICS

In the figure above one can see that the SELTICS experiment takes only ½ rotor period compared to the 2 rotor periods required in the TOSS experiment.

Use glycine or tyrosine.HCl at a reasonable spinning speed.

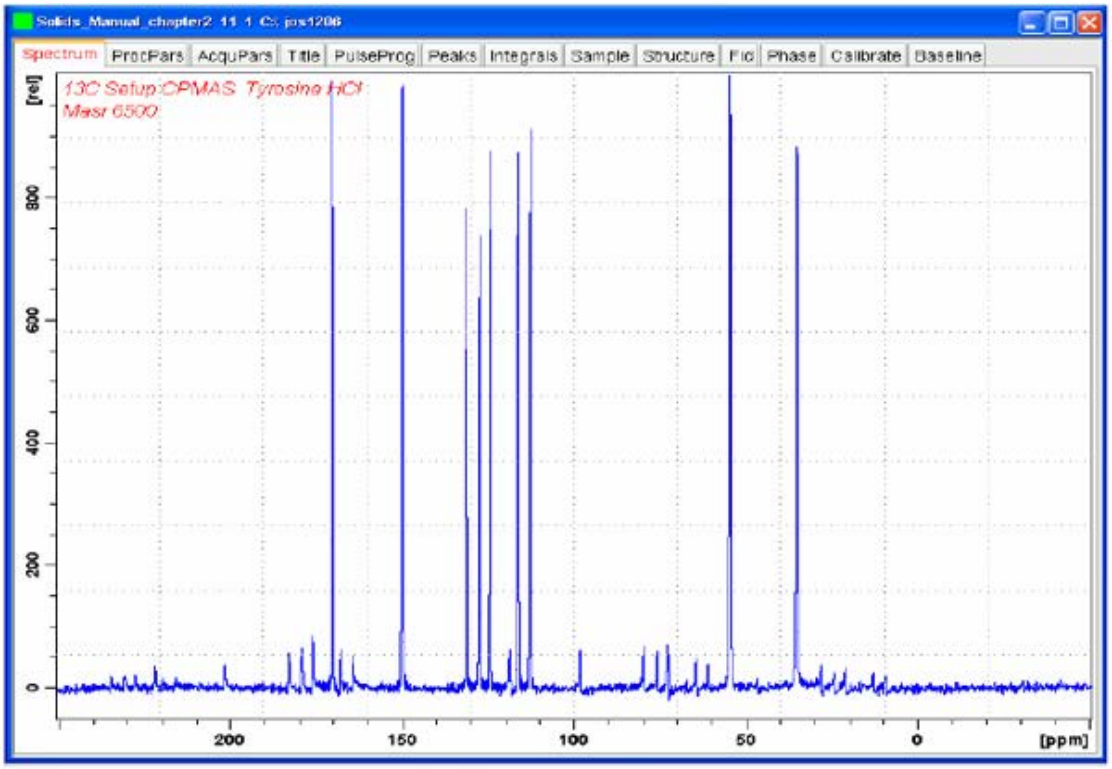


Figure 7.7: SELTICS at 6.5 kHz Sample Rotation on Tyrosine HCl

In the figure above the amplitude of the spinning sidebands are reduced to more than 10% compared to the original spectrum without sideband suppression, 256 transients were recorded.

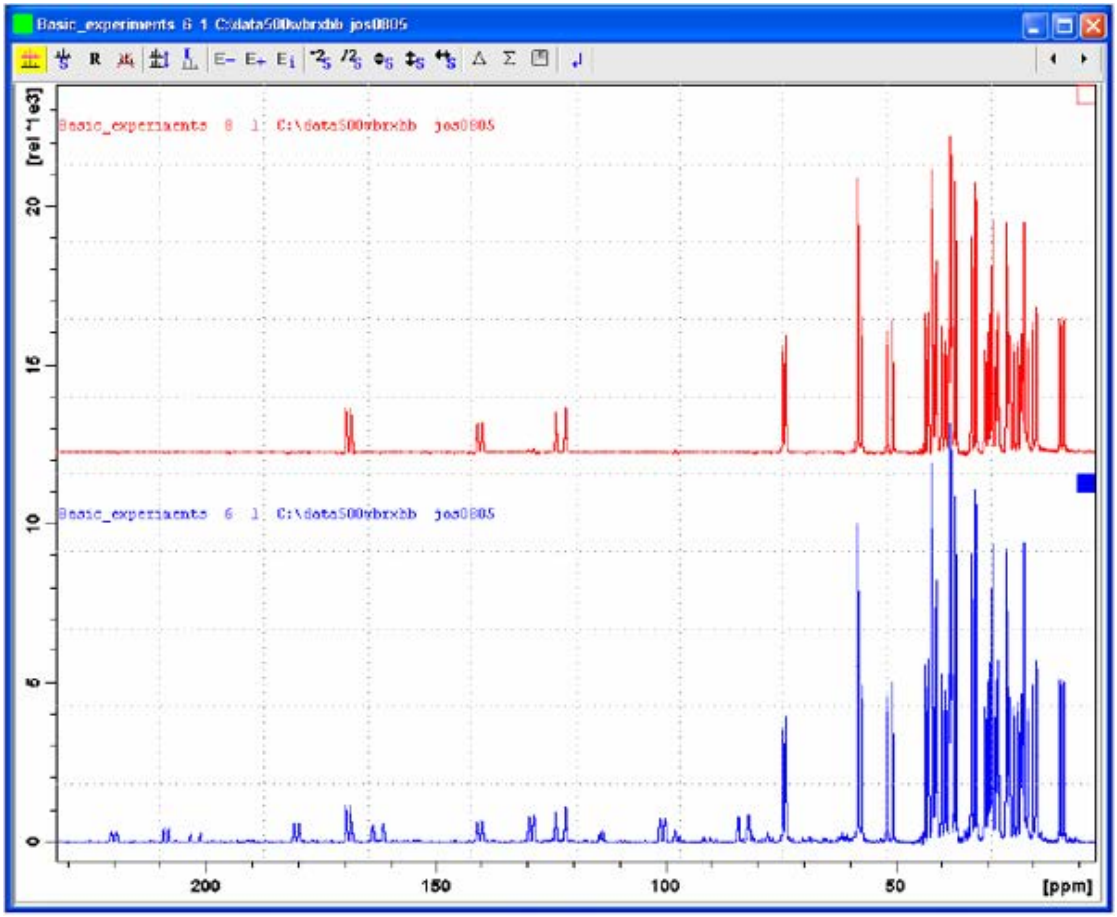


Figure 7.8: Cholesterylacetate Spectrum Using Sideband Suppression

In the figure above is a cholesterylacetate spectrum using sideband suppression with the SELTICS sequence at 5 Hz sample rotation (upper spectrum). The lower spectrum is the CP-MAS spectrum at 5 kHz sample rotation.

7.4 Non-Quaternary Suppression (NQS)

The NQS experiment is a simple spectral editing experiment. It relies on the fast dephasing of rare spins coupled to 1H spins through the heteronuclear dipolar interaction. For the dephasing delay **d3** one uses between 30 and 80 µs.

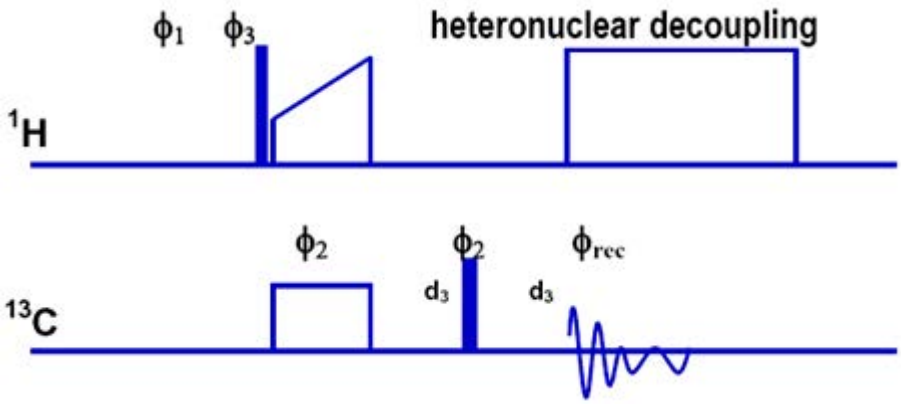


Figure 7.9: Block Diagram of the Non-quaternary Suppression Experiment

The **non-quatemary suppression** experiment is also called the **dipolar dephasing** experiment.

Use glycine or tyrosine spinning at 11 kHz as before.

Parameter	Value	Comments
pulprog	cpnqs, cptoss_nqs	
p2		180° pulse on X nucleus.
pl11		Power level driving P2 on X-channel.
d3	30 – 80 µs	Dephasing delay.

Table 7.3: Acquisition Parameters

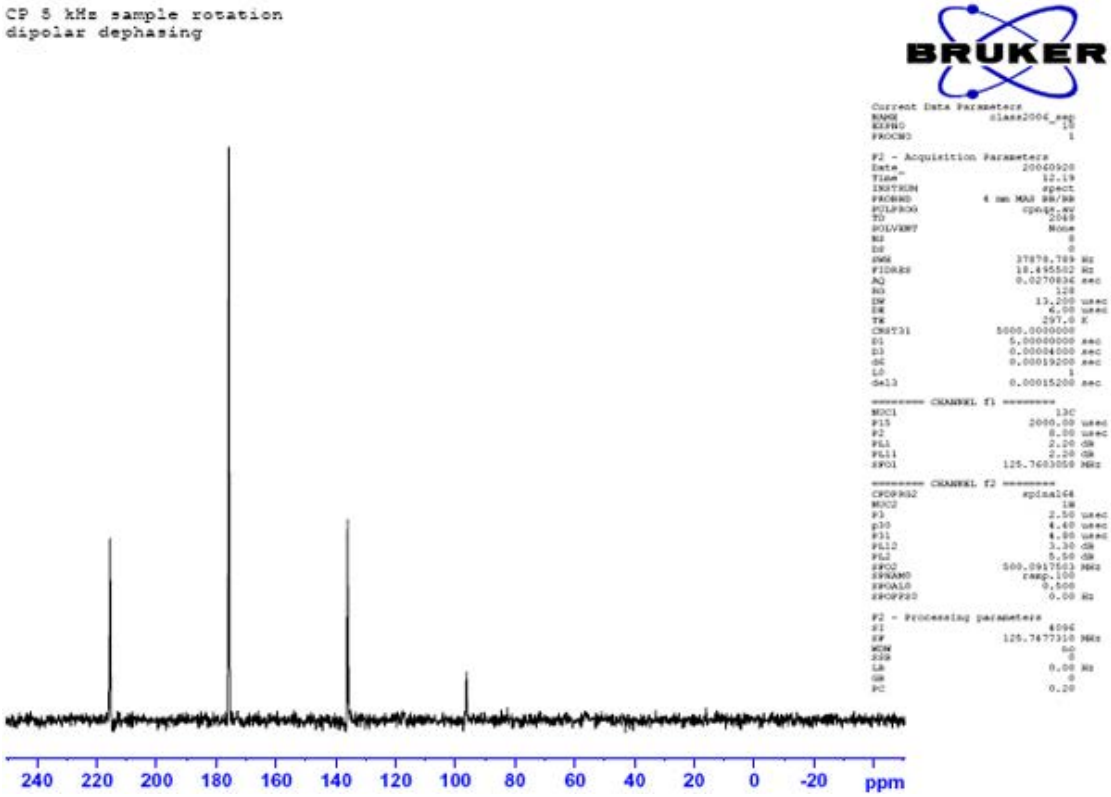


Figure 7.10: Glycine 13C CP-MAS NQS Experiment with a Dephasing Delay

In the figure above is a glycine ^{13}C CP-MAS NQS experiment with a dephasing delayd3 = 40 µs so that the total dephasing time is 80 µs. Spinning sidebands are still visible.

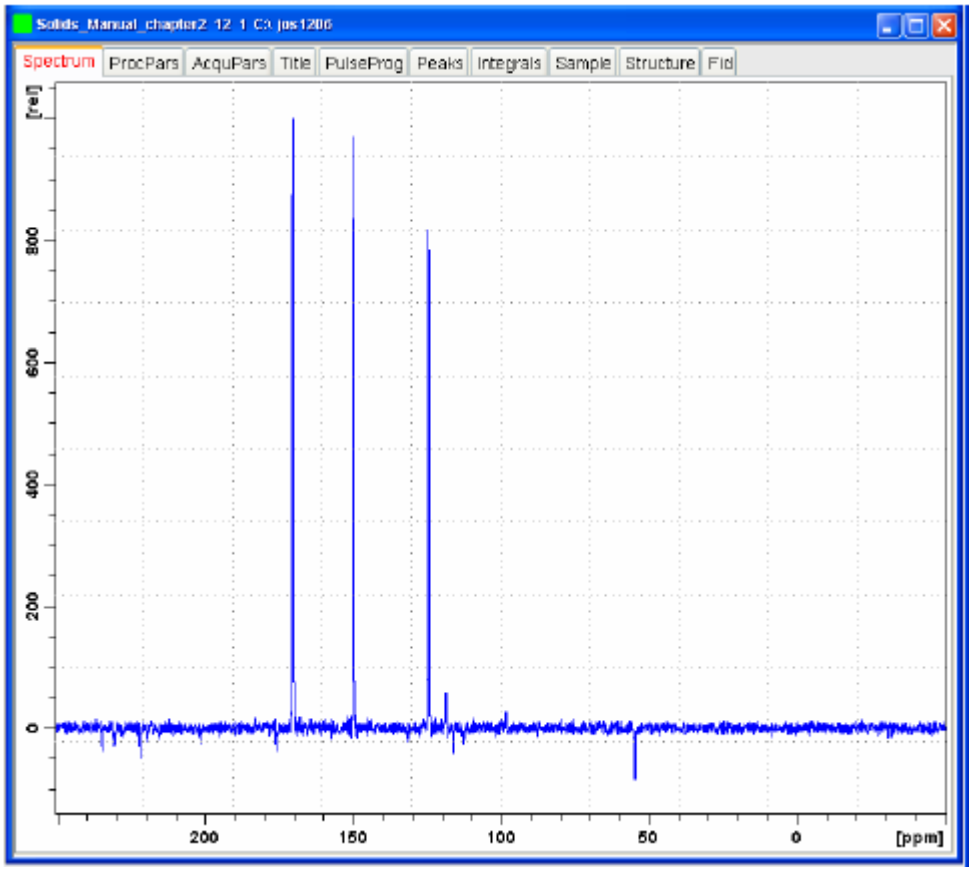


Figure 7.11: Tyrosine 13C CP-MAS NQS Experiment with TOSS

In the figure above is a tyrosine 13C CP-MAS NQS experiment with TOSS using a dephasing delay d3 = 60 µs. Spinning sidebands are suppressed for a clean spectrum. In this experiment the total dephasing time is 20 µs shorter than that used for the CPNQS experiment on glycine in the previous figure.

7.5 Spectral Editing Sequences: CPPI, CPPISPI and CPPIRCP

These spectral editing sequences help to distinguish CH, CH2, CH3 and quaternary carbons in 13C spectra. Common to all are various polarization and depolarization times, which properly mixed and combined give a series of spectra, which can be added and subtracted in order to obtain the various sub-spectra. All these sequences use constant-amplitude CP, which should be adjusted for maximum signal intensity. For the CPPI and CPPISPI sequences, the only parameter needed in addition to the CP parameters is `p16`. A good starting value is 40 µs to give a null signal for CH, a negative signal for CH2, and a positive signal for C and CH3. If necessary, this value can be optimized on a sample similar to the sample of interest for better editing. The p17, for the repolarization step in the CPPIRCP experiment, is about 10 – 20 µs.

For this experiment to succeed reliably, one should use moderate spinning speeds around (up to) 10 kHz. At slow rotation rates, no advantage was found in measuring the exact HH match. Running the experiment at constant amplitude CP, optimized for maximum signal, proved to be sufficient.

References for these experiments:

X. Wu, K. Zilm, Complete Spectral Editing in CP-MAS NMR, J. Magn. Reson. A 102, 205-213 (1993);

X. Wu, K. Zilm, Methylene-Only Subspectrum in CP-MAS NMR, J. Magn. Reson. A 104, 119-122 (1993);

X. Wu, S.T. Burns, K. Zilm, Spectral Editing in CP-MAS NMR. Generating Subspectra Based on Proton Multiplicities, J. Magn. Reson. A 111, 29-36 (1994);

R. Sangill, N. Rastrup-Andersen, H. Bildsoe, H.J. Jakobsen, and N.C. Nielsen, Optimized Spectral Editing of 13C MAS NMR Spectra of Rigid Solids Using Cross-Polarization Methods, J. Magn. Reson. A 107, 67-78 (1994).

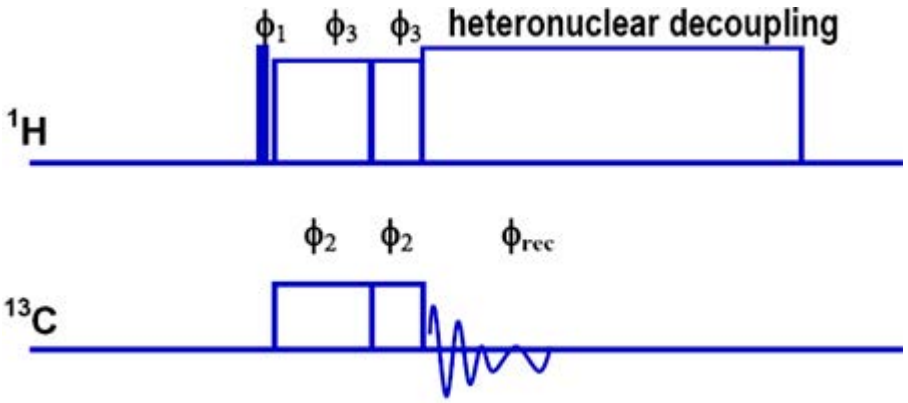


Figure 7.12: Block Diagram of the CPPI Experiment

Typical pulse widths for the second part of the CP pulse with the phase inversion are 40 µs.

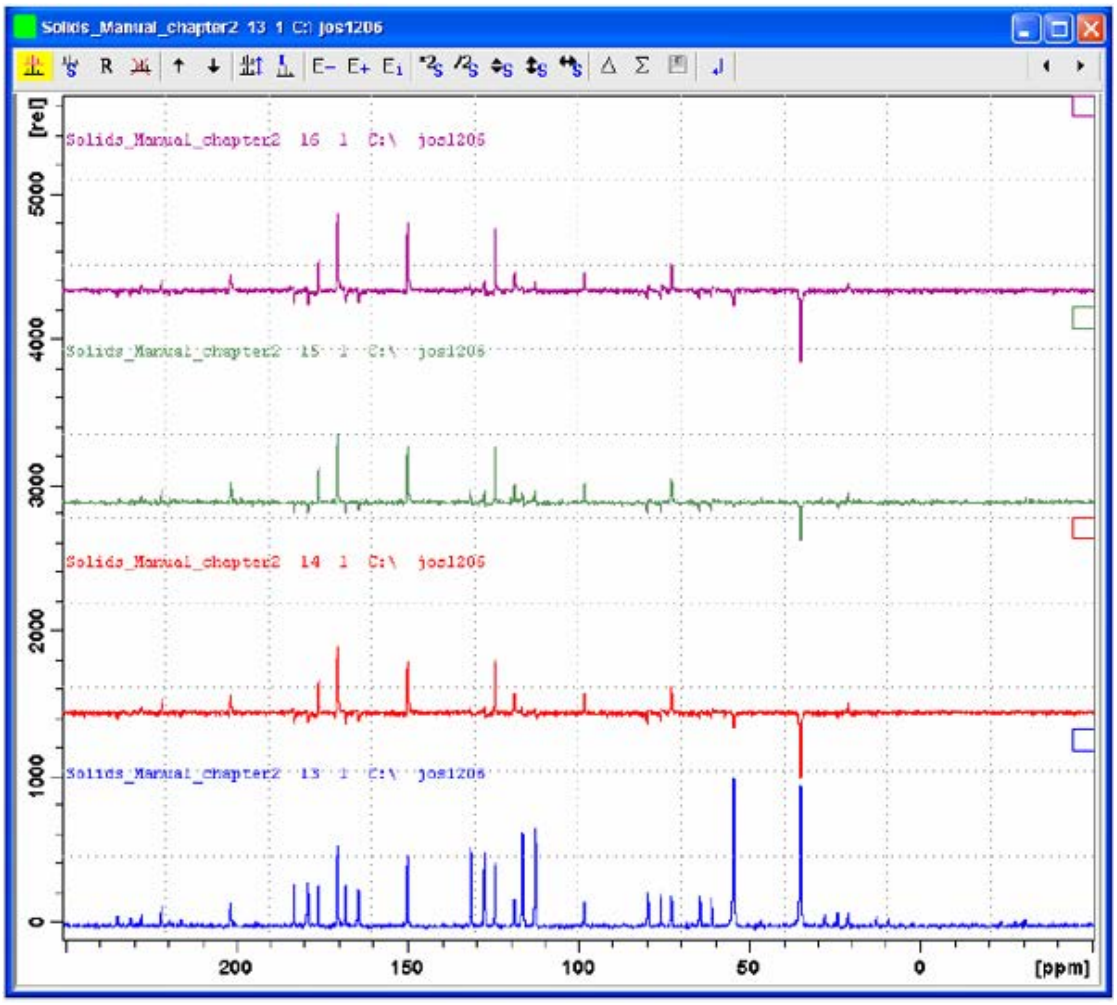


Figure 7.13: CP-MAS Spectrum of Tyrosine.HCl at 6.5 kHz

The figure above shows a CP-MAS spectrum of tyrosine.HCl at 6.5 kHz sample rotation obtained on a 500 WB spectrometer using a 4 mm CP-MAS double resonance probe. The red (third) spectrum is a CPPI spectrum where we see the CH_2 resonance at 35 ppm with a negative intensity. The aromatic CH resonances are clearly suppressed, where the Ca shows a slightly negative intensity. The polarization inversion pulse p16 was 40 µs long. The green (second) spectrum is a CPPIRCP experiment with p16=40 µs and p17=10 µs for better nulling of CH resonances, but in this case the aromatic CH resonances gained some intensity back. The purple (first) spectrum is a CPPISPI experiment with a similar performance as for the CPPI spectrum. Our experience is that one can adjust p15 (= 1 ms in this spectrum), p16 (= 30 µs in the purple spectrum) and so edit for pure CH resonances for example. Such tuning needs to be done of course on a known sample, which behaves similarly to the one under investigation, for the editing to be conclusive and correct.

Note: For more editing experiments consider the *Solid State Attached Proton Test* experiment, using the `sostapt` pulse program name, or look at 2D editing sequences, based on the FSLG HETCOR experiment.

8 FSLG-HETCOR

This chapter discusses setup and use of the **F**requency **S**witched **L**ee **G**oldburg **Het**eronuclear **Cor**relation (FSLG HETCOR) experiment.

The FSLG Hetcor experiment correlates ^{1}H chemical shifts with X-nuclei (e.g. ^{13}C, ^{15}N) chemical shifts. The experiment provides excellent ^{1}H resolution in the indirect dimension. Homonuclear decoupling in the 1H evolution period is achieved with an FSLG pulse train. FSLG permits relatively high spinning speeds and makes this experiment available for high field systems, requiring high spinning speeds in order to move spinning sidebands out of the spectral region. Decoupling the protons from the coupled X-nucleus during evolution is not essential, since the high spinning speed already achieves that. One can however improve the heteronuclear decoupling by a π-pulse in the middle of the evolution period (see A. Lesage et. al.).

Mixing is achieved during the cross polarization contact time. Since magnetization transfer from protons to X (e.g. ^{13}C) occurs rapidly, contact times should be kept short in order to avoid long range transfer, leading to unspecific cross peak patterns since the magnetization then has time to flow from any proton to any X-nucleus. A modification of the basic sequence uses cross polarization under a LG frequency offset for the protons. In this case, the proton magnetization detected by the X-nucleus comes from close protons only, since the proton spin lock at an LG offset interrupts the “communication” between the proton spins. A third modification of the basic sequence uses phase modulated pulses instead of frequency shifts. These three modifications to the basic sequence are described in the next chapter.

References

H.J.M. deGroot, H. Förster, and B.-J. van Rossum, Method of Improving the Resolution in Two-Dimensional Heteronuclear Correlation Spectra of Solid State NMR, United States Patent No 5,926.023, Jul. 20, 1999.

B.-J. van Rossum, H. Förster, and H.J.M. deGroot, High-field and high-speed CP-MAS 13C NMR Heteronuclear dipolar-correlation spectroscopy of solids with frequency-switched Lee-Goldburg homonuclear decoupling, J. Magn. Reson. A 120, 516-519 (1997).

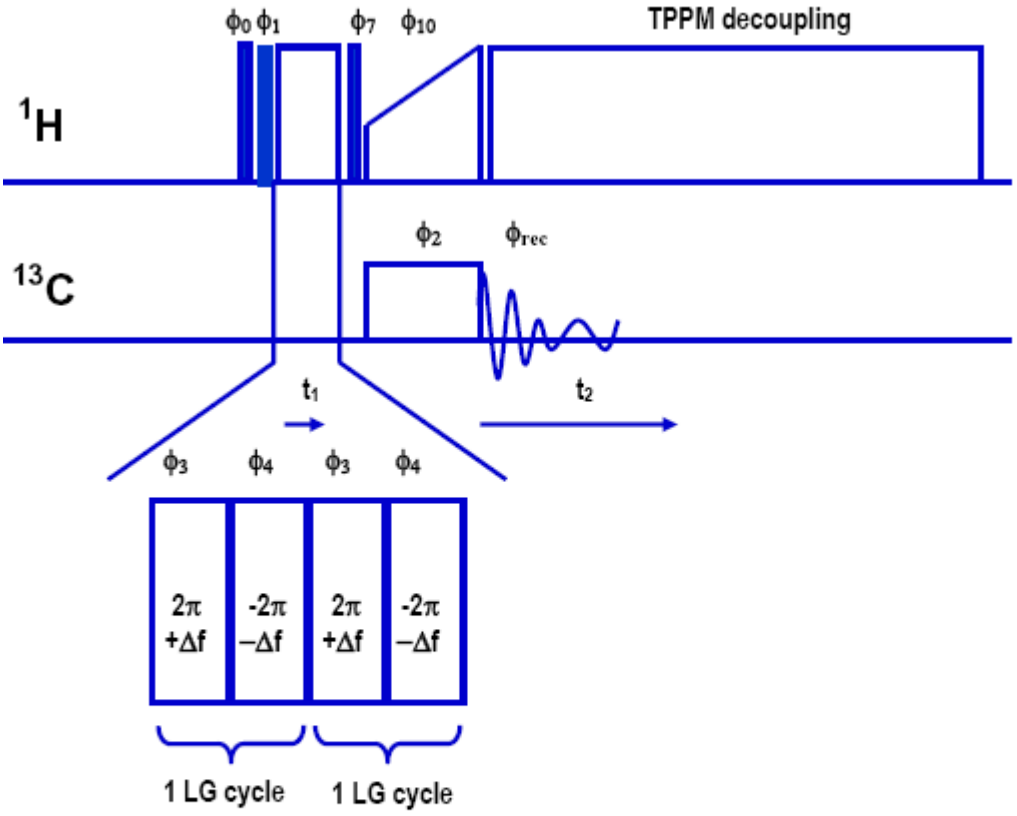
B.–J. van Rossum, Structure refinement of photosynthetic components with multidimensional MAS NMR dipolar correlation spectroscopy, Thesis, University of Leiden, Holland; (2000).

B.–J. van Rossum, C.P. deGroot, V. Ladizhansky, S. Vega, and H.J.M. deGroot, A Method for Measuring Heteronuclear (1H-13C) Distances in High Speed MAS NMR, J. Am. Chem. Soc. 122, 3465-3472 (2000).

D.P. Burum and A. Bielecki, An Improved Experiment for Heteronuclear-Correlation 2D NMR in Solids, J. Magn. Res. 94, 645-652 (1991).

A. Lesage and L. Emsley, Through-Bond Heteronuclear Single-Quantum Correlation Spectroscopy in Solid-State NMR, and Comparison to Other Through-Bond and Through-Space Experiments, J. Magn. Res. 148, 449-454 (2001).

8.1 Pulse Sequence Diagram for FSLG HETCOR



ϕ_0= 1 3 + STATES-TPPI (t_1)	ϕ_1= 1	ϕ_{10}= 0
ϕ_2 = 0 0 2 2 1 1 3 3	ϕ_4 = 2 (-LG)	ϕ_3 = 0 (+LG)
ϕ_7= 3	ϕ_{rec}= 0 2 2 0 1 3 3 1	

Figure 8.1: The FSLG Hetcor Experiment

The FSLG Hetcor experiment consists of 3 basic elements, the homonuclear decoupling sequence during which the ^{1}H chemical shifts evolve, the cross polarization sequence, during which the information of the ^{1}H spin magnetization is transferred to the X-spins, followed by observation of the X-spins under proton decoupling.

8.2 Setting up FSLG HETCOR

- This experiment requires a probe of 4 mm spinner size or smaller. One can run it on a 7 mm probe, but the results will not be very convincing.
- Start from a data set with well adjusted cross polarization and proton decoupling at fairly high RF-fields. Unlike standard multiple pulse decoupling, which only works well at very high RF-fields, FSLG requires only moderately high RF fields. Decent performance is achieved at 80-100 kHz proton field. At lower magnetic fields (200-300 MHz proton frequency) lower RF-fields are adequate, RF fields of 100 kHz and higher perform better at higher magnetic fields (500 MHz and up).
- Insert a suitable test sample, spin at a suitable speed. We recommend ^{13}C labeled tyrosine hydrochloride, since it has a wide spread (2.5-12ppm) of proton shifts, a short proton T_1, a well resolved ^{13}C-spectrum with many lines, and it is readily available. The unlabeled sample can also be used, but requires a few more scans (8-32).
- Optimize the spin rate such that no overlap occurs between centerbands and sidebands (especially with the labeled sample, in order to avoid rotational resonance broadening). Re-optimize decoupling and HH-condition. Check the proton RF-field via the proton 90° pulse **p3**. Set **pl13** = **pl12**, set **cnst20** to RF-field in Hz as calculated from **p3**.
- Generate a new data set with **edc**, **new**. **Set pulprog lghetfq** and change to a 2D parameter set using the **123** button in **eda**. Set **FnMode** to STATES-TPPI. Type **ased** or click the pulse symbol in **eda**.

Figure 8.2: The "12..." icon, and the ased icon in eda

- Performing **ased** will show all parameters which are essential for the acquisition, rather than all available parameters. In addition, it performs calculations which are specified in the pulse program. Note that all parameters which are calculated are not editable, and will show only, if explicitly used during the main pulse program between ze and exit. In this sequence, the proton chemical shift evolution is influenced by the RF field (**cnst20**) under which the shifts evolve and the type of homonuclear decoupling sequence (FSLG in this case) which scales chemical shifts (by about 0.578 in this case). In order to obtain proton chemical shifts at the standard scale, both parameters are taken into account and an increment along F1 is calculated which yields correct chemical shifts for protons. Transfer this increment to **IN_F1** in **eda** (a button in ased). This will set the sweep width along F1. Note that the time increment here is generated by a loop counter, counting the periods of FSLG. The loop counter l3 is used to multiply this increment. Usually, l3 is set to 2-4 in order to reduce the F1 sampling width to a reasonable value. **Cnst24** is usually set to -1000 - -2000 in order to move the spectrum away from the center ridge in F1.

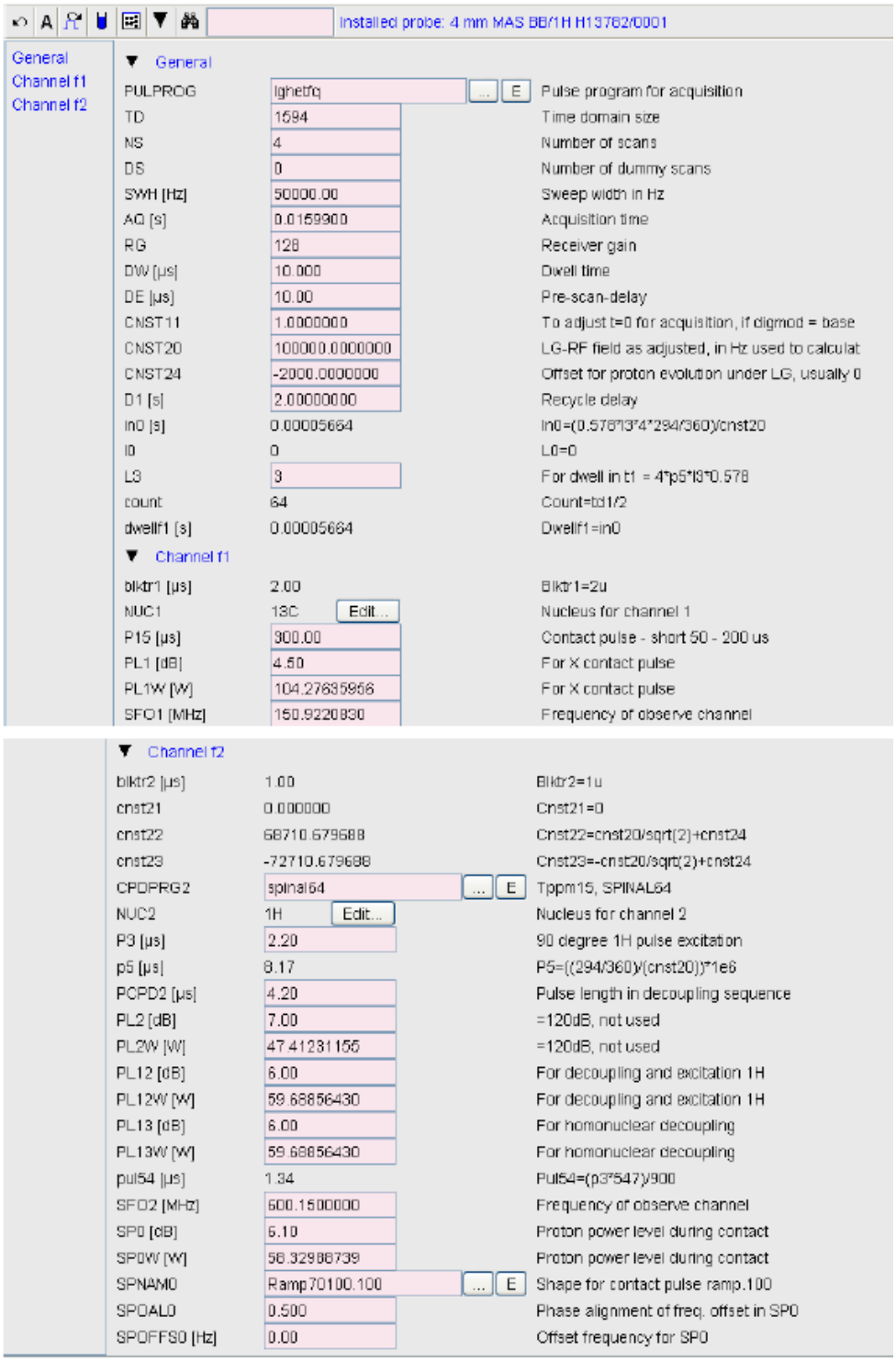


Figure 8.3: The ased Display

In the figure above the frequency offsets for the FSLG part are shown as **cnst22**, **cnst23**. They are different because **cnst24** shifts the center frequency by 2000 Hz.

Parameter	Value	Comments
pulprog	lghetfq	FSLG program.
nuc1	13C	
o1p	100 ppm	
nuc2	1H	
cnst20	70-100000	Proton spin nutation frequency with PL13.
cnst24	0 - -2000	Place carrier off during evolution.
pl1		Power level channel 1 for contact pulse.
pl12		Power level channel 2 TPPM/SPINAL decoupling.
pl13		Power level channel 2 FSLG decoupling.
sp0		Power level channel 2 for contact pulse.
spnam0	ramp.100 or simil.	Shape for contact pulse channel f2.
p3	2.5 – 3 µsec	90º pulse channel 2 at pl12.
p15	50 -500 µsec	Contact pulse width.
pcpd2	≈ 2*p3-0.2	SPINAL64 /TPPM decoupling pulse.
cpdprg2	SPINAL64/ TPPM15	Decoupling sequence.
F1 1H indirect		
l0	0	Start value 0, incremented during expt.
l3	2- 4	Multiples of FSLG-periods, increment per row.
in_f1	=in0 as calculated	Set according to value calculated by ased.
F2 13C acquisition		
d1	2s	Recycle delay.
sw	310ppm	Sweep width direct dimension.
aq	16-20 msec	
masr	10000-15000 Hz	At 100 kHz RF, 15 kHz is okay.

Table 8.1: Acquisition Parameters for FSLG-HETCOR (on tyrosine-HCl)

Parameter	Value	Comment
F2 direct dim 13C		
si	2-4 k	
wdw	QSINE	SSB 2, 3 or 5
ph_mod	pk	
F1 indirect 1H		
si	256-1048	
mc2	STATES-TPPI	
wdw	QSINE	
ssb	3, 5	
ph_mod	pk	

Table 8.2: Processing Parameters for FSLG-HETCOR (on tyrosine-HCl)

8.3 Results

A full plot of a FSLG-HETCOR on labeled tyrosine-HCl is shown in the next figure.

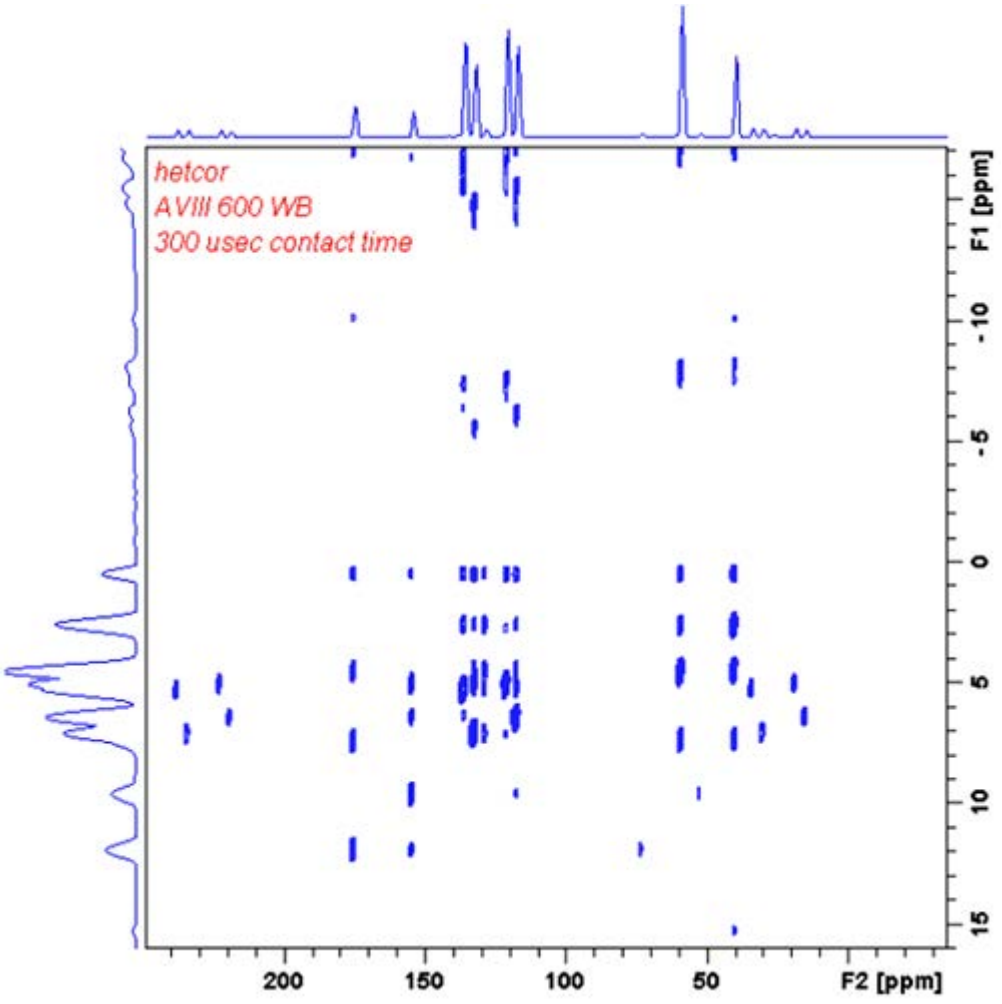


Figure 8.4: FSLG Hetcor Spectrum Tyrosine HCl

The figure above shows a FSLG Hetcor Spectrum Tyrosine HCl with parameters as shown in the tables in section *Setting up FSLG HETCOR [▶ 107]*. Full transform with slight resolution enhancement, **qsine/SSB**=3. Proton shifts calibrated as 2.5 and 12 (most high field/low field peak). Center ridge at 0 ppm along F1 is spin locked signal which does not follow the FSLG-evolution. **Cnst24** is used to separate the proton spectrum from this ridge. The contact time of 300 µsec shows many long range couplings. The next figure shows the region of interest excluding the center ridge and the spinning sidebands.

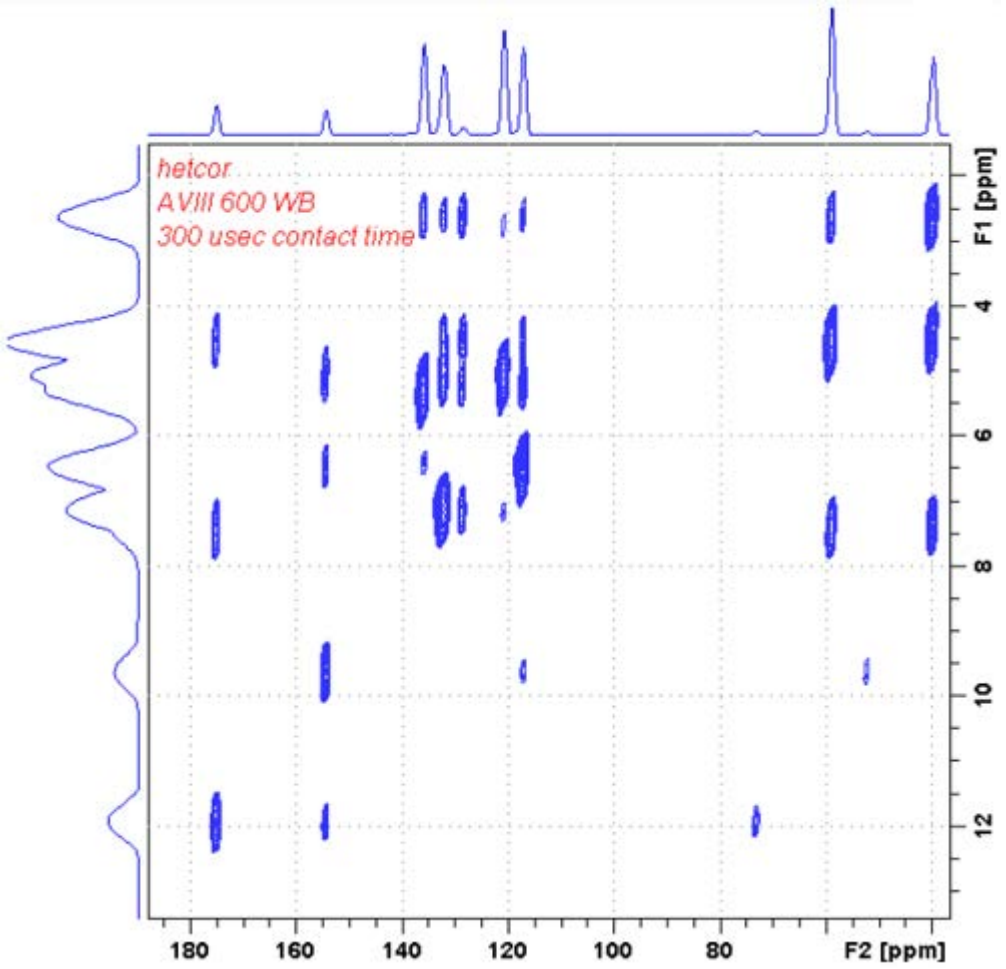


Figure 8.5: FSLG Hetcor Spectrum Tyrosine HCl

The figure above shows a FSLG Hetcor Spectrum Tyrosine HCl with parameters as shown in the tables in *Setting up FSLG HETCOR [▶ 107]*. Full transform with slight resolution enhancement, qsine/SSB=3. Proton shifts calibrated as 2.5 and 12 (most high field/low field peak) expansion plot.

9 Modifications of FSLG HETCOR

The basic HETCOR sequence can be improved in several respects. The protons which are observed are all coupled to ^{13}C carbons (since we observe these). So the proton shifts also evolve under the residual dipolar coupling and the J-coupling to ^{13}C. This can be refocused by a ^{13}C π-pulse in the middle of the proton evolution. The pulse program `lghetfqpi` will serve this purpose.

Furthermore, it may be desirable to compare the proton shift spectrum obtained with X-observation (HETCOR) with the proton spectrum obtained by CRAMPS techniques (refer to the chapters on CRAMPS later in this manual), observing the protons directly. Usually, these experiments use phase modulated shapes (PMLG or DUMBO). In order to make both experiments comparable, it is favored to use the same type of proton shift evolution in both sequences. The pulse programs which use phase gradient shapes to achieve homonuclear proton dipolar decoupling are `pmlghet` and `wpmlghet`. If DUMBO decoupling is desired, the pulse programs are `dumbohet` or `edumbohet`. These pulse programs use either windowless pulse trains, or windowed pulse trains which can be timed in exact analogy to the CRAMPS-type sequences wpmlg*2 and dumbo*2. These sequences also suppress the center ridge efficiently so that the carrier frequency need not be shifted out of the proton range during evolution (cnst24=0). In contrary, it is possible to shift the carrier to the proton shift range center.

A third modification addresses the problem of poor discrimination between sites which are strongly and weakly coupled to protons. In the standard sequence this is solely achieved setting contact times short. Of course, this reduces cross peaks from remote couplings more than it reduces cross peaks from directly bonded protons. However, the remote couplings are always present through the homonuclear coupling between all protons. These couplings can however be suppressed by executing the contact with a Lee-Goldburg proton offset. Then the protons are homonuclear decoupled, and the transfer of protons to X only follows the heteronuclear dipolar coupling between these. The pulse program `lghetfqlgcp` works completely in analogy to `lghetfq`, but executes the contact at a proton offset calculated from the proton RF field during the spin lock contact pulse. This modifies the HH condition, which must be reestablished using the pulse program lgcp.

In the following sections, the specifics of these modified sequences are discussed.

9.1 Carbon Decoupling During Evolution

The only difference between `lghetfq` and `lghetfqpi` is the decoupling π-pulse at the center of the evolution period. All that needs to be set in addition is the X- π-pulse **p2** at power level **pl1**. At fast spin rates and in fully labeled samples, the narrowing effect on the proton spectrum may be small or not noticeable, but on samples with natural abundance it may be noticeable. At long contact times and transfer from many different protons, the line width in the proton spectrum may also be insensitive. In the following figure are two columns with the highest field aliphatic peak in tyrosine-HCl are shown. The π-decoupled trace (red) is clearly narrower.

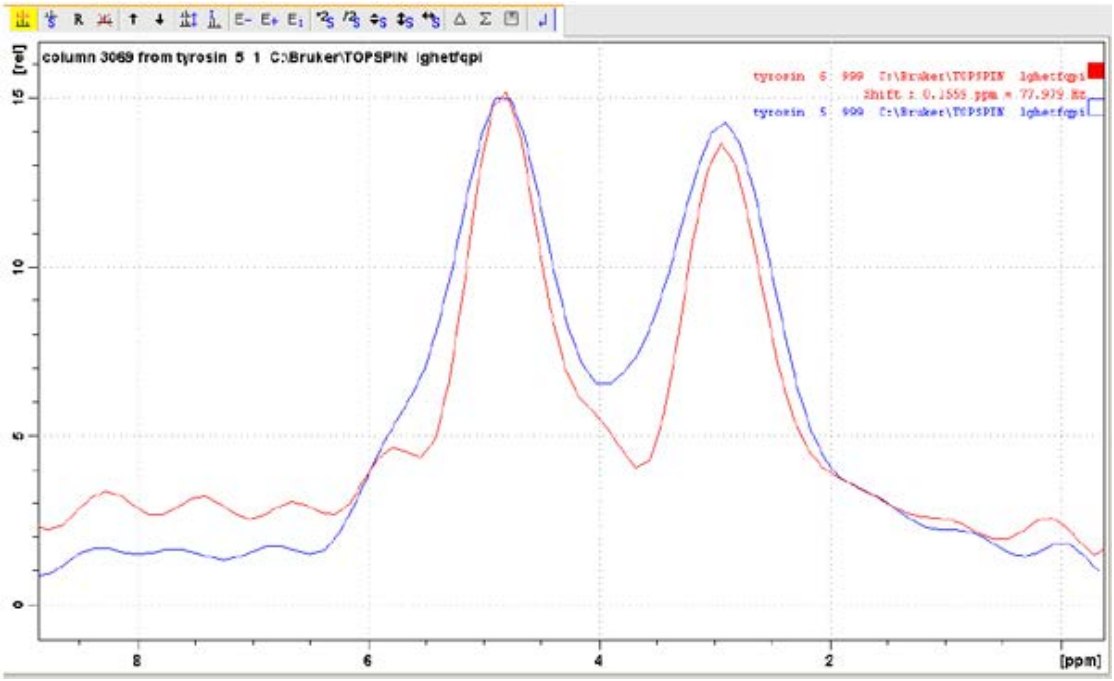


Figure 9.1: Comparison of HETCOR with and without 13C-decoupling

The figure above shows a comparison of HETCOR with and without 13C-decoupling. Natural abundance tyrosine-HCl was run with 50 µsec contact time.

Reference:

A. Lesage and L. Emsley, Through-Bond Heteronuclear Single-Quantum Correlation Spectroscopy in Solid-State NMR, and Comparison to Other Through-Bond and Through-Space Experiments, J. Magn. Res. 148, 449-454 (2001).

9.2 HETCOR with DUMBO, PMLG or w-PMLG, Using Shapes

These sequences use phase modulated shapes for homonuclear proton decoupling. Apart from some smaller differences, the sequences are in complete analogy to the HETCOR sequence using frequency shifts.

The only differences between these sequences lie in the length and type of shape used for homonuclear decoupling. DUMBO and e-DUMBO (Emsley et al.) use principles known from multiple pulse NMR operating on resonance, whereas pmlg and w-pmlg (Vega et al.) use phase ramps which act like frequency offsets and are therefore derivatives of FSLG.

References

D. Sakellariou, A. Lesage, P. Hodgkinson and L. Emsley, *Homonuclear dipolar decoupling in solid-state NMR using continuous phase modulation*, Chem. Phys. Lett. 319, 253 (2000).

Vinogradov, E.; Madhu, P. K.; Vega, S., *High-resolution proton solid-state NMR spectroscopy by phase modulated Lee-Goldburg experiment*, Chem. Phys. Lett. (1999), 314(5,6), 443-450.

E. Vinogradov, P.K. Madhu and S. Vega, *Proton spectroscopy in solid state NMR with windowed phase modulated Lee-Goldburg decoupling sequences*, Chem. Phys. Lett. (2002), 354, 193.

Leskes, Michal; Madhu, P. K.; Vega, Shimon, *A broad banded z-rotation windowed phase modulated Lee-Goldburg pulse sequence for 1H spectroscopy in solid state NMR*, *Chem. Phys. Lett. (2007), 447, 370-374.*

Leskes, Michal; Madhu, P. K.; Vega, Shimon, *Supercycled homonuclear decoupling in solid state NMR: towards cleaner 1H spectrum and higher spinning rates*, J. Chem. Phys. (2007) in press.

9.2.1 The Sequence pmlghet

This sequence uses windowless phase ramped shapes. One can write these shapes as multiples of FSLG cycles to manipulate the length of the T_1-increment. Usually, 2 FSLG cycles make sense. The pulse program calculates the required shape pulse length from the RF-field during the FSLG-evolution. In pmlghet, a shape with 2 cycles is assumed in the calculation. The sequence is optimized for a simple twofold linear phase ramp (supplied as **lgs-2**). The carrier may be placed in the middle of the proton spectrum during evolution which may allow using fewer increments and therefore saving time. However, one should be aware of the presence of proton spinning sidebands along F1 which may inappropriately fold in if the spectrum window selected along F1 is too small.

Processing is done in complete analogy to the FSLG-experiment, as for all following sequences.

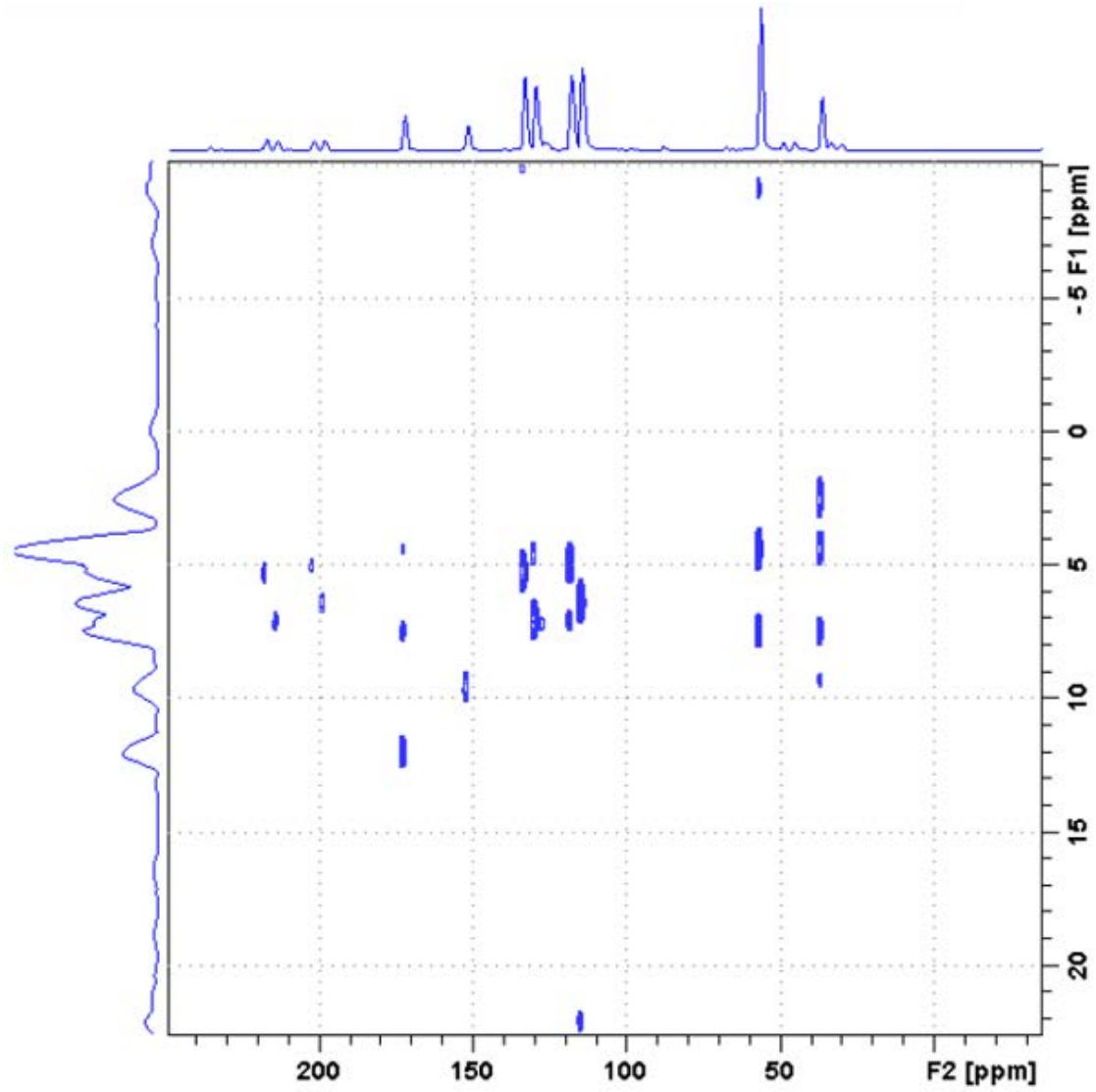


Figure 9.2: HETCOR Using Windowless Phase Ramps

The figure above shows the HETCOR using windowless phase ramped shapes for proton homonuclear decoupling during evolution. The carrier was placed in the middle of the proton spectrum and the usual carrier ridge was suppressed by phase cycling (Leskes et al., Chem. Phys. Lett.). This allows reduced measurement times. `Pmlghet`, `wpmlghet`, `dumbohet` and `edumbohet` should give rather similar spectra.

Parameter	Value	Comments
pulprog	pmlghet	Using phase ramps.
nuc1	13C	
o1p	100 ppm	
nuc2	1H	
cnst20	80-100000	Proton spin nutation frequency with PL13.
cnst24	1000-3000	Place carrier within proton spectrum for evolution.
pl1		Power level channel 1 for contact pulse.
pl12		Power level channel 2 TPPM/SPINAL decoupling.
pl13		Power level channel 2 PMLG decoupling.
sp0		Power level channel 2 for contact pulse.
spnam0	ramp.100 or similar	Shape for contact pulse channel f2.
sp1	set to pl13	To match cnst20.
spnam1	lgs-2	To match ppg calculation of in0, in_f1.
p3	2.5 – 3 µsec	90º pulse channel 2 at pl12.
p15	50 -500µs	Contact pulse width.
pcpd2	≈ 2*p3	SPINAL64 /TPPM decoupling pulse.
cpdprg2	SPINAL64/ TPPM15	Decoupling sequence.
F1 1H indirect		
l0	0	Start value 0, incremented during expt.
l3	2- 4	Multiples of FSLG-periods, increment per row.
in_f1	=in0 as calculated	Set according to value calculated by ased.
F2 13C acquisition		
d1	2s	Recycle delay.
sw	310ppm	Sweep width direct dimension.
aq	16-20 msec	
masr	10-15 kHz	15 kHz ok at 100 kHz RF-field.

Table 9.1: Acquisition Parameters for pmlg-HETCOR (on tyrosine-HCl)

Parameter	Value	Comment
F2 direct dim 13C		
si	2-4 k	
wdw	QSINE	SSB 2, 3 or 5.
ph_mod	pk	
F1 indirect 1H		
si	256-1048	
mc2	STATES-TPPI	
wdw	QSINE	
ssb	3, 5	
ph_mod	pk	

Table 9.2: Processing Parameters for pmlg-HETCOR (as for FSLG on tyrosine-HCl)

9.2.2 w-pmlghet

If one wants to compare a solids proton spectrum, acquired via ^{13}C detection (using HETCOR) and a direct detect proton spectrum (using `w-pmlg`), it is useful to do this using exactly the same parameters (power levels and timings) in both experiments. If `w-pmlg` is used and optimized for the direct detect experiment, the same parameters will work with `w-pmlghet`, provided that both experiments are done with the same probe.

Parameter	Value	Comments
pulprog	wpmlghet	Phase ramp allows detection window.
nuc1	13C	
o1p	100 ppm	
nuc2	1H	
cnst20	usually >=100 kHz	As prepared with wpmlgd proton detect exp.
cnst24	1000-3000	Place carrier within proton spectrum for evolution.
pl1		Power level channel 1 for contact pulse.
pl12		Power level channel 2 TPPM/SPINAL decoupling.
pl13		Power level channel 2 w-PMLG decoupling.
sp0		Power level channel 2 for contact pulse.
spnam0	ramp.100 or similar	Shape for contact pulse channel f2.
sp1	set to pl13	To match cnst20.
spnam1	m5m or m5p	Both include one FSLG cycle.
p3	2.5 – 3 µsec	90º pulse channel 2 at pl12.
p15	50 -500 µsec	Contact pulse width.
pcpd2	≈ 2*p3	SPINAL64 /TPPM decoupling pulse.
cpdprg2	SPINAL64/ TPPM15	Decoupling sequence.
F1 1H indirect		

l0	0	Start value 0, incremented during expt.
l3	2- 4	Multiples of FSLG-periods, increment per row.
in_f1	=in0 as calculated	Set according to value calculated by ased.
F2 13C acquisition		
d1	2s	Recycle delay.
sw	310ppm	Sweep width direct dimension.
aq	16-20 msec	

Table 9.3: Acquisition Parameters for wpmlg-HETCOR (on tyrosine-HCl)

9.2.3 edumbohet

DUMBO decoupling is as efficient as PMLG decoupling. As will be discussed in chapters *CRAMPS 1D [▷ 247]* and *Modified W-PMLG [▷ 257]*, it requires a bit slower spin (up to 12-13 kHz) and to place the carrier closer to resonance. Faster spinning is possible with higher power pulses and shorter pulse widths (24 µsec or 16 µsec). The library of AU-programs in TopSpin includes `dumbo`, which calculates the desired shapes for windowed and windowless DUMBO shapes. If the windowless version is desired, the e-dumbo shape is preferred. Typing `xau dumbo` starts a dialog, in which e (for e-dumbo 22), 1 for the number of cycles, 64 for the number of steps, and 0 for an added angle (this value would be added to all phases in the shape). The program sets **p20** to 32 µsec as default. This is appropriate for an RF field of 100 kHz. `Spnam2` is set to edumbo22_1+0.

Parameter	Value	Comments
pulprog	edumbohet	Windowless dumbo shape.
nuc1	13C	
o1p	100 ppm	
nuc2	1H	
cnst24	1000-3000	Place carrier within proton spectrum for evolution.
pl1		Power level channel 1 for contact pulse.
pl12		Power level channel 2 TPPM/SPINAL decoupling.
pl13	for 100 kHz	Power level channel 2 DUMBO decoupling.
sp0		Power level channel 2 for contact pulse.
spnam0	ramp.100 or similar	Shape for contact pulse channel f2.
sp2	set to pl13	100 kHz for default duration 32 µs.
spnam2	edumbo22_1+0	Both include one e-DUMBO cycle.
p3	2.5 – 3 µsec	90° pulse channel 2 at pl12.
p15	50 -500µs	Contact pulse width.
p20	32 µs	32 µs for 100 kHz field, set by xau dumbo.
pcpd2	≈ 2*p3	SPINAL64 /TPPM decoupling pulse.
CPDPRG2	SPINAL64/ TPPM15	Decoupling sequence.
F1 1H indirect		

l0	0	Start value 0, incremented during expt.
l3	2- 4	Multiples of e-DUMBO period, increment per row.
in_f1	=in0 as calculated	Set according to value calculated by ased.
F2 13C acquisition		
d1	2s	Recycle delay.
sw	310 ppm	Sweep width direct dimension.
aq	16-20 msec	
masr	12000-13000	

Table 9.4: Acquisition Parameters for e-DUMBO-HETCOR (on tyrosine-HCl)

9.2.4 dumbohet

This is the windowed version of the previous experiment, analogous to wpmlghet. Run `xau dumbo`, select d (for dumbo), 1 (for 1 cycle), 0 (for added angle). The calculated shape, dumbo_1+0 will be entered as spnam1, p10 will be set to 32 µsec. The projection of this experiment can be compared to the result of a direct proton detected CRAMPS experiment, using dumbod2. At higher fields, p10 may have to be set to 24 rather than 32 µsec for better resolution. The same pulse length p10, shape, window (p9) and power level should be used in both experiments.

The resolution along the proton dimension is comparable for all these experiments. The FSLG experiment is the most forgiving, only requiring knowledge of the RF power level for decoupling at a certain RF field. Setting **cnst20** to this RF-field (+ 5 or 10%) is all that needs to be set, if the ^{13}C observe parameters are well adjusted.

Parameter	Value	Comments
pulprog	dumbohet	Windowed dumbo shape.
nuc1	13C	
o1p	100 ppm	
nuc2	1H	
cnst24	1000-3000	Place carrier within proton spectrum for evolution.
pl1		Power level channel 1 for contact pulse.
pl12		Power level channel 2 TPPM/SPINAL decoupling.
pl13		Power level channel 2 DUMBO decoupling.
sp0		Power level channel 2 for contact pulse.
spnam0	ramp.100 or similar	Shape for contact pulse channel f2.
sp1	set to pl13	100 kHz for default duration 32 µs.
spnam1	dumbo_1+0	Both include one DUMBO cycle.
p3	2.5 – 3 µsec	90º pulse channel 2 at pl12.
p15	50 -500µs	Contact pulse width.
p10	32 or 24µs	32 µs for 100 kHz field, 24 µs for better resolution at high magnetic fields (>500 MHz).
pcpd2	≈ 2*p3	SPINAL64 /TPPM decoupling pulse.

cpdprg2	SPINAL64/ TPPM15	Decoupling sequence.
F1 1H indirect		
l0	0	Start value 0, incremented during expt.
l3	2- 4	Multiples of DUMBO-periods, increment per row.
in_f1	=in0 as calculated	Set according to value calculated by ased.
F2 13C acquisition		
d1	2s	Recycle delay.
sw	310ppm	Sweep width direct dimension.
aq	16-20 msec	
masr	12000-13000	

Table 9.5: Acquisition Parameters for DUMBO-HETCOR (on tyrosine-HCl)

9.3 HETCOR with Cross Polarization under LG Offset

Usually, the cross polarization step is executed at less power than what is used for the initial excitation pulse and decoupling during observe. Therefore, a second LG condition must be set for the power level during contact. Furthermore, the HH condition must be re-established, since the proton spin lock pulse now must be a square pulse, not a ramp. The ramp can be transferred to the carbon (F1) side.

The following steps are involved:

1. The RF field for protons during contact must be measured and adjusted. With linearized transmitters, the required power level can be calculated from a known reference pulse width.
2. The LG frequency offset must be calculated from the RF field (RF/sqrt(2)).
3. The HH condition must be reestablished, varying the F1 (13C)-RF field.

As the effective field is the vector sum of RF-field and frequency offset, which is required to match for both nuclei at the HH condition, a higher power pulse is required on F1 with increasing offset. So it is recommended not to set the proton power higher than the 50 kHz RF field. For the setup, the pulse program `lgcp` is used. It contains an include file lgcalc2.incl which will calculate the LG offset from a given RF field specified as **cnst17**. It will set the calculated offset as **cnst19** during contact. In **ased**, it will also calculate **cnst16**, which shows the effective field under **cnst17** RF field and **cnst18**, **cnst19** offset. The X contact pulse is executed as (ramp)-shape. Any standard ramp is possible, but a flat ramp (70-100% or 90-100%) is preferred. Usually, calculating the required RF field for the HH match can be done in the following way:

1. Load the pulse program lgcp into a standard CP/MAS data set with all parameters set and optimized. Set p14 to about 54° flip angle. Use tyrosine *HCl or the sample of interest.
2. From p3, calculate the power level **pl2** for 50 kHz RF field, using calcpowlev.
3. Enter **cnst17**= 50000, type **ased**. Read the value of **cnst16** (=effective field under **cnst18**, **cnst19** offset (about 61000 Hz). Add or subtract the spin rate used, for instance spinning at 13 yields 74000 or 48000 Hz. If a 30% ramp is used, and **sp0** is set to 74000 (with a safety margin of 1000), the HH condition will cover 75000 down to 45000 Hz RF field. This includes both HH sidebands n=±1.
4. Optimize the HH-condition varying **sp0**, and **p14**, **phcor1** for maximum signal.

5. Run a variation of **p15** between 20 and 5000 ms on your sample (or tyrosine *HCl). One should see intensity variations (“dipolar oscillations”) which are normally smeared out by extensive homonuclear dipolar couplings between protons.
6. Set up a 2D data set with the pulse program `lgcphetfq`. The figure below compares spectra taken with and without LG-offset during cp.

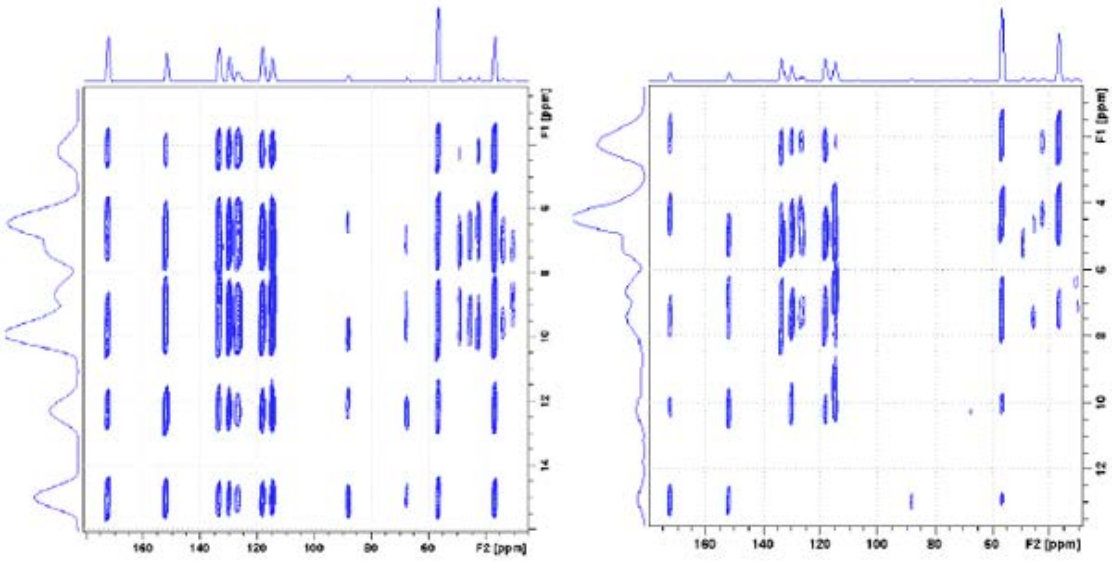


*Figure 9.3: HETCOR on tyrosine *HCl without (left) and with LG contact (1msec contact)*

10 RFDR

Radio Frequency-Driven Recoupling (RFDR) with longitudinal magnetization exchange is a homonuclear dipolar recoupling experiment. This easy setup technique is a zero-quantum recoupling sequence that achieves chemical shift correlation under MAS conditions. The time-dependence of the cross peak amplitudes can be employed to determine inter-nuclear distances. With short dipolar recoupling times, only spins in close spatial proximity lead to cross peak facilitating assignment of ^{13}C resonances in uniformly labelled peptides for instance. RFDR may also be used in order to correlate chemical shifts and crystallographic sites on materials samples.

The homonuclear dipolar recoupling is implemented via the application of rotor-synchronised 180-degree pulses (one inversion pulse per rotor period). The phases of the 180-degree pulses are cycled with Gullion's compensated XY-8 echo sequence in order to achieve efficient recovery of single spin magnetization and to generate an effective dipolar recoupling Hamiltonian during the mixing period. The critical experimental point is to avoid ^{1}H-X recoupling induced by interference between the ^{1}H decoupling RF field and ^{13}C RF recoupling field. This effect can be removed using a ^{1}H decoupling RF field 3 times as strong as the ^{13}C RF field used for recoupling or by using Lee-Goldburg ^{1}H decoupling during the mixing period.

References

T. Gullion, D. B. Baker and M. S. Conradi, *New, compensated Carr-Purcell sequences*, J. Magn. Reson. 89, 479-484 (1990).

A. E. Bennett, J. H. Ok, R. G. Griffin and S. Vega, *Chemical shift correlation spectroscopy in rotating solids: Radio frequency-driven dipolar recoupling and longitudinal exchange*, J. Chem. Phys. 96, 8624-8627 (1992).

A. E. Bennett, C. M. Rienstra, J. M. Griffith's, W. Zhen, P. T. Lansbury and R. G. Griffin, *Homonuclear radio frequency-driven recoupling in rotating solids*, J. Chem. Phys. 108, 9463-9479 (1998).

B. Heise, J. Leppert, O. Ohlenschläger, M. Görlach and R. Ramachandran, *Chemical shift correlation via RFDR: elimination of resonance offset effects*, J. Biomol. NMR 24, 237-243 (2002).

10.1 Experiment

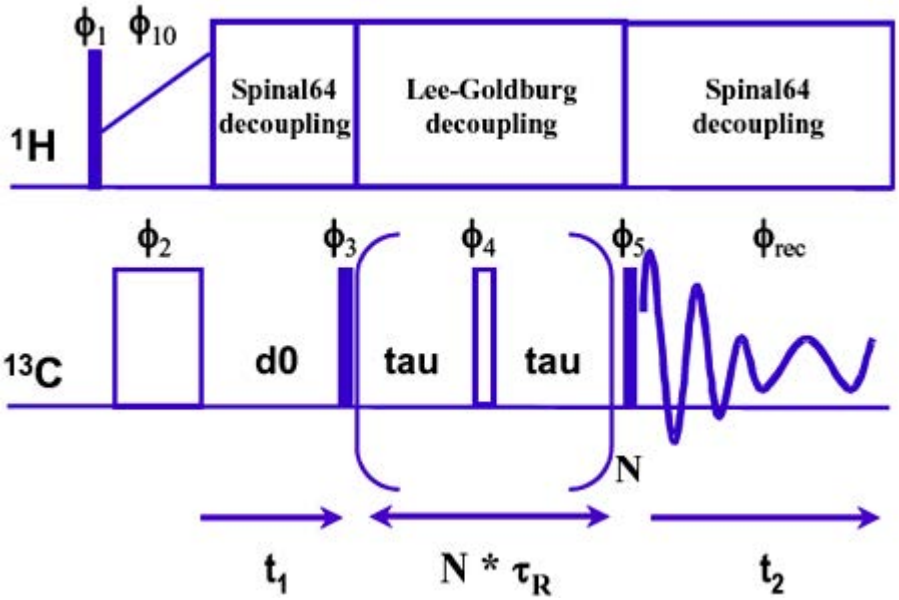


Figure 10.1: RFDR Pulse Sequence for 2D CPMAS Exchange Experiment

φ1 = 1	φ4 = 0101 1010	
φ2 = 1	φ5 = 0303 0303, 1010 1010	
φ3 = 0123 0321	φ10 = 0	φrec = 0220 0220, 1331 1331

10.2 Set-up

Sample: ^{13}C fully labelled histidine.

Experiment time: Less than 1 hour.

Experiment Setup

First setup the ^{1}H-^{13}C cross polarization and the Hartmann-Hahn match according to the procedures described in *Basic Setup Procedures [49]*.

An important experimental consideration of the RFDR experiment is that the RF field strengths used in the recoupling channel (**pl11**) and the RF field on the ^{1}H decoupling channel must be sufficiently different, ca. a factor of 3, to avoid rapid depolarization of the carbon signal during the mixing time. This is usually not achievable, so it should be set as high as possible, using a LG offset.

During the mixing period (**cpds3=cwlg**), as shown in the figure in the section *Experiment [124]*, the Gullion compensated echo sequence used for the mixing period is a XY-8 phase cycling (f4=XYXY YXYX). Consequently, the number of rotor periods for the mixing time (**L1**) must be a multiple of 8.

10.3 Data Acquisition

Sample: ^{13}C fully labelled histidine.

Experiment time: Several hours.

10.3.1 Set-up 2D Experiment

After 1D parameter optimization as described previously, type **iexpno** to create a new data file and switch to the 2D mode using the **123** button. Set the appropriate **FnMode** parameter in **eda**. Pulse program parameters are indicated below (the figure in the section *Experiment [124]* shows the pulse sequence).

Figure 10.2: The 123 Icon in the Menu Bar of the Data Windows Acquisition Parameter Page

The **123** icon in the menu bar of the data windows acquisition parameter page (figure above) is used to toggle to the different data acquisition modes, 1D, 2D and 3D if so desired.

Parameter	Value	Comments
pulse program	cprfdr.av	Pulse program.
nuc1	^{13}C	Nucleus on f1 channel.
o1p	100 ppm	^{13}C offset, to be optimized.
nuc2	^{1}H	Nucleus on f2 channel.
o2p	2-3 ppm	^{1}H offset, to be optimized.
pl1		Power level for contact time on f1 channel.
pl11		Power level for f1 recoupling and excitation.
pl2		Power level for contact time on f2 channel.
pl12		Power level decoupling f2 channel and excitation.
pl13		Power level LG decoupling f2 channel.
p1		90° excitation pulse on f1 channel.
p2		180° excitation pulse on f1 channel.
p3		90° excitation pulse on f2 channel.
p15		Contact pulse on f1 and f2 channel.
d0	3µ	t1 evolution period.
d1		Recycle delay.
cpdprg2	Spinal64	Spinal64 decoupling on f2 channel.
cpdprg3	cwlg	cwlg decoupling on f2 channel.
ns	16	Number of scans.
cnst20	≈100000	Proton RF field in Hz to calculate LG parameters.
cnst21	0	f2 channel offset.
cnst24		Additional Lee-Goldburg offset.
cnst31		Spinning speed in Hz.
l1	for 2-40 msec	Number of rotor cycles for mixing time.

F2 direct 13C		Left column.
td	4k	Number of complex points.
sw	200 ppm	Sweep width direct dimension.
F1 indirect 13C		Right column.
td	256	Number of real points.
sw		Rotor synchronized sweep width, or = 2 sw.
in_f1	= dw or rotor period	Synchronized sampling avoids sidebands.
FnMode		STATES or STATES TPPI.

Table 10.1: Acquisition Parameters

10.4 Spectral Processing

Parameter	Value	Comment
F1 acquisition ^{13}C		Left column.
si	4k	Number of points and zero fill.
ph_mod	pk	Phase correction if needed.
bc_mod	quad	DC offset correction.
F2 indirect ^{13}C		Right column.
si	256	Zero fill.
mc2		STATES or STATES TPPI.
ph_mod	pk	Phase correction if needed.
bc_mod	no	Automatic baseline correction.

Table 10.2: Processing Parameters

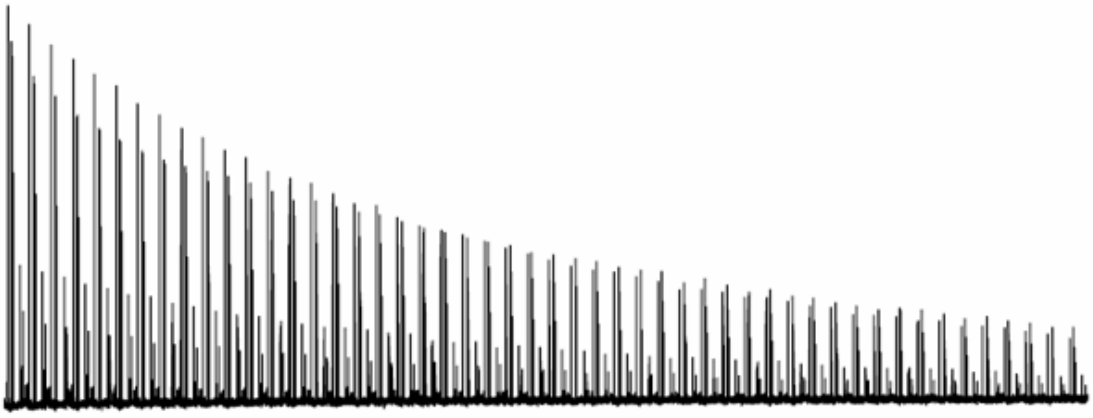


Figure 10.3: 13C Histidine Signal Decay as a Function of the RFDR Mixing Time

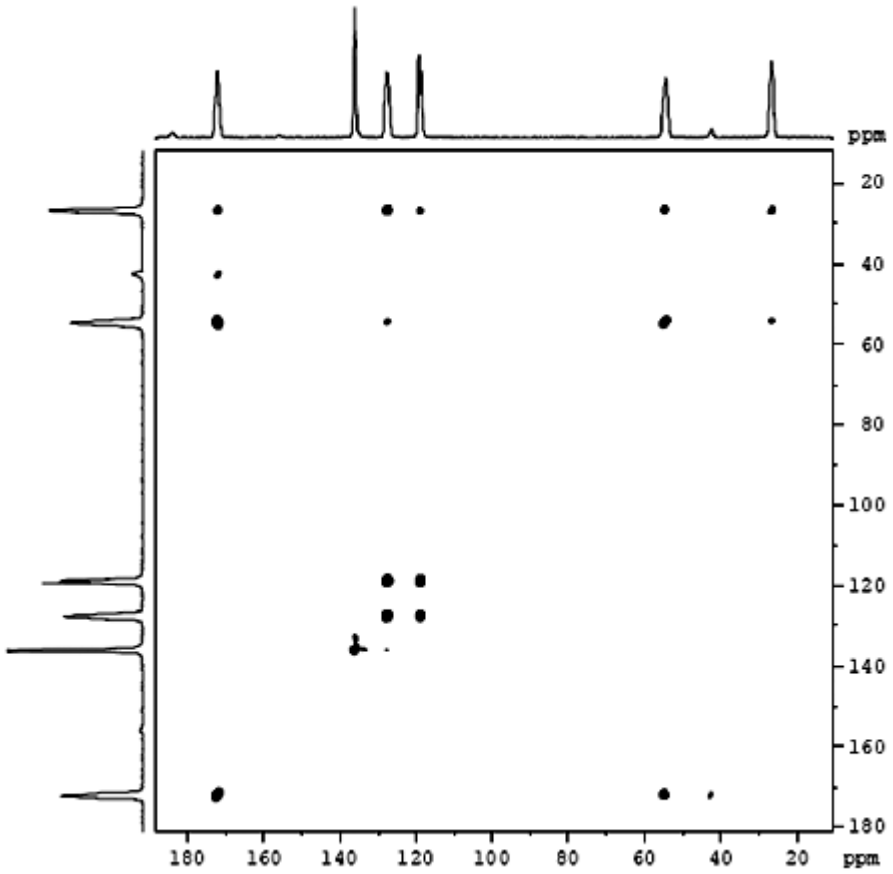


Figure 10.4: 2D RFDR Spectrum of 13C fully Labelled Histidine (RFDR mixing time 1.85 ms)

11 Proton Driven Spin Diffusion (PDSD)

PDSD is a 2D experiment that correlates a spin 1/2 nucleus to another spin of the same species via homonuclear dipolar coupling or chemical exchange. The experiment resembles the NOESY (**N**uclear **O**verhauser **E**ffect **S**pectroscopy**Y**) pulse sequence in the liquid state by replacing the initial 90° pulse with a cross polarization scheme. Since spin diffusion between X-nuclei is measured, cross peak intensities depend on the probability of interaction between different sites, which is low with low natural abundance of NMR-active nuclei. Therefore, these experiments usually require enrichment for nuclei like ^{13}C or ^{15}N in order to allow sensible measurement times.

The pulse program **cpspindiff** allows several types of PDSD experiments to be run. The CP preparation period excites the X nuclei. During the evolution time the X magnetization evolves under the effect of the chemical shift interaction. The evolution time ends with an X 90° pulse that stores the chemical shift information along the z axis and marks the beginning of the mixing time. The X spins communicate through chemical exchange or spin diffusion, depending on the properties of the material and the duration of the mixing time. At the end of the mixing period, another X 90° (read) pulse and the data acquisition follow. The usually strong ^{1}H-X dipolar interaction is removed by high power ^{1}H decoupling during the preparation and the acquisition time. The ^{1}H decoupling is switched off during the mixing to dephase the residual X transverse magnetization. Spin diffusion between X-nuclei is usually very slow and requires very long mixing times since the dipolar coupling between all X-nuclei is small. However, turning the proton decoupling field off during the mix period allows another process to take place: spin diffusion via the proton spin system. Since the rare nuclei are strongly coupled to protons and all protons are strongly coupled to each other, the flip flop transition rate is high along the X_1-H_1-H_2-X_2-pathway. In fact, the spin exchange is almost solely due to proton mediated mechanisms except when chemical exchange is present. At high spin rates, spin diffusion may however still be slow since the X-H spins are decoupled. A simple procedure to recouple the X-H interaction is to irradiate the protons at an RF field of n times the spin rate. These modified sequences are DARR (Dipolar Assisted Rotational Resonance, T. Terao et al.) or RAD (Rf Assisted Diffusion, see C.R. Morcombe et al).

The setup for all these sequences is rather robust, requiring only the ^{1}H to X Hartmann-Hahn condition and the X 90° hard pulse to be set. For RAD and DARR, it is usually sufficient to calculate an RF power level corresponding to n times the spin rate, which is then applied during the mixing period. Rotor synchronization of the mixing period is recommended in some cases, where cross peaks due to sidebands need to be suppressed (de Jong et al.) or where spin diffusion is enhanced by matching the spin rate with the chemical shift difference between the sites to be correlated (M. Ernst et al.). PDSD is typically applied to high abundance nuclei or labeled materials to detect through space proximity between spins. This experiment has been often used on proteins as an alternative to Radio Frequency DRiven spin diffusion (see *RFDR [▶ 123]*). RFDR provides similar information to PDSD but with a different mixing period. Here the term *frequency driven* relates to recoupling pulses on the X channel, whereas in DARR or RAD the radio frequency that drives the recoupling is the proton RF.

An important aspect of this experiment is that the mixing time is a simple delay and no pulse, or only weak RF irradiation (DARR, RAD) is required. Therefore it can be made very long because no technical or experimental problems can arise. So the effects even of small dipolar couplings (requiring long spin diffusion times) can be observed. However the information from this experiment may be ambiguous, because a rather non-selective transfer (within the proton spin system) is utilized.

Nevertheless, even complex molecules like proteins can be surprisingly well characterized by PDSD experiments with different mixing times. The buildup of cross peak intensities can be studied and correlated, for instance, to the structure of a macromolecule in the solid state. The same approach has been used to compare different states of a protein, i.e. bound to a membrane or free, as can be found in the recent literature on solid state NMR applied to protein studies.

More elaborate derivatives of PDSD are also known in bio-molecular NMR, where the unspecific spin diffusion within the proton spin system is filtered through a double quantum selection (Lange et al.).

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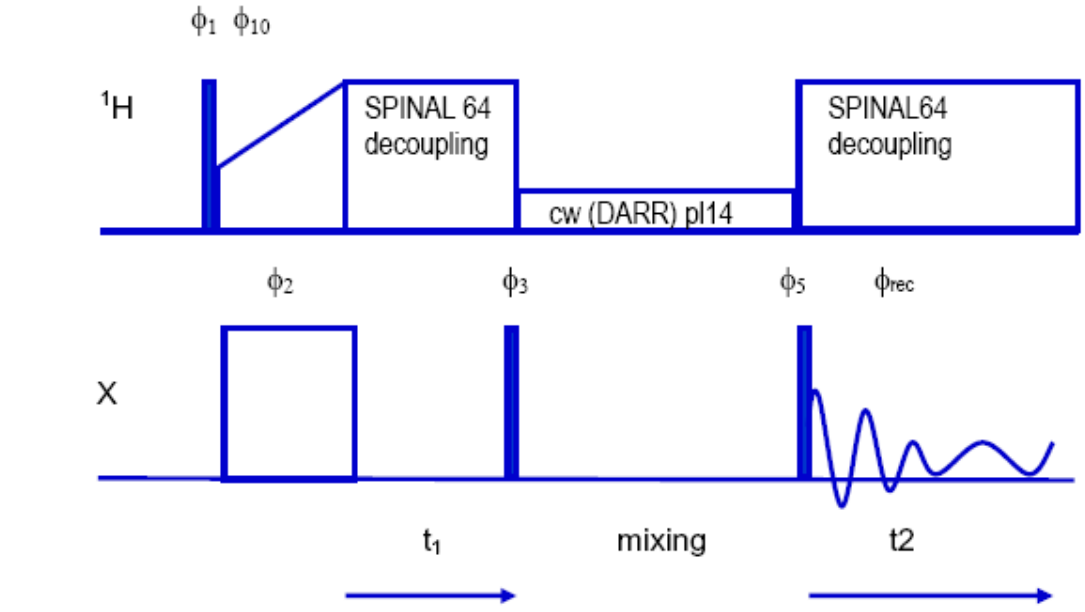
11.1 Pulse Sequence Diagram



ϕ_1 = 1, 3	ϕ_2 = 1	ϕ_3= 0 0 0 0 0 0 0 0 2 2 2 2 2 2 2 2
ϕ_5= 0 0 2 2 1 1 3 3	ϕ_{10} = 0	ϕ_{rec}= 0 2 2 0 1 3 3 1 2 0 0 2 3 1 1 3

Figure 11.1: CPSPINDIFF Pulse Sequence

11.2 Basic Setup

1. On a standard sample (i.e. glycine) determine HH match and decoupling parameters.
2. Check the X 90° hard pulse with **cp90** on the standard sample.
3. If the real sample tunes and matches very differently than the setup sample, verify the HH conditions briefly and eventually the X 90° hard pulse.
4. Set the spin rate as high as possible, making sure to avoid rotational resonance conditions (overlap between center bands and sidebands), then recheck the HH condition. Set **cnst31**=spin rate.
5. Optimize the contact time, o1 and o2 on a ^{13}C 1D CP experiment if necessary.
6. Create a new experiment with either **iexpno** or **edc**.
7. Change to a 2D data set.

11.2.1 2D Experiment Setup

1. Type **iexpno** to create a new data file and switch to the 2D mode using the **123** button. Load the pulse program **cpspindiff**.
2. Recheck the pulse widths and power levels, using **ased**.
3. Go into **eda** and set parameters for sampling in the indirect dimension, the spectral width 1 **swh**. Note, that in TopSpin 2.1 or later the parameter **IN_F1** replaces the parameters **in0** and **nd0**. Usually, 1 **swh** equals **swh**. Choose a suitable spin rate such that no RR condition occurs and sidebands do not overlap with peaks, if possible set the sweep width in F1 **1 swh** equal to the spin rate or such that sidebands folding in along F1 do not interfere.

4. Make sure the correct nucleus is selected in the F1 dimension; make sure to choose an appropriate quadrature detection mode in **FnMode** (usually STATES-TPPI).
5. Choose the appropriate sampling time (**td1**) so that the required resolution (**FIDRES**) in the indirect dimension is achieved.
6. Set the desired mixing time to **d8**. The required multiple of spin periods (from **cnst31**) is calculated as **l1**, the real mixing time may deviate by fractions of a rotor period. The required mixing time may vary widely depending on the sample properties, from a few milliseconds to hundreds of milliseconds, if long distance correlations in a mobile sample need to be observed. Note that longer mixing times will result in S/N deterioration, as the mixing time approaches the T_1 of the observed nuclei.
7. Set **pl14**, if DARR/RAD is desired, or else make sure **pl14 = 120 dB**. For DARR/RAD calculate the required power level using **calcpowlev**, or use the setup procedure shown in *Basic Setup [▷ 131]*.
8. Start the experiment.

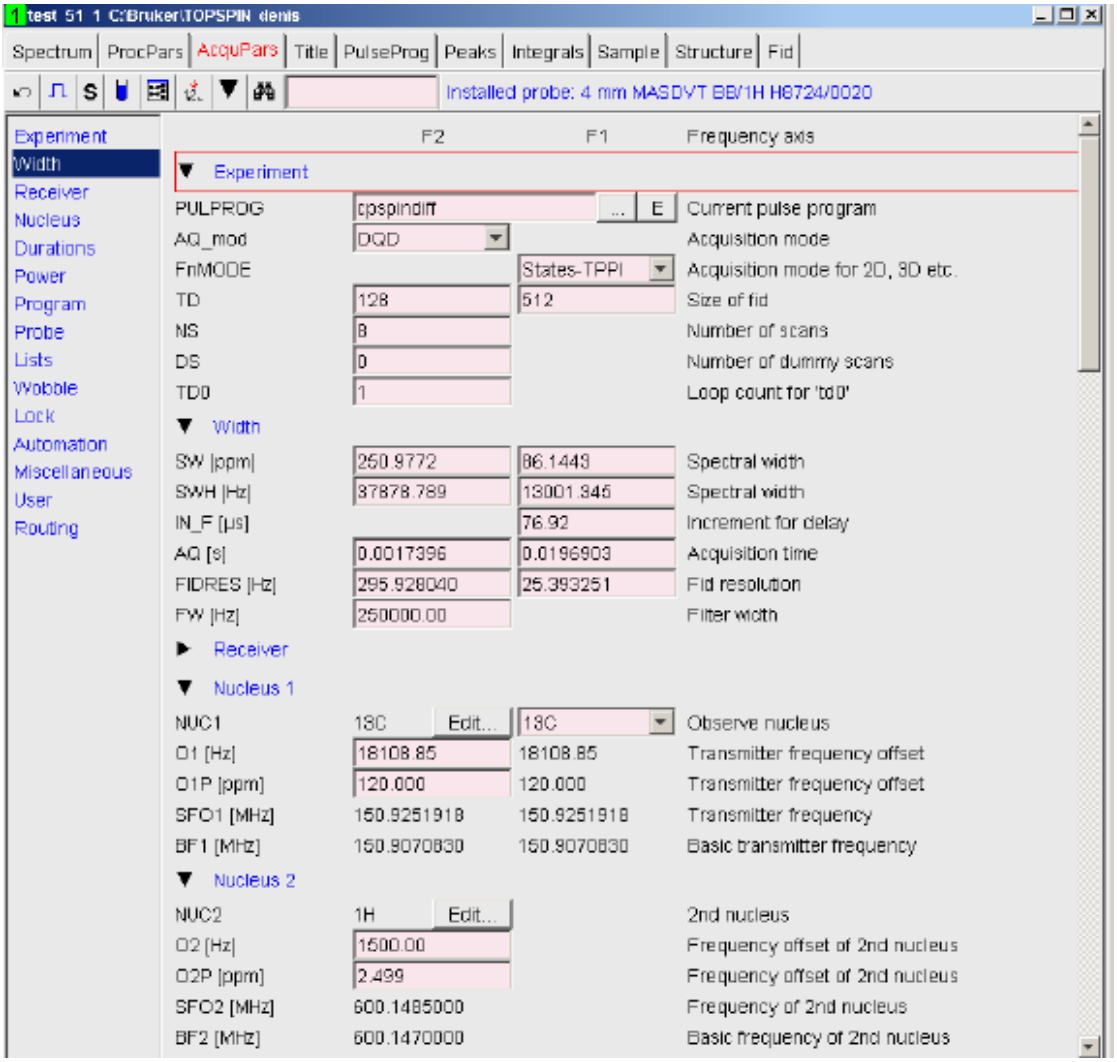


Figure 11.2: The Acquisition Parameter Window eda

11.3 Acquisition Parameters

Sample: ^{13}C labelled histidine (labelled tyrosine-HCl).

Experiment time: 90 min. (20 min.).

Parameter	Value	Comments
PULPROG	cpspindiff (old: cpnoesy)	Pulse program.
NUC1	^{13}C	Set ^{13}C in both F2 and F1 column.
SW	250 ppm	To be optimized.
O1P	120 ppm	To be optimized.
NUC2	^{1}H	For CP/decoupling only.
O2P	3 ppm	To be optimized for dec.
PL1		For ^{13}C contact.
PL11		For ^{13}C flip pulses.
PL12		For ^{1}H excitation and decoupling.
PL14	for n*spin rate (DARR) or 120	Recoupling.
SP0		For ^{1}H contact using shape.
SPNAM0	ramp.100 or ramp70100.100	For ^{1}H-^{13}C contact.
CPDPRG2	SPINAL64	At PL12.
P1		^{13}C excitation (flip) pulse.
P3		^{1}H excitation pulse.
P15		^{13}C-^{1}H Contact pulse.
PCPD2		Decoupling pulse for spinal64.
D1		Relaxation delay.
D8	5-500 msec	Depending on sample.
CNST31	MAS speed	Used to calculate d31 (rotation period).
L1	calculated from cnst31 and d8	Number of rotor cycles for mixing time.
AQ_MOD	DQD	
TD {F1}	512	Number of points.
SW{F1}	usually = SW	=MASR if possible.
NUC1{F1}	=NUC1	
TD {F2}	128	Number of points.
ND0	1	Not required in TopSpin 2.1.
NS	4*n	
FnMode	TPPI/States/States-TPPI	

Table 11.1: Acquisition Parameters

11.3.1 Processing Parameters

Process with **xfb.**

Parameter	Value	Comment
F2 acquisition ^{13}C	**********	Left column.
SI	1k	Number of complex points in direct dimension.
WDW	QSINE	Apodization in t2.
SSB	2-3	
PH_mod	pk	
F1 indirect ^{13}C	**********	Right column.
SI	512	Number of complex points in indirect dimension.
WDW	QSINE	Apodization in t1.
SSB	2-3	
PH_mod	pk	

Table 11.2: Processing Parameters

11.4 Adjust the Rotational Resonance Condition for DARR/RAD

1. Load the Adamantane sample, spin at the same speed as desired for your sample, match and tune, use a suitable **cp** setup (same as in section *Acquisition Parameters [▶ 133]*).
2. Set **CPDPRG2** to **cw.**
3. Use the au program calcpowlev to calculate the power level required for a proton decoupling RF field of n × **masr**, using p3 and pl12 as reference values. Refer to chapter *Basic Setup Procedures [▶ 49]* for more information).
4. Vary the decoupler power level **pl12** used with cw decoupling as indicated in the figure above from a power level value **pl12** 1 dB below the calculated n = 1 condition to 1 dB above the calculated n = 2 condition. Bandwidth considerations favor the n = 2 condition, sample heating considerations favor the n = 1 condition. An RF field of 2 × proton chemical shift range is on the safe side.
5. Enter the power level determined above as **pl14** recoupling power for DARR or RAD.
6. Using DARR or RAD shorter mixing times are possible.

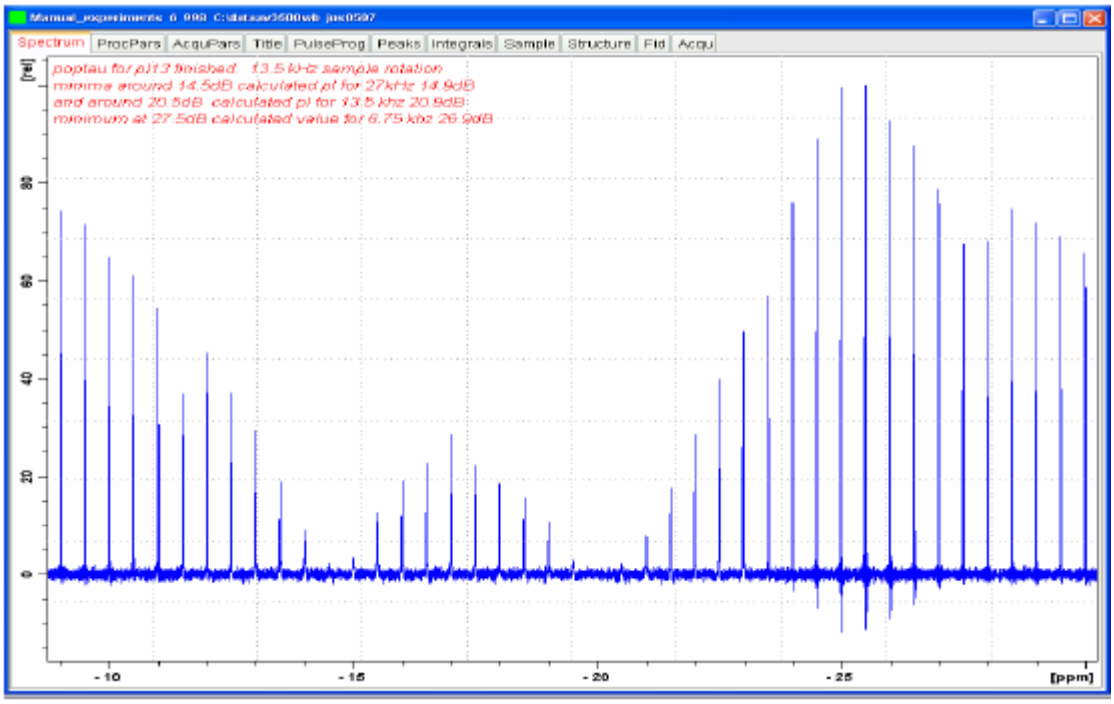


Figure 11.3: POPT Result for the cw Decoupling Power Variation

The figure above shows the POPT result for the cw decoupling power variation from about 50 kHz RF field to about 50 kHz RF field, spinning the adamantane sample at 13 kHz. The minima at 14.5 and 20.5 dB indicate the n = 2 and n = 1 RR conditions (26 and 13 kHz RF field).

11.5 Example Spectra

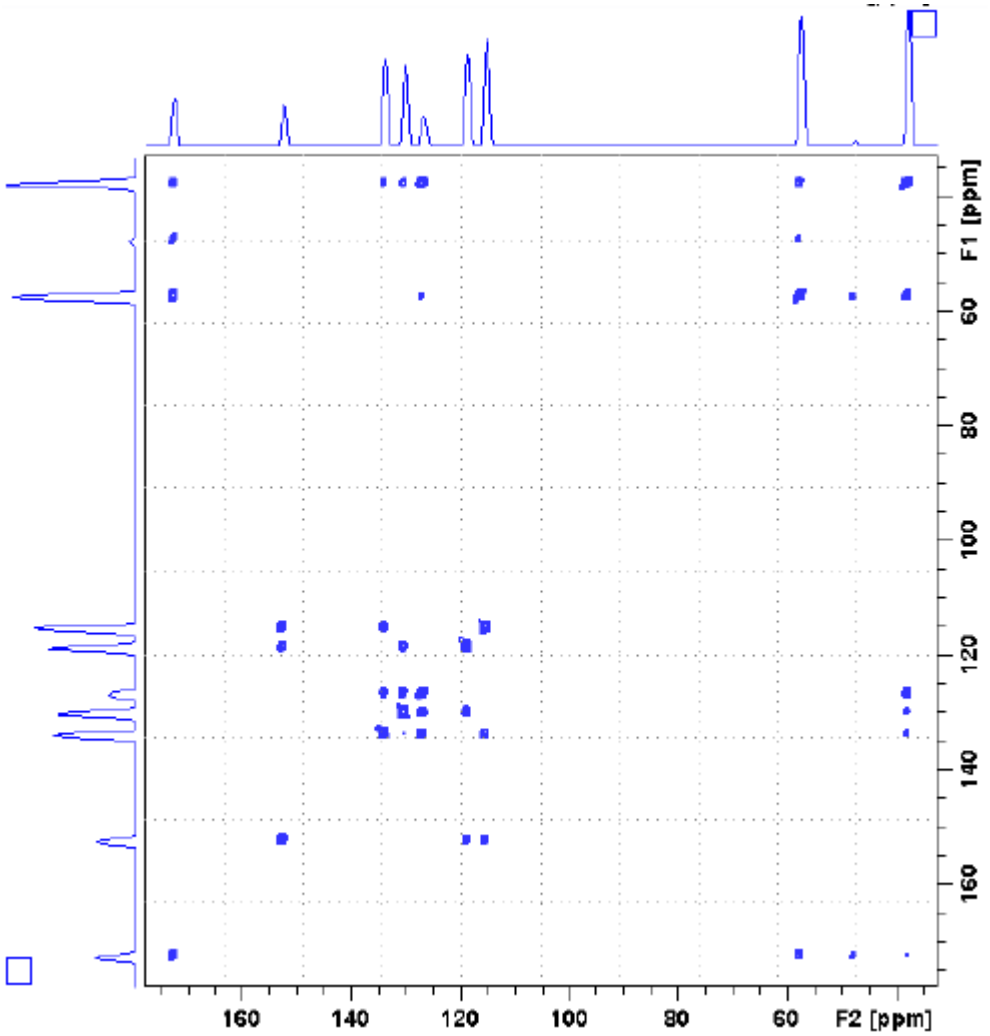


*Figure 11.4: 13C CPSPINDIFF of fully labeled tyrosine*HCl, spinning at 22 kHz, 4.6 msec mix. Upper: PDSD, lower: DARR*

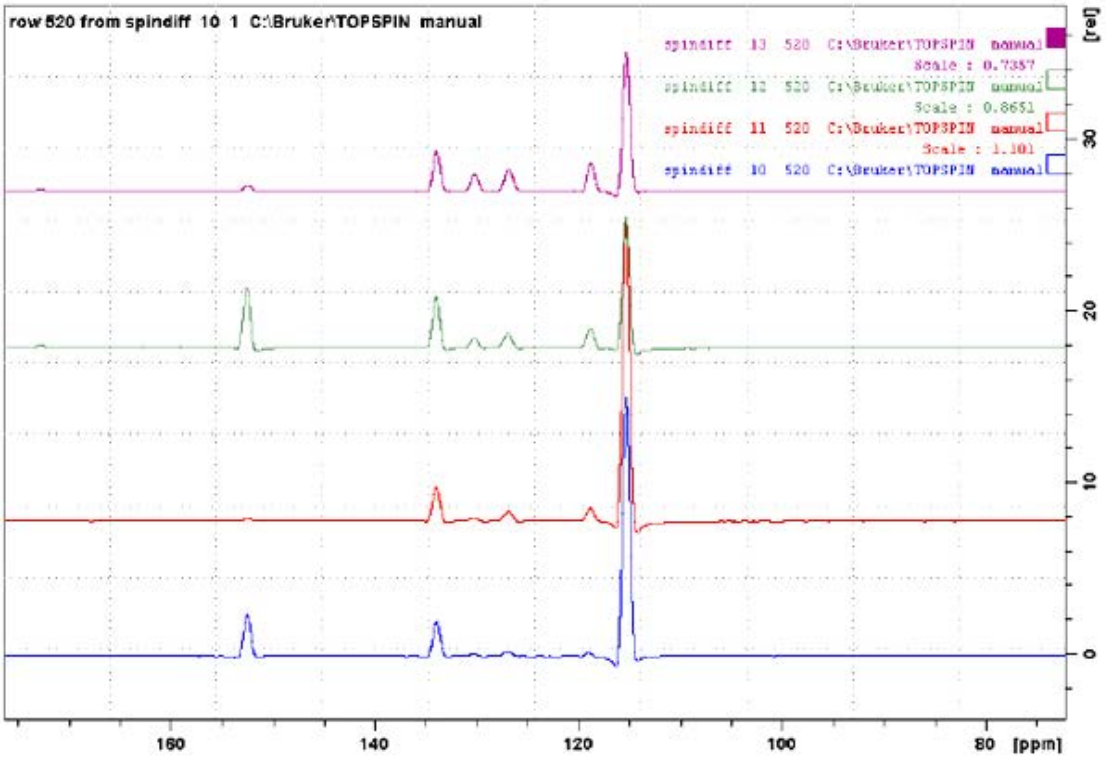


Figure 11.5: Comparison of DARR/PDSD

The figure above is a comparison of DARR/PDSD, with 4.6 and 20 msec mixing time, sample tyrosine-HCl spinning at 22 kHz. Traces through peak at 115 ppm, most high field aromatic carbon. Traces from below: DARR at 4.6 msec mix, PDSD at 4.6 msec mix, DARR at 20 msec mix, and PDSD at 20 msec mix. Note that some cross peak intensities differ substantially!

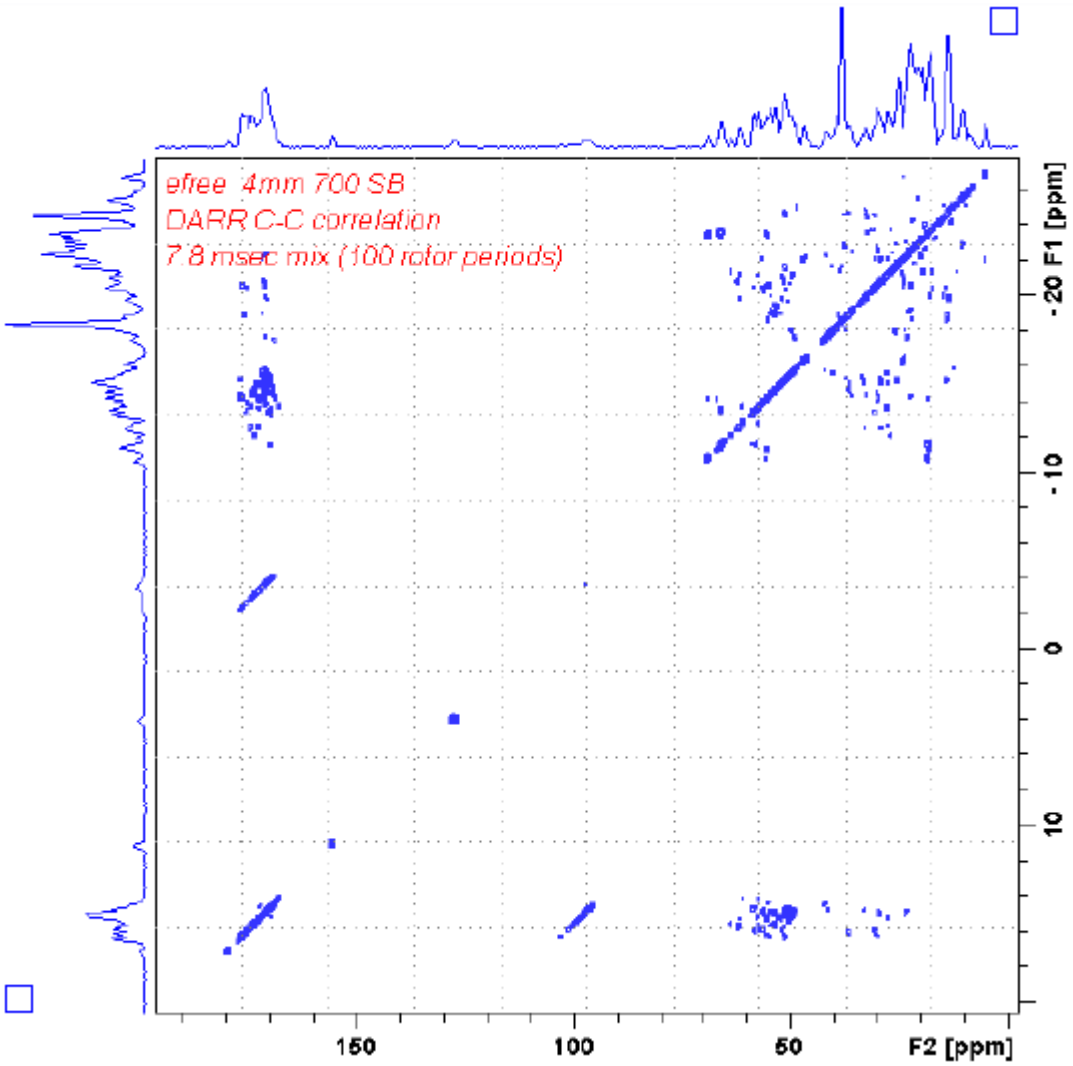


Figure 11.6: 13C DARR of Fully Labelled Ubiquitine Spinning at 13 kHz

12 REDOR

Rotational **E**cho **DO**uble **R**esonance is an experiment based on the heteronuclear dipolar coupling between the observed nuclei. The REDOR sequence investigates this coupling under high resolution MAS conditions.

Dipolar couplings between I (decoupled spin) and S (observed) spins are spun out under MAS if there are no strong homonuclear interactions and if the heteronuclear coupling is not too big. In the case of couplings between most hetero-nuclei (like ^{13}C, ^{15}N and ^{29}Si) this is usually the case (small couplings of a few kHz and dilute spins), if the coupled nucleus is ^{31}P, ^{19}F or even ^{1}H, the coupling may not easily be spun out and the standard REDOR sequence may not be applicable in these cases.

The REDOR sequence reintroduces the heteronuclear dipolar coupling between the spin S and I by applying p pulses every half of a rotor period on the second channel (I), while the S channel is observed. A p pulse at half a rotor period will refocus the dipolar interaction averaged by spinning and dephase the magnetization, leading to an attenuation of the observed signal. Evaluation requires the acquisition of 2 data sets, one with refocusing pulse, the other one without, so that the natural dephasing can be subtracted out from the dipolar dephasing due to the refocusing pulse. The reference experiment (without I refocusing pulse) and dephased experiment (with I refocusing pulse) are subtracted and evaluated. Reference experiment and dephased experiment can be acquired consecutively or in an interleaved mode so that experimental drifts will not cause large errors. Usually, the experiment is set up as a pseudo-2D experiment where the number of rotor periods with p pulses is increased before detection.

The experiment can either be used to investigate isolated spin systems or multi speed systems. In both cases the time dependent difference of the echo S_0 (without the reintroduction of the heteronuclear dipolar coupling) and the second echo experiment S' (with the reintroducing p pulses applied on the I channel) can be used for calculating the distance information for the two involved spins or the second moment of the spin system respectively.

In isolated two spin systems the measured REDOR (dephasing)-curve can be used to determine distance information between the two involved spins. In case of investigating multispin systems the experimental REDOR dephasing curves can only be used to determine the second moment (M_2).

$$\propto \sum 1/r^6$$

By the relation of the second moment to the distance by, this information in combination with theoretical simulations can be used to determine a mean distance between the involved spins as well.

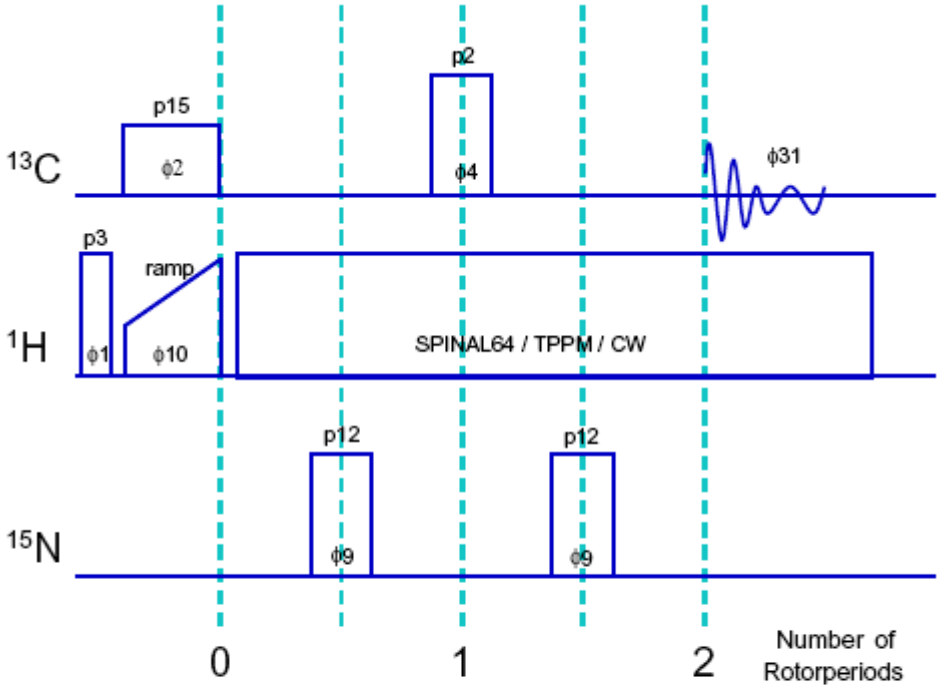
In the case of very strong dipolar couplings the normal REDOR approach for multispin systems can not be used without introducing severe errors into the calculated second moment. With very strong dipolar couplings, the signal intensity may be lost after very few or only one refocusing pulse, so the decay curve cannot be measured, and the reference experiment does not show a decay independent of the heteronuclear dipolar coupling. Faster spinning may solve the problem. Alternatively, the constant time (CT-)REDOR sequence may be used, which can help to enhance the performance in these cases and to reduce the experimental time needed dramatically.

Here the refocusing pulse is stepped in small time increments from the beginning of the rotor period, where it has a small effect on the signal intensity, to the middle of the rotor period, where it has the maximum effect.

If one or both of the involved nuclei do have a quadrupolar moment, the REDOR sequence should not be used, but there are several different REDOR like experiments in literature which have an enhanced performance on observing quadrupolar nuclei. These are, for example, the **R**otational **E**cho **A**diabatic **P**assage **DO**uble **R**esonance (REAPDOR) or **TRA**nsfer of **P**opulation in **DO**uble **R**esonance (TRAPDOR) sequences.

Usually, the setup is chosen such that the more sensitive nucleus is observed. The measured coupling is independent from the choice of the nucleus, but there may be reasons to consider carefully, which nucleus is observed. Of course it is tempting to observe the nucleus with higher isotopic concentration, but this is usually not recommended, since it is more difficult to observe a small intensity change on a strong signal (which might be caused by fluctuations) than a bigger difference on the low abundance nucleus caused by the high abundance of the coupling partner. An example: The measurement will be more precise observing ^{13}C and defocus with ^{31}P pulses than vice versa, because the effect on ^{31}P, caused by 1.1% of ^{13}C would be max. 1.1%, a very small change which would require extremely high S/N and extremely high stability in signal generation and spin rate.

12.1 Pulse Sequence



ϕ_1=1 3	ϕ_9=0 1 0 1 1 0 1 0
ϕ_2=0 0 1 1 2 2 3 3	ϕ_{10}= 0
ϕ_4=0 0 1 1 0 0 1 1	ϕ_{31}=0 2 1 3 2 0 3 1

Figure 12.1: REDOR Pulse Sequence

12.2 Setup

The example of setup given here is based on a biological sample, but of course the procedure will not change if you are going to analyze different samples with different combinations of coupled spins.

Sample: fully labelled $^{15}N/^{13}Ca$ Glycine (diluted in natural abundance glycine to 10%). Dilution will reduce long range dipolar interactions strongly and lead to a well defined direct interaction between ^{15}N and ^{13}Ca, so that a single frequency dipolar modulation is obtained. A triple resonance probe with the X channel tuned to ^{13}C and the Y channel tuned to ^{15}N is required. Set ^{13}C observe with cross polarization from protons. It is recommended that separate preamps are used for C and N so observation can be changed between C and N without rewiring. There should be an X-low pass filter or ^{13}C bandpass filter on ^{13}C, a ^{15}N low pass or bandpass on ^{15}N.

MAS rate: 5-10 kHz. The MAS spinning speed should be stable within 1 to 2 Hz in order to get a well refocused echo.

Overall Experimental time, including setup procedure: 3-5 hours.

Packing the sample is critical for the success of the experiment, check that your sample is within the central region of the spinner or use a 12 ml spinner. The quality of the refocusing p pulses is essential, this can only be achieved with a center packed sample. The coil of a 4 mm MAS probe has a length of 10 mm, the sample should be no longer than 5 mm, preferably 3 mm (CRAMPS spinner or 12 ml HR-MAS spinner).

Setup the CP conditions for the ^{1}H magnetization transfer on both coupled nuclei (X and Y). Use cp90 to determine precise p pulses on both the X and Y channel of your probe.

Accurately setup the p/2 and p pulses for both 13C and 15N according to the standard setup procedures with an accuracy of at least 0.1ms.

After the setup, the pulse lengths for the different channels should be within the same duration and short enough to not exceed an overall duty cycle of about 5%.

One can reoptimise the refocusing pulse on the coupled nucleus using a 1D version of redor, setting the number of experiments (1 td) to 1. Here, the number of refocusing pulses **l0** value is selected so that there is a noticeable decrease in the dephasing experiment (normally this is between **l0**=1-15), and then the refocusing p pulse is optimized for minimum signal intensity in the dephasing experiment.

Parameter	Value	Comments
pulse program	cpredori	pulse program
nuc1	^{13}C	nucleus for f1 channel
nuc2	^{1}H	nucleus for f2 channel
nuc3	^{15}N	Nucleus for f3 channel
p3	according to specs	π/2 pulse on f2 channel
p15	2000	Contact time between f1 and f2
pcpd2	about 2*p3	pulse length for decoupling sequence
p2	according to specs, 6-10 µsec	π pulse on f1
p12	about 10 µsec	π pulse on f3
cnst31	=masr	MAS spinning rate
l0	1	starting value for the desired evolution time, value must be odd

d1	4s	recycle delay
pl1	for HH condition	f1 power level for contact pulse
pl11	according to specs	f1 power level for π pulse
sp0	for HH condition	power level for 1H ramp
pl2	-	not used
pl12	adequate	power level 1H decoupling
pl3	adequate	power level for f3 π pulse
pl22	120	power level in S0 experiment for the recoupling pulses
spnam0	ramp	ramp file name for CP
td(f1)	32-256	depending on coupling
aq	20-40 msec	acquisition time in f2
ns	8	number of scans per experiment
fnmode	QF	phase correction mode in the f1 dimension
rg	16-64	receiver gain level
digmod	baseopt or digital	digitizing mode

Table 12.1: Acquisition Parameters for a 13C observed C/N REDOR

12.2.1 Data Acquisition

Setup of the 2D data set:

After the optimization, as described above, type **iexpno** to create a new data set, afterwards switch to 2D data mode by using the **123** button. Set the time domain for the F1-Dimension according to the maximum desired evolution time of the final REDOR curve. To calculate the value for td1 you have to keep in mind that the REDOR–Program is organized in multiples of two rotation periods (compare with the pulse program scheme). For example, for a given MAS rate of 10 kHz and a desired overall evolution time of 10 ms, you have to set **td1** to a value of 200. This will record 100 sets of S_0 and S' experiments with an maximum evolution time in the last two data rows of the desired 10 ms. The pulse program library supplies the following REDOR sequences:

cpredor: Standard REDOR for dipolar couplings > 500 Hz, incrementing in 2 rotor period intervals. Two data sets are required with pl3 set for pulsing (REDOR experiment) or no pulse (=120 dB, reference experiment).

cpredori: Stores these data sets in the 2D data frame by interleaving scans of the S_0 and S' experiments for the same evolution period (see the figure in *Data Processing [▶ 143]*).

cpredorxy8: Increments in units of 16 rotor periods, for small couplings. Here the XY-8 scheme is used for recoupling pulses which is insensitive to offsets. While in cpredor and cpredori the p pulses should be executed close to resonance, this is not required nor desired in the XY-8 version because the offset dependence is well compensated.

In addition to the cross-polarization pulse programs, all REDOR sequences also exist as a direct excitation version without a cp-Step. These are the programs **redor**, **redori** and **redorxy8**. The setup procedure of these sequences is identical to the ones explained earlier, but you can skip the CP optimization procedure, which is replaced by a p/2 pulse on the observe channel.

12.2.2 Data Processing

The following refers to the sequence **cpredori**, where the REDOR-experiment and the reference experiment are executed in an interlaced mode (which is less likely to be subject to systematic errors).

The acquired NMR data are arranged in a 2D like structure, every odd row (1, 3, 5,…) contains the REDOR data set with additional p pulses on channel I (S' experiment), every even row (2, 4, 6,…) contains the corresponding echo experiment (S_0 experiment) with the same evolution period t, after the RAW-Data is processed with the **XF2** command (see the figure below). For further processing you can then either use the function **T1/T2 Relaxation** in the analysis part of TopSpin in order to do the integration and/or find the peak maxima of the S_0 and S' intensities automatically or use your favorite deconvolution program for data analysis.

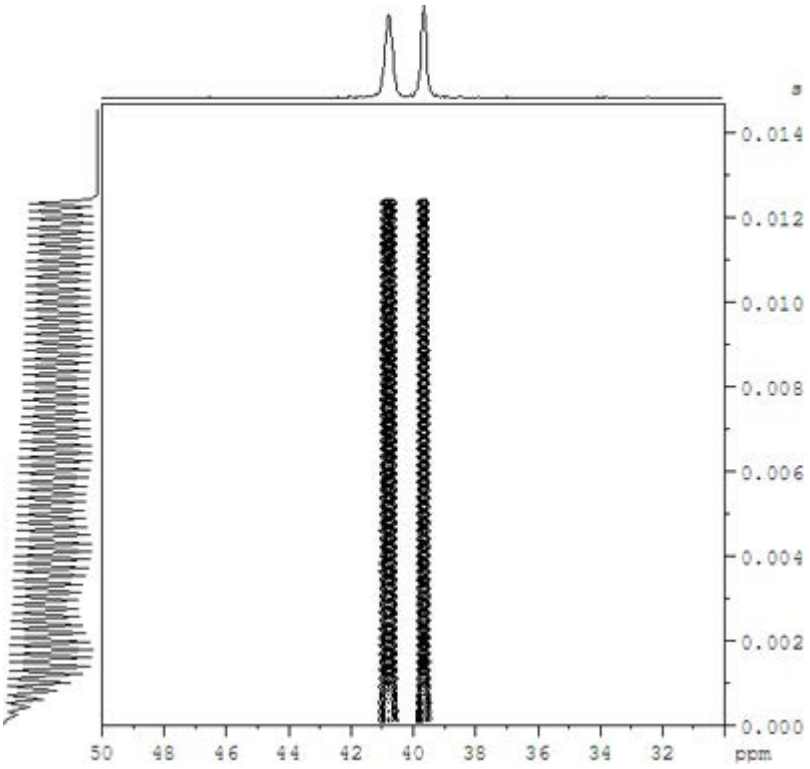


Figure 12.2: 2D data set after "xf2" processing

In the figure above the data set contains the alternating S' and S_0 experiments.

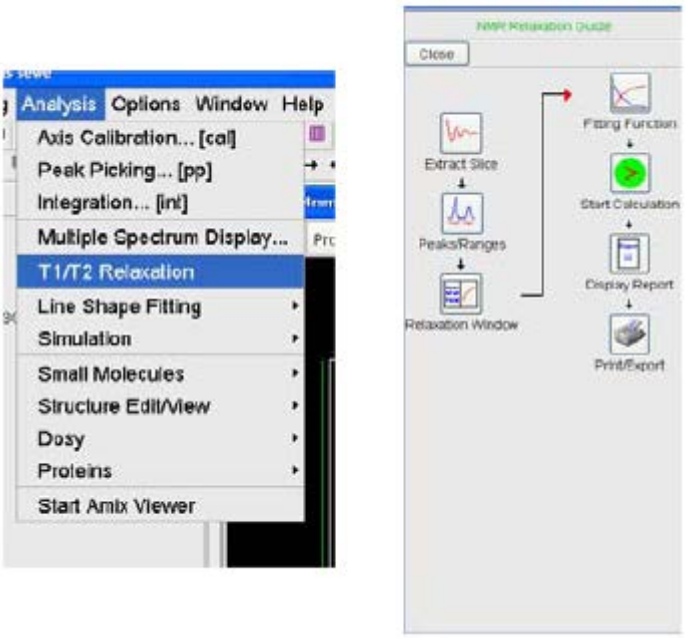


Figure 12.3: T1/T2 Relaxation for further Analysis of the Data Figure and the Analysis Interface

If you are going to use TopSpin you have to choose **T1/T2-Relaxation** in the Analysis Menu (see the figure above) to open the graphical interface for the data analysis.

To begin extracting the first spectrum, use the **extract slice** button and select the desired peak by manual peak picking. To save the data press the **Save** button:

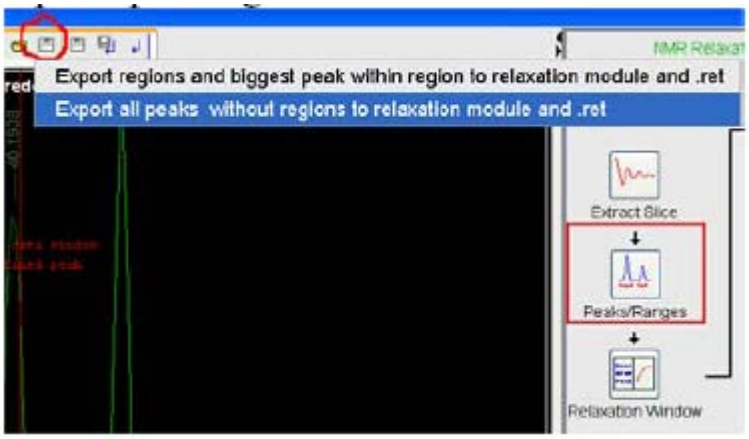


Figure 12.4: Saving Data to Continue to the Relaxation Window

Now after switching to the relaxation window, TopSpin will show the parameter window which can also be accessed later by clicking the marked button in the upper icon bar (see figure below). Here the value for the number of points has to be set to the td2 value and the list file name has to be switched to auto, otherwise the data preparation will fail (compare with the figure below). By clicking **OK**, TopSpin will automatically pick all the intensities for your measured REDOR experiments. These values are then saved in the processed data directory (\data\user\nmr\experiment\exp#\pdata\1) in the file “t1t2.dx”. The intensities are saved in two columns for each peak, while the first column represents a arbitrary x-scale, the measured intensities are within the second column. Remember, because of the used pulse program `cpredori` every odd line contains an intensity value for an REDOR spectrum (S’) and every even line the corresponding intensity of the ECHO experiment (S_0).

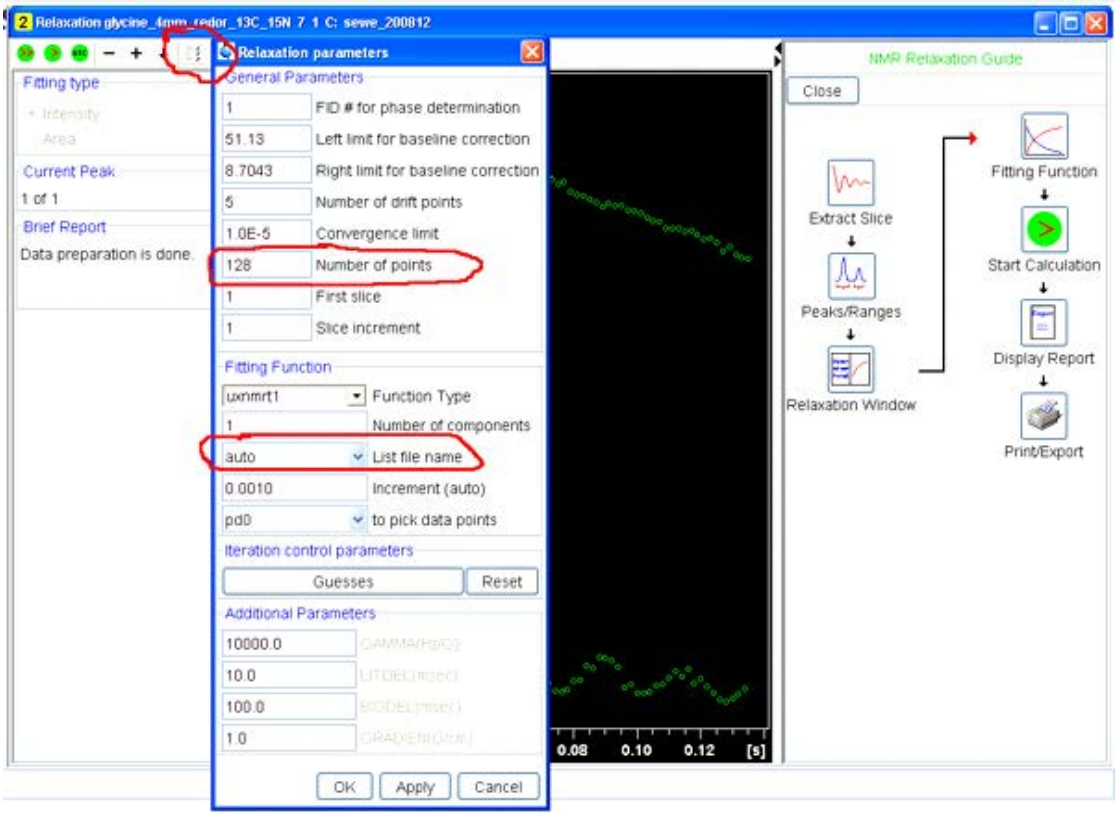


Figure 12.5: Setting the Correct Analysis Parameter

After importing this file into Excel or any other program (using either “Origin” or “Igor” is recommended) for calculating the values for $(S_0\text{-}S')/S_0$, these normalized intensities are plotted versus the evolution time of each spectrum (which is a multiple of “TR*2”) as it is shown in the figure below. The x-axis is therefore calculated by

$$(1/\mathrm{MASrate}) \cdot (2i-2)$$

Where **i** is the number of the corresponding data point $(S_0\text{-}S')/S_0$ (the number of the S_0/S' set within the 2D data set). In this experiment MAS spinning speed was 10 kHz resulting in data points every 200 ms. Note, only the analysis of the low field shifted peak is shown here, corresponding to the alpha-glycine modification of the measured sample.

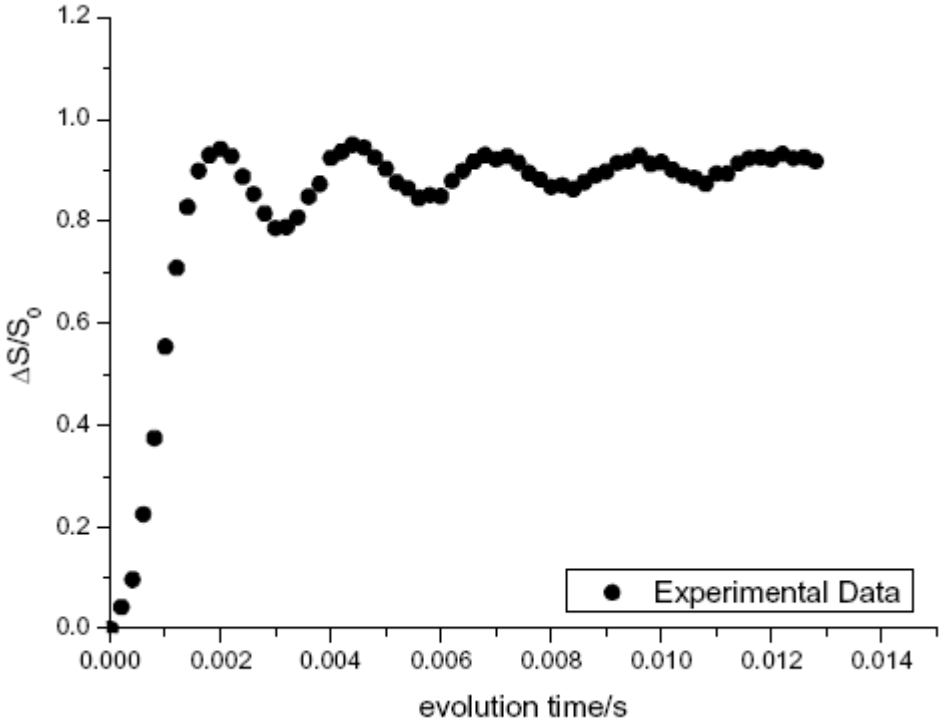


Figure 12.6: Plot of the Normalized Signal Intensity Versus the Evolution Time

There are different methods for the interpretation of the experimentally measured REDOR curves. In the case of isolated two spin systems, like in this case of ^{15}N-^{13}Ca-glycine, it is generally possible to fit the experimental dephasing curve by using a combination of bessel functions. This is called the “REDOR transformation” and gives you direct access to the dipolar coupling information for the measured spin system (for details check reference 6 in *Final Remarks [▷ 150]*).

The more common way for the interpretation of the experiment is the second moment approach, which is also suitable for multiple spin systems. Here the beginning of the REDOR curve can be fitted by a parabolic approximation up to normalized signal intensities of about 0.2-0.3 (for details see e.g. reference 10 in *Final Remarks [▷ 150]*). In the case of very strong dipolar couplings this approach may be restricted to very high MAS spinning speeds, because otherwise it will not be possible to get enough data points within the 0.3 area of the curve (it here may be useful to use a more efficient REDOR technique for strong dipolar coupled systems, like CT-REDOR, see e.g. reference 7 in *Final Remarks [▷ 150]*).

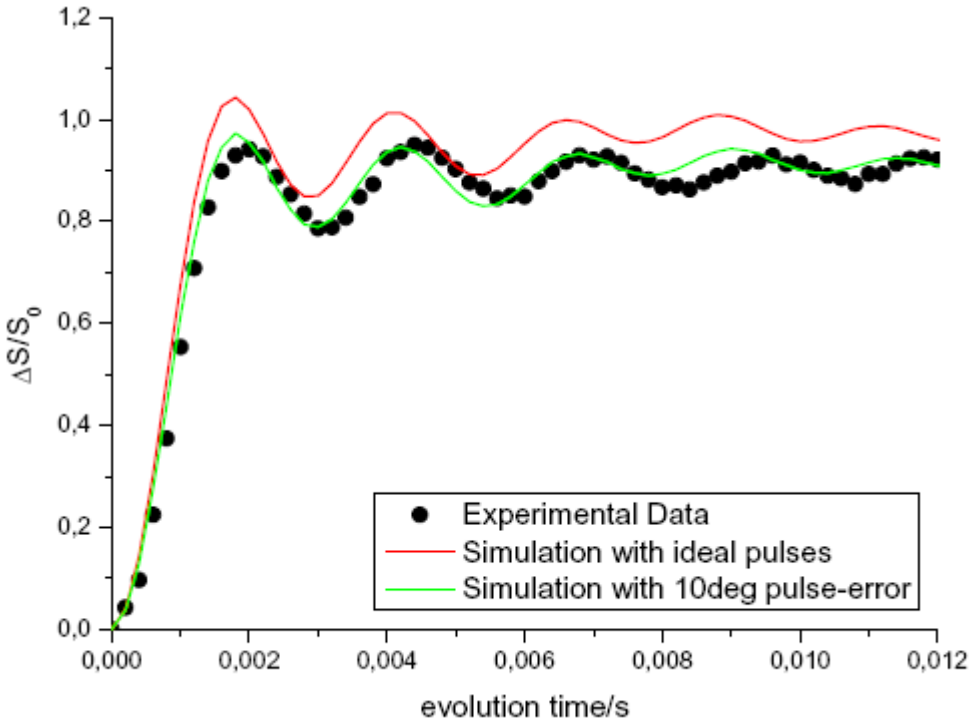


Figure 12.7: Experimental data for the glycine 13C{15N}-REDOR

Red and green curves are the results of different Simpson simulations.

Finally, the complete REDOR curve can be simulated using the Simpson (reference 8 in *Final Remarks [150]*) NMR simulation package. The following pages will explain the M_2 approach as well as the Simpson interpretation of the glycine REDOR data. The figure above shows the experimental data points together with two different Simpson simulations (for details of the geometry and distance information of the labelled ^{15}N-^{13}Ca spin pair of glycine see reference 9 in *Final Remarks [150]*). The red simulation shows the time dependent evolution assuming ideal p pulse lengths on both the S and I channel (corresponding to an experiment without any errors on both frequency channels for ^{15}N and ^{13}C), leading to a slightly too high theoretical REDOR curve compared to the actual experiment. The green curve shows the same simulation assuming pulse errors of 10% on both channels, corresponding very well with the experimental data.

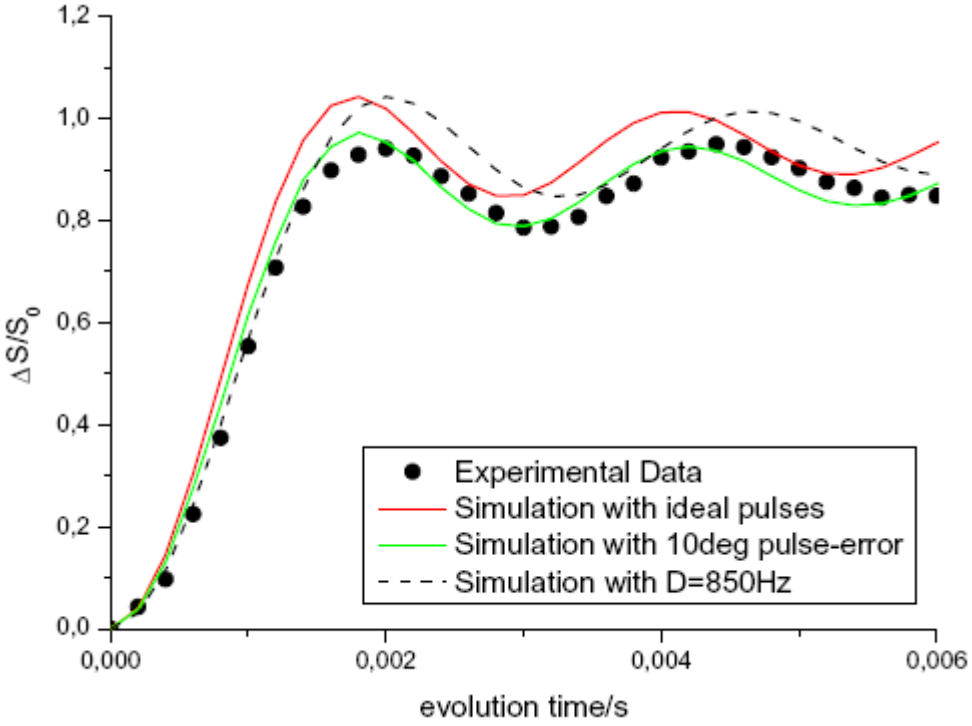


Figure 12.8: Comparison of Experimental Data to a Simulation with Reduced Dipolar Coupling

The figure above shows a zoomed view of the beginning REDOR curve for the glycine sample. The third, simulated line (broken black line) shows a Simpson simulation with ideal pulses on both channels, but varying the dipolar coupling between I and S in order to fit to the experiment. The simulation using a dipolar coupling constant of 850 Hz fits the experimental data points very well. This coupling can be transformed into a distance between ^{15}N and ^{13}C of about 1.53 Å, which is compared to the theoretical value of 1.47 Å (964 Hz) an error within 10%. As you can see, in an unknown spin system, the interpretation by using Simpson will always suffer from the fact that an un-optimal setup of the experiment will introduce the same error like a reduced dipolar coupling between the analyzed spins; these two effects cannot be easily separated from each other during the interpretation process. Therefore it is useful to perform a calibration run like the glycine measurement before analyzing an unknown sample and to set up the complete experiment very carefully in order to calibrate the experiments.

The above shown simulated REDOR curves can now be used to demonstrate the second moment approach in analyzing REDOR experiments. Here a parabolic fit is used to describe the first few points of each curve. In the case of the isolated two spin system of the glycine this parabola is defined by:

$$\Delta S / S_0 = 4 / 3\pi^2 (NT_r)^2 M_2$$

Using this equation you end up with the values given in for the distances and the second moments. The M2 values given in brackets for the simulated curves are calculated using the second moment approach, in order to demonstrate the error margins you have to expect by the M2 approach.

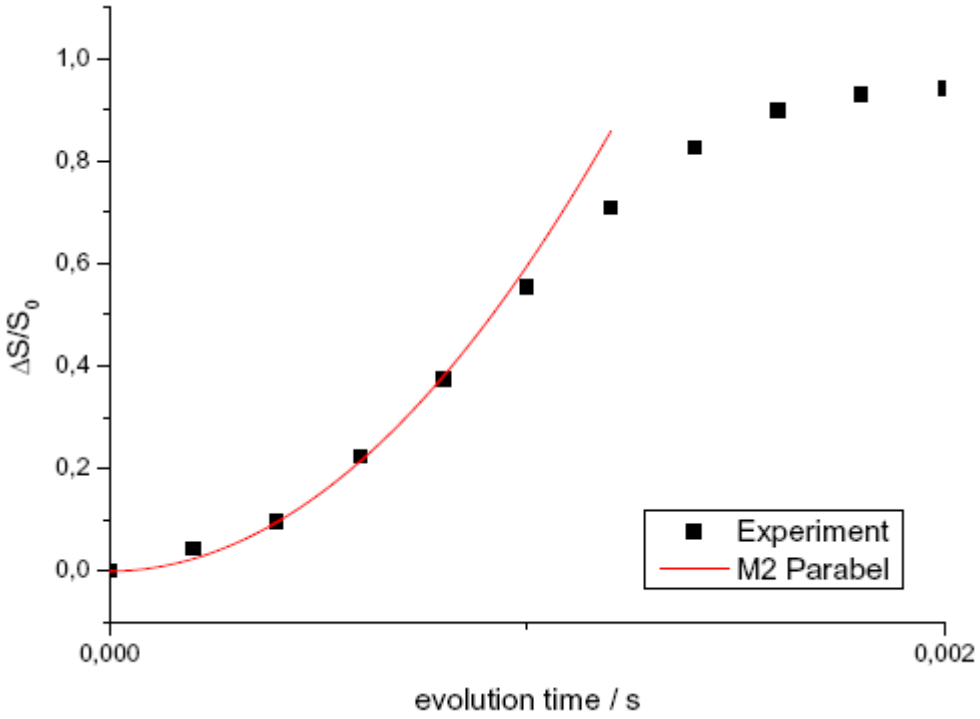


Figure 12.9: Experimental data with the corresponding M2 parabolic analysis

The figure above shows the calculated parabola for the experimental data set. As you can see in the corresponding table the experimental setup reflects the theoretical M2 within an error margin of 40% while this transforms to an overall distance error of about 10%.

Afterwards the calculated M2 can be transformed into the dipolar coupling constant by using:

$$\sqrt{15/(16\cdot\pi)^2/0.5/(0.5+1)\cdot M_2}$$

The comparison of the second moment directly calculated from the experimental data with the second moment extracted from the best fitting Simpson simulation, gives you in this example 747 Hz compared to 850 Hz. Of course this result can be improved by introducing more data points within the interpretation region of the parabolic fit, which can easily be done by running the experiment again with higher spinning speeds. Afterwards the data sets can be combined before running the interpretation process.

Measurement	M2 [s-2]	Dipolar coupling [Hz]	Distance [Å]
experiment	4.4E6	747	1.6
ideal pulses	7.3 / (6.3)E6	964 / (896)	1.47 / (1.5)
10% pulse error	7.3 / (5.7)E6	964 / (850)	1.47 / (1.53)
850 Hz dipolar coupling	5.7 / (5.0) E6	850 / (800)	1.53 / (1.56)
theoretical ( )	7.3E6	964	1.47

Table 12.2: Results for the M2 Calculation and the Simulations

12.3 Final Remarks

The REDOR sequence is a powerful tool to measure distances in different spin systems. But, as seen above, a small error during the setup procedure will lead to severely stretched distances calculated in the interpretation process of the measured data. Although the results using Simpson simulations are in a much better agreement with the expected values, the simulations cannot compensate for the errors introduced by a faulty setup. Additionally it is not always possible to use the simulations for the interpretation of experimental data, e.g. in the case of multispin systems or amorphous systems it may not possible to get reliable input data for the simulations setup.

In any case, in order to be sure of the correct setup of the experiment it is absolutely necessary to proof the experimental setup on a known spin system like the glycine in order to check the robustness of the overall sequence setup. After this validation and calibration process the sequence can then be used to determine distances or M2 values for unknown samples by using the calibrated experimental setup. Experiments on unknown samples should be measured as close to the calibration run as possible to minimize the influence of experimental fluctuations (pressure changes and consecutive spin rate changes, temperature changes and consecutive pulse power changes and the like).

Of course a qualitative comparison within a set of samples is always possible with the same set of experimental parameters without doing a full calibration run.

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M. Bak, J. T. Rasmussen, N. C. Nielsen, Journal of magnetic resonance 147, 296-330 (2000).

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M. Bertmer, H. Eckert, *Dephasing of spin echoes by multiple heteronuclear dipolar interactions in rotational echo double resonance NMR experiments*, Solid State Nuclear Magnetic Resonance 15,139-152 (1999).

13 SUPER

13.1 Overview

Separation of Undistorted Chemical-Shift Anisotropy Powder patterns by Effortless Recoupling (SUPER) correlates CSA powder patterns in the F1 dimension with the isotropic chemical shift in the F2 dimension. The SUPER experiment is based on Tycko's CS – CSA correlation experiment, but provides better compensation for experimental imperfections such as B_1 in-homogeneities and pulse imperfections. Also, both experiments produce scaled powder patterns in F1, and the scaling factor is more favorable in SUPER than the factor 0.39 in Tycko's version. As a consequence, the SUPER experiment does not require high spinning speeds (to fit the F1 line shape into the rotor-synchronized spectral window) or very strong ^{13}C pulses.

SUPER has several advantages. First of all, it covers a large bandwidth for the isotropic chemical shift. Secondly, no requirements exist for 1H decoupling during the recoupling pulses, because it uses 360º pulses instead of the 180× pulses in Tycko's experiment. Exact 360º pulses automatically decouple the heteronuclear dipolar interaction so that no, or only weak, ^1H decoupling is required during the recoupling pulses. The scaling factor is normally 0.155 so that a spectral width over 40 kHz can be achieved in the indirect dimension. As a consequence, moderate spinning speeds of up to 6.5 kHz can be chosen so that experiments can be performed without serious problems on high field instruments. The limiting factor in the choice of the spinning speed is the rotor synchronization requirement of the recoupling 360º pulses:

$$\nu_{RF} = 12.12\nu_{rot}$$

References

S-F. Liu, J-D Mao, and K. Schmidt-Rohr, *A Robust Technique for Two-Dimensional Separation of Undistorted Chemical-Shift Anisotropy Powder Patterns in Magic-Angle-Spinning NMR*, J. Magn. Reson. 155, 15-28 (2002).

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13.2 Pulse Program

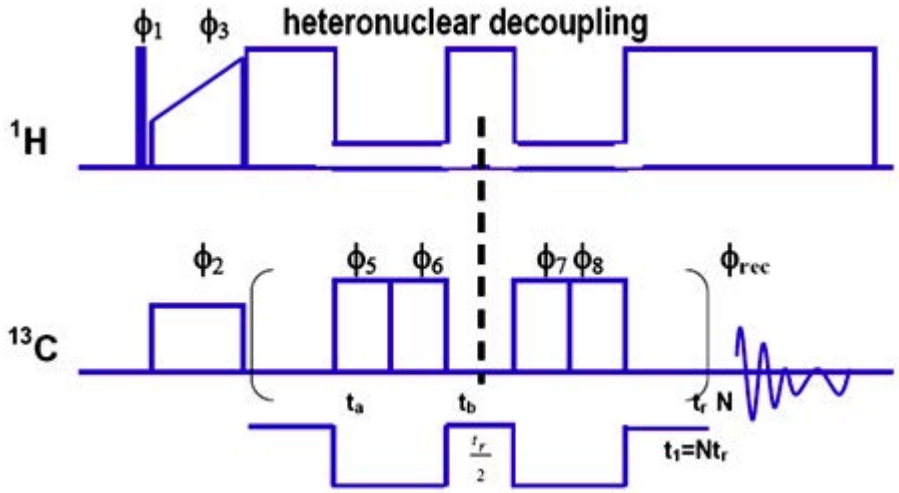

Figure 13.1: Pulse Sequence for 2D CPMAS Exchange Experiment

13.3 2D Experiment Setup

Sample: Tyrosine HCl natural abundance.

Setup time: Less than 1 hour.

13.3.1 Experiment Setup

1. In order to setup the experiment, determine ^{1}H-^{13}C and parameters with variable amplitude on ^{1}H according to *Basic Setup Procedures [▶ 49]*.
2. Verify the pulse parameters on the ^{13}C channel (see *Pulse Calibration with CP [▶ 93]*) and calculate the power level required for the recoupling pulses, i.e. $f_{rf} = 12.2 \times f_{rot}$.
3. Verify the pulse width.
4. Calculate power level required for heteronuclear decoupling during the recoupling pulses, **pl23**, i.e. 20 – 30 kHz or > 25 x f_{rot}. Low power decoupling during the recoupling pulses is permitted because the 360 degree pulses act like heteronuclear decoupling pulses. **pl22** during delays should be high.
5. The experiment requires a minimum of 64 transients to complete the phase cycle. Between 32 and 64 experiments are needed for a 2D data set. Depending on the choice for the gamma integral, more transients per slice may be required. The recommended value is 4, which increases the number of required transients per experiment to 256.
6. Run 1D experiment and make sure everything is set properly.
7. Create a new experiment with either **iexpno** or **edc**.
8. Change to 2D data set:

13.3.2 Setup 2D Experiment

After 1D parameter optimization is completed as previously described, type **iexpno** to create a new data file and switch to the 2D mode using the **123** button. Set the appropriate **FnMode** parameter in **eda**. Pulse program parameters are detailed as follows (the figure in the previous section shows the pulse sequence).

Figure 13.2: The 123 Icon in the Menu Bar of the Data Windows Acquisition Parameter Page

The **123** icon in the menu bar of the data windows acquisition parameter page is used to toggle to the different data acquisition modes, 1D, 2D, and 3D if so desired.

1. Make sure the correct nucleus is selected in F1 dimension, and an appropriate quadrature detection mode is selected in **FnMode** (TPPI, STATES-TPPI or STATES).
2. Choose the appropriate sampling time (**td1**) so that the required resolution (**FIDRES**) in the indirect dimension is achieved.
3. Set **pl11** to give a pulse nutation frequency of12.12* rotation rate (see chapter *Introduction [▶ 9]*).
4. Set **d4**, the z-filter delay, to about 1 ms (integer number of rotor periods if possible).
5. Set **p2** to be a 180× pulse at **pl1** for the TOSS sequence.
6. Set **l5** for the gamma integral, typically = number of spinning sidebands normally 4.
7. Start the experiment.

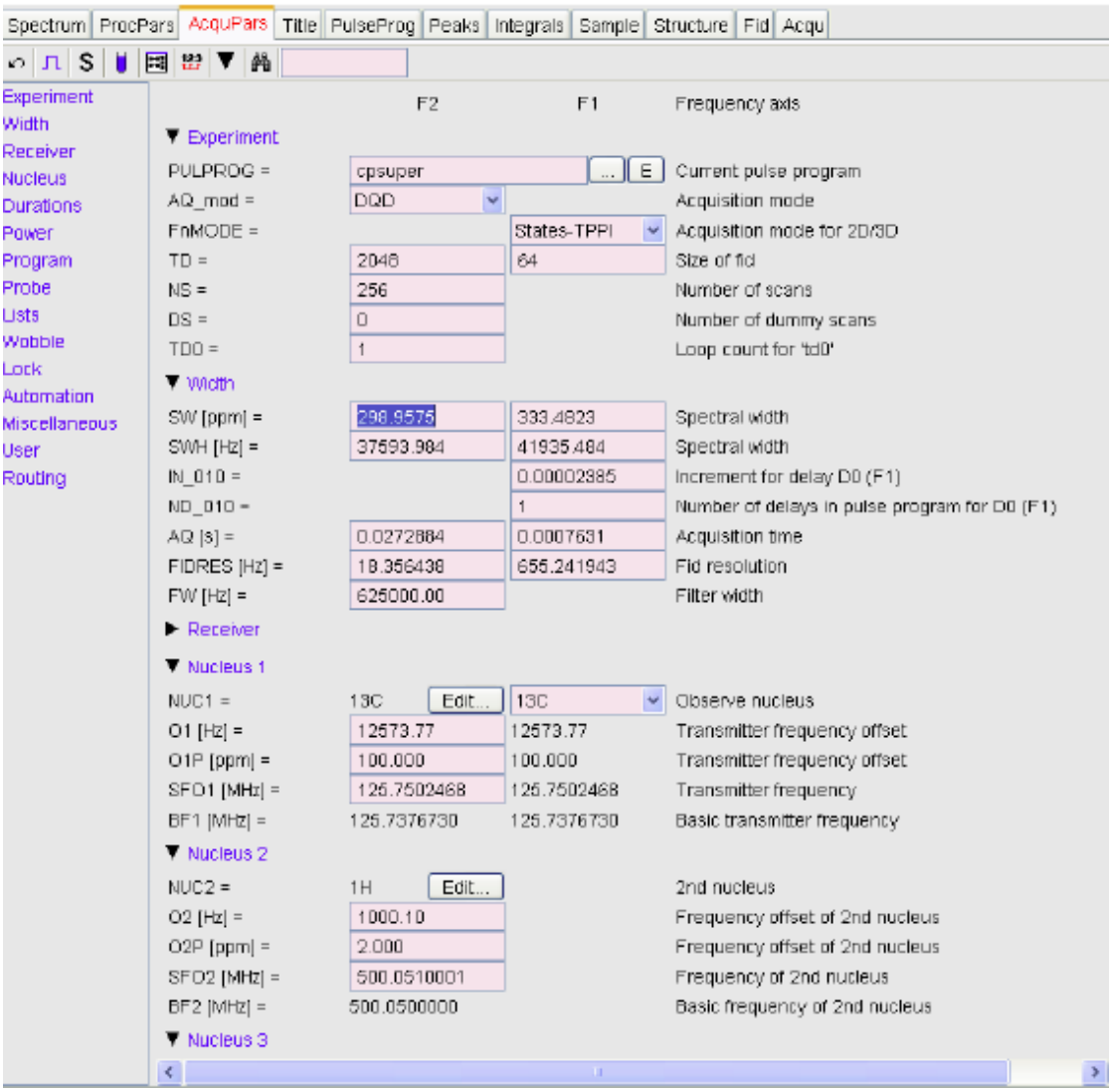

Figure 13.3: The Acquisition Parameter Window eda

13.4 Data Acquisition

Sample: Tyrosine HCl

Experiment time: Several hours

Parameter	Value	Comments
Pulse program	cpsuper	Pulse program.
NUC1	^{13}C	Nucleus on f1 channel.
O1P	100 ppm	^{13}C offset.
NUC2	^{1}H	Nucleus on f2 channel.
O2P	0 ppm	^{1}H offset (can be optimized for best decoupling).
PL1		Power level for f1 channel.
PL11		Power level for f1 recoupling.
P2		180° pulse on F1 during TOSS with PL1.
PL2		Power level for f2 channel.
PL12		Power level decoupling f2 channel and excitation.
P3		Excitation pulse f2 channel.
P15		Contact pulse – first contact.
CPDPRG2		TPPM or SPINAL64.
NS	64*n*l5	Number of scans.
CNST31		Spinning speed in Hz.
L5		L5/cnst31 counter for increment in t1 and number of gamma integral – typically number of SSB's.
F2 direct ^{13}C		Left column.
TD	2048	Number of complex points.
SW	300 ppm	Sweep width direct dimension.
F1 indirect ^{13}C		Right column.
TD	32 - 64	Number of real points.
FnMode		TPPI, STATES or STATES-TPPI.

Table 13.1: Acquisition Parameters

13.5 Spectral Processing

Parameter	Value	Comment
F1 acquisition ^{13}C		Left column.
SI	4096	Number of points and zero fill.
WDW	QSINE	Squared sine bell.
SSB	2	90º shifted sine bell.
PH_mod	pk	Phase correction if needed.
BC_mod	quad	DC offset correction.
Alpha	-1	For shearing the spectrum.
F2 indirect ^{13}C		Right column.
SI	128	Zero fill.
MC2	STATES-TPPI	
WDW	QSINE	Squared sine bell.
SSB	2	90º shifted sine bell.
PH_mod	pk	Phase correction if needed.
BC_mod	no	Automatic baseline correction.

Table 13.2: Processing Parameters

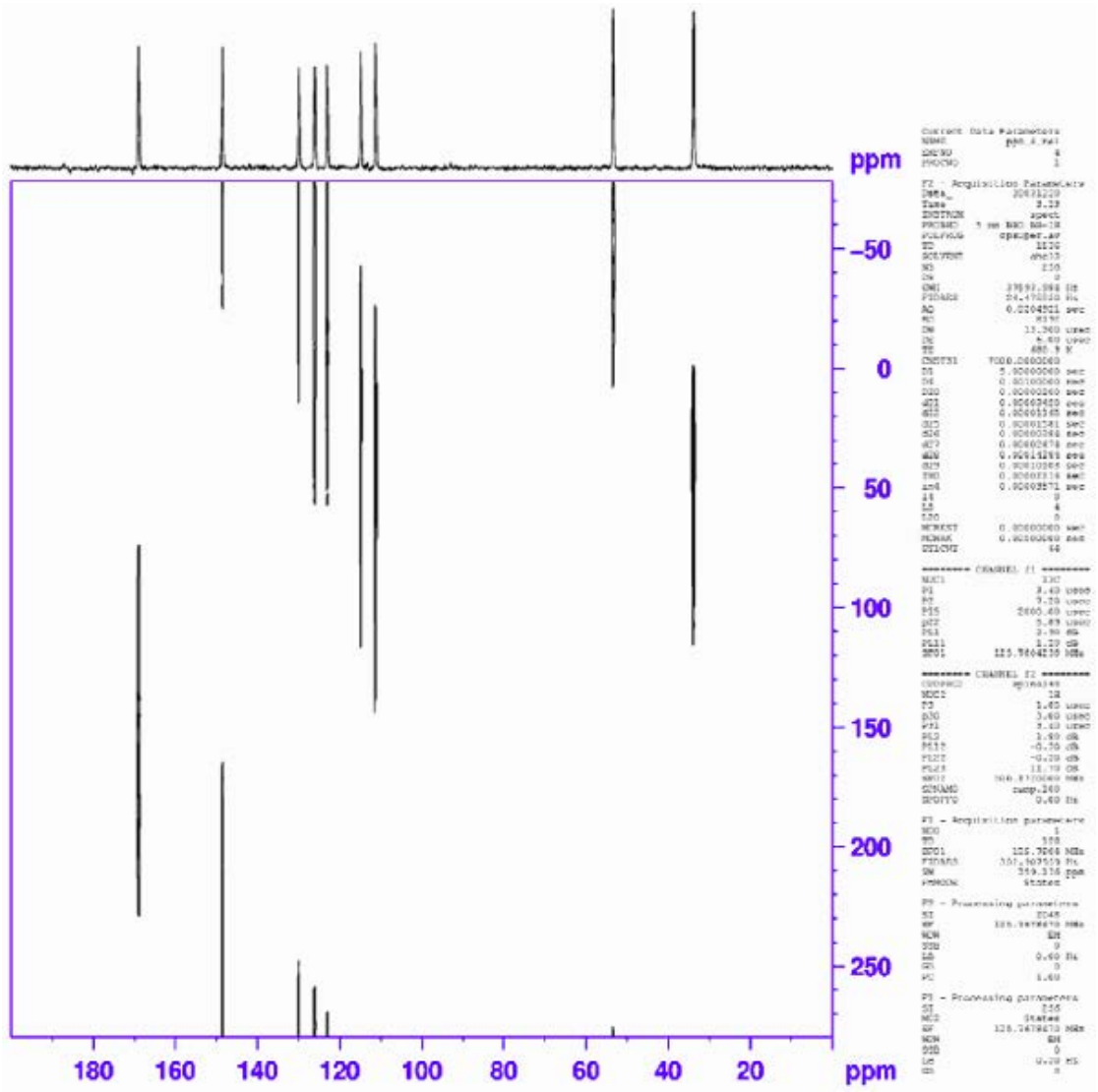

Figure 13.4: The SUPER Spectrum of Tyrosine HCl After Processing Using "xfb"

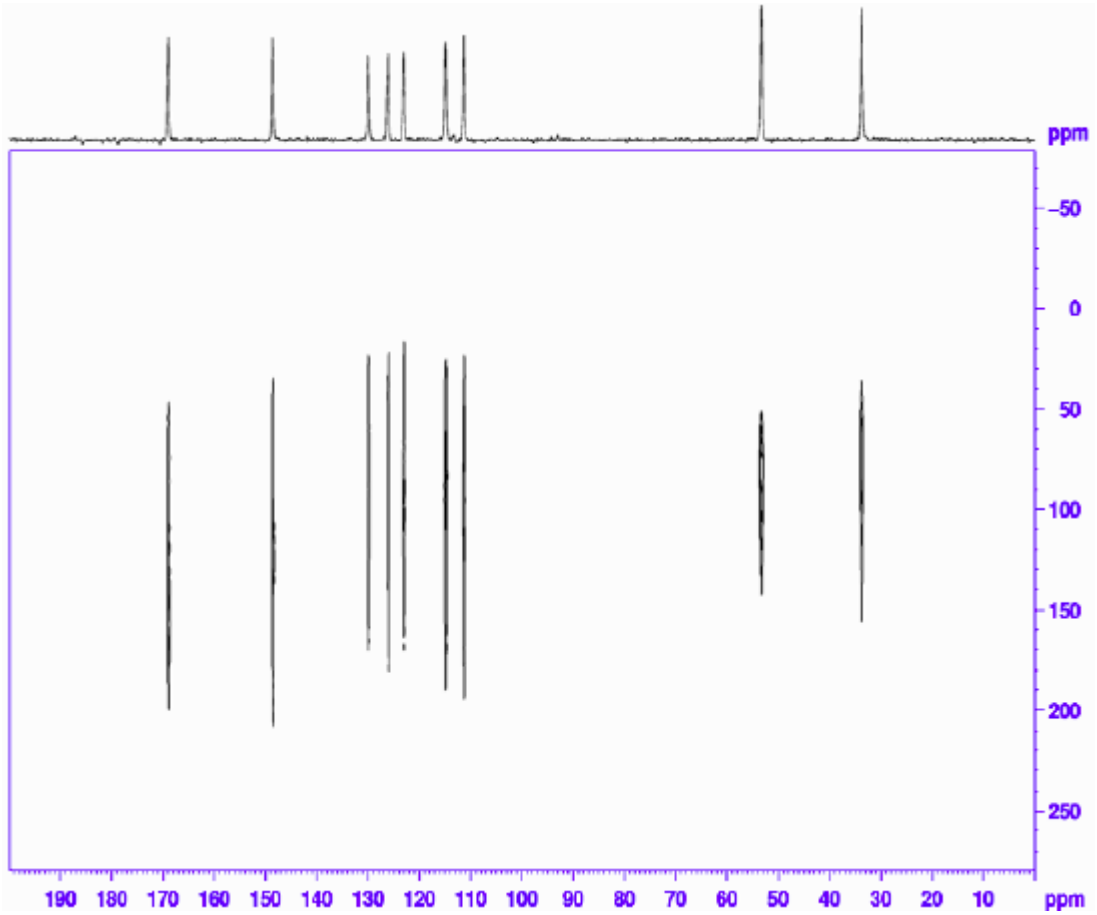

Figure 13.5: SUPER spectrum after tilting the spectrum setting “1 alpha” = -1

The figure above is a SUPER spectrum after tilting the spectrum setting “1 alpha” = -1 and using the command **ptilt1** repeatedly until the CSA lines are within the spectral range.

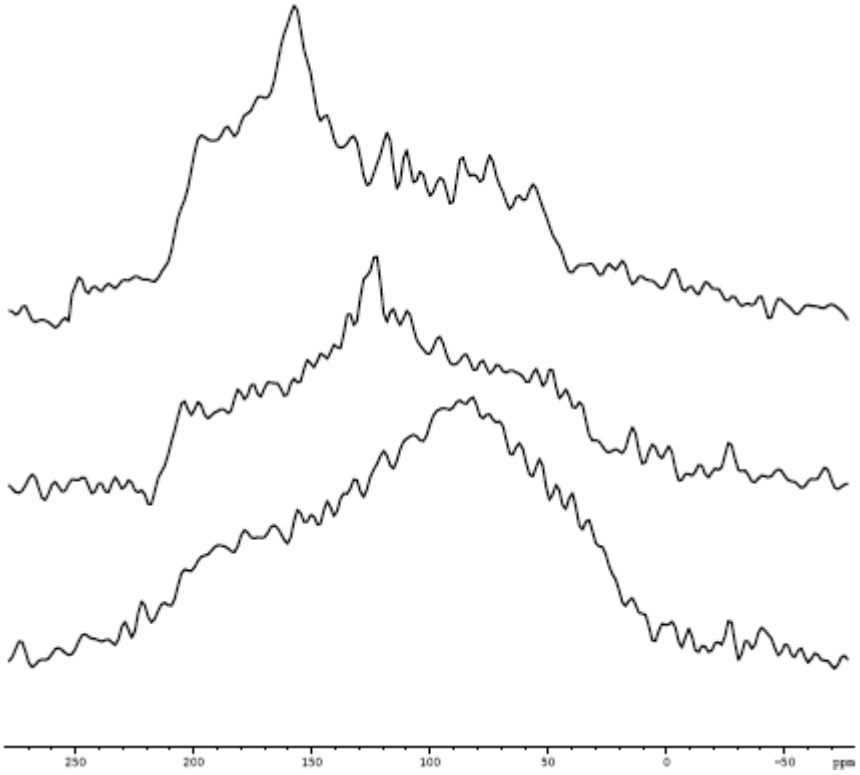

Figure 13.6: Various Cross Sections from the Upper 2D Experiment

The figure above illustrates various cross sections from the upper 2D experiment, from which CSA parameters can be determined.

14 Symmetry Based Recoupling

Sample rotation averages most anisotropic interactions, and therefore removes the information available from them. Therefore, selective recoupling of anisotropic interactions is desired for structural analysis (re-coupling, reintroduction of anisotropic interactions, like e.g. dipolar coupling), in order to regain specific information. The topic has been thoroughly reviewed, by E.A. Bennett et al, and by S. Dusold et al. One strategy is the use of symmetry based recoupling sequences; see M. Hohwy et al (1998) et al. and A. Brinkmann et al. (2000). In these sequences, double quantum coherence are excited via the dipolar homonuclear dipolar coupling. Single quantum coherence are suppressed by phase cycling. The size of the dipolar coupling can be determined from the build-up rate of DQ signal intensity, measured after reconversion into SQ coherence. It should also be mentioned that there are recoupling sequences that do not generate double quantum coherences (DRAWS, DRAMA, and MELODRAMA).

Symmetry-based recoupling sequences recouple specific spin interactions, using cyclic sequences composed of N phase-shifted repetitions of either 2π (C sequences) or π (R sequences) rotation elements. Which interaction(s) are recoupled by a given sequence is determined by the relationship between the sample rotation rate, the spin rotation rate, and the rate of phase shift between the elements. The sequences are denoted as e.g. CN^{v}_{n}, where N is the number of elements in the cycle, n is the number of rotor periods spanned by the N elements, and the total phase rotation between the elements is 2π/v. In the simplest implementation of a C sequence, the 2π rotation element is simply a 2π pulse, but other elements are possible. Thus the sequence $C7^{1}_{2}$ consists of 7 consecutive 2π pulses, with the phase of each pulse shifted by 2π/7 from the previous one. The whole sequence takes two rotor periods, each 2π pulse thus takes 2/7th rotor period. The spin nutation frequency and sample rotation frequency are thus related by $v_{RF} = (7/2)*v_{rotor}$. In practice, the original C7 sequence uses an additional π-phase alternation for every second pulse, so that 14 pulses are executed during 2 rotor periods, requiring $vRF = 7*v_{rotor}$.

For all C and R sequences, the spin nutation frequency must be accurately matched to the sample rotation rate. Since X-X dipolar couplings are usually small, long mixing times are required to reintroduce the dipolar coupling. When ^{1}H decoupling is required, it is important to avoid any transfer of magnetization to or from the proton spin system (HH condition), which would destroy the desired information. This means that the effective fields on X and H must be very different. However, proton decoupling must still be efficient as well. It has been shown that the two RF fields should differ by a factor of 3, which in practice is extremely difficult to meet. It has also been shown that at very high spin rates (>16 kHz) decoupling is not necessary at all. A possible trick is also to use off-resonant LG decoupling during the recoupling sequence. This enhances the effective proton field (vector sum of RF field and offset), and sharpens the HH condition since the homonuclear couplings are suppressed.

Another important parameter to observe is the required excitation bandwidth of these sequences. Naturally, going to higher magnetic fields, the higher chemical shift spread requires higher RF fields for the recoupled X-nuclei, requiring even higher RF fields for protons. So the tendency is going to high spin rates (also desired to get rid of spinning sidebands) and turning the decoupling off during recoupling, which represents a much lower RF load to the probes and increases experimental stability substantially.

The table in *Setup [▶ 163]* shows the sample rotation rate and the required spin nutation frequencies for the X-nucleus. The spin nutation frequency must be 7 times the sample rotation rate for C7, 5 times the sample rotation rate for SPC5 and 3.5 times the sample rotation rate for SC14. Be careful to obey the maximum allowed spin nutation frequencies for the hardware in use.

It is essential that all these parameters are considered carefully in context with the properties of your sample before the experiment is started, so that the appropriate hardware is used. The choice of the MAS-probe is also essential in achieving a sensible setup. The table in *Setup [▶ 163]* shows the selection parameters for three standard recoupling sequences.

References

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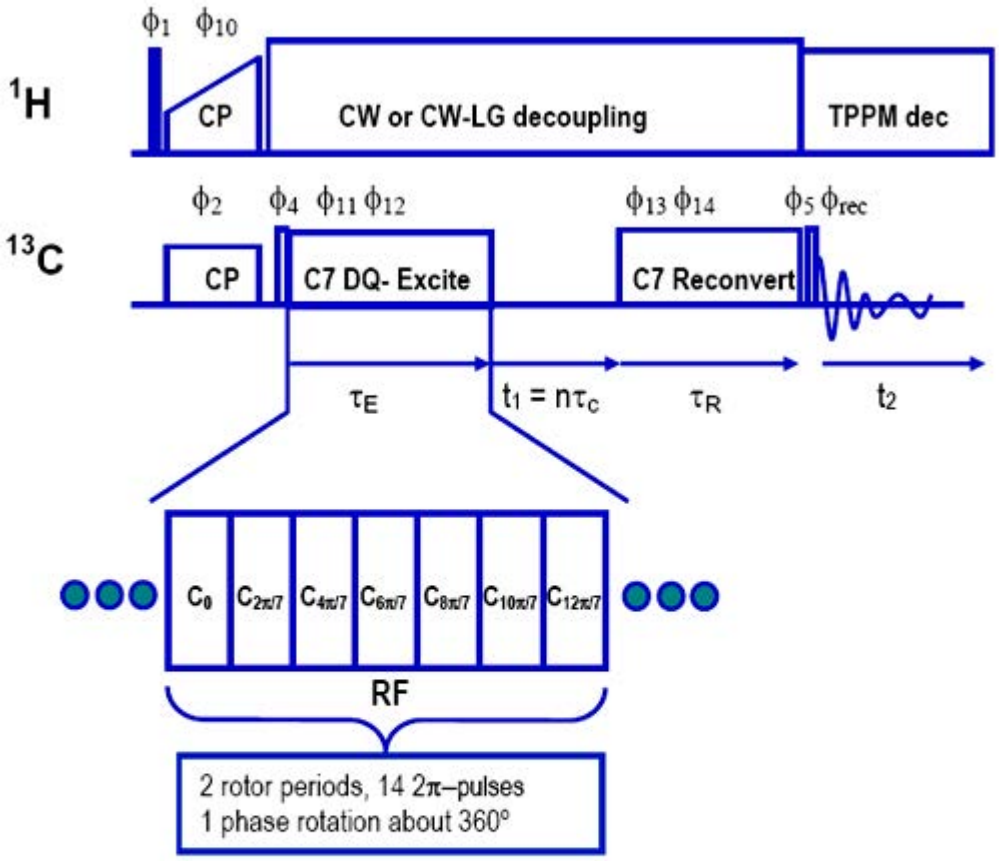
14.1 Pulse Sequence Diagram, Example C7



ϕ_1= 0	ϕ_2 = 0	ϕ_4= 3
ϕ_5 = 1 2 3 0	$\phi_{11,12}$ = φ,φ+180° +QD-phase	$\phi_{13,14}$ = 90° + (φ,φ+180°) +θ_{DQ}
ϕ_{DQ}= 0 1 2 3	ϕ_{rec}= 0 3 2 1	

Figure 14.1: C7 SQ-DQ Correlation Experiment

14.2 Setup

As mentioned before, it is essential that the parameters of your sample of interest are considered before the experiment is started. The table below provides information on the proper choice of hardware for the observe nucleus ^{13}C. Obviously, observation of DQ coherence requires samples with reasonable dipolar couplings and reasonable probability of coupled species. So, running this experiment on 13C samples requires reasonable enrichment. Usually, fully enriched samples are used, sometimes diluted in natural abundance samples to reduce nonspecific long range interactions. As always, rotary resonance conditions (overlap of side- and center bands) should be avoided unless specifically desired.

Running the experiment on enriched ^{15}N samples is of course possible, but one should consider that most samples will not have nitrogen atoms directly attached to each other, so small couplings will prevail, requiring long DQ excitation and DQ reconversion times, with a nucleus that requires high RF power levels to achieve a certain RF-field. On the other side, the shift range is not large, allowing relatively slow spinning. Considering a nucleus like ^{31}P, there is no need for enrichment, but cases with directly bonded ^{31}P-atoms are rare. Phosphates are usually easy, since the shift range is small (couplings are also rather small). If however a large shift range (possible with ^{31}P) needs to be covered, there may be a substantial problem.

Sequence	n= (vRF/ masr)	Rotor diameter/ masr max.	Masr max. rec.4 (Hz)	vRF(13C) max. (kHz/µs) 5	vRF(H) max. (kHz/µs) 6	Bo max. (MHz) 8
POST-C7 [1]	7	7/6000	5000	35 / 7.15	70 / 3.5 +LG	300
		4/15000	9000	63 / 4	100 / 2.5 +LG [7]	500
		3.2/24000	12000	84 / 3	110 / 2.27 +LG [7]	600
		2.5/35000	14000	100 / 2.5	130 / 1.95 +LG [7]	800
SPC5 [2]	5	7/6000	5000	25 / 10	70 / 3.5	300
		4/15000	13000	65 / 3.85	100 / 2.5 +LG [7]	500
		3.2/24000	17000	85 / 3	none	700
		2.5/35000	20000	100 / 2.5	none	900
SC14 [3]	3.5	7/6000	6000	21 / 12	70 / 3.5	200
		4/15000	15000	52.5 / 4.75	100 / 2.5 +LG [7]	500
		3.2/24000	22000	77 / 3.25	none	700
		2.5/35000	28000	100 / 2.5	none	950

Table 14.1: Recommended Probe/Spin Rates for Different Experiments and Magnetic Field Strengths

1. C7 is not recommended due to restricted excitation bandwidth.
2. SPC5 can be recommended as a standard sequence for 4 mm probes and not too high fields.
3. SC14 or sequences with similar RF-field requirements are recommended for small spinners/high fields.
4. Maximum speed results from max. possible RF-field.
5. Maximum ^{13}C RF fields taken from ^{13}C RF field specification, or ^{1}H RF field specification, considering the requirement of an off HH condition.
6. Maximum RF field for decoupling.
7. +LG means cw decoupling with optimized LG-offset frequency at the given RF field in order to avoid HH contact.
8. Maximum magnetic field as proton resonance frequency in MHz. This results from spin rate requirements for ^{13}C observation (to avoid rotary resonance conditions) as well as excitation bandwidth considerations.

14.2.1 Spectrometer Setup for 13C

1. Load a CPMAS parameter set for ^{13}C.
2. Load a uniformly labeled glycine sample and rotate at the desired rotation rate (see table in *Setup [▶ 163]*), depending on the recoupling experiment planned and the sample under investigation. Consider possible rotational resonance conditions in the sample of interest!
3. Tune and match the probe, optimize the ^{13}C and ^{1}H pulse parameters for excitation and decoupling.
4. Use the **cp90** pulse program with **pl11= pl1** to measure the nutation frequency for ^{13}C, in order to calculate the recoupling conditions (see chapter *Basic Setup Procedures [▶ 49]*). Calculate the power levels required by the spin speed (see table in *Setup [▶ 163]*) using **calcpowlev**.
5. Set **pl11** back to 120 dB (**p1** to zero) and run 1 experiment with 16 (4) scans as a reference.

14.2.2 Setup for the Recoupling Experiment

1. Create a new experiment and load the appropriate pulse program (**spc5cp1d**), use the same routing. Set the appropriate sample rotation rate, set **cnst31** equal to the rotation rate.
2. Load the power level calculated for the necessary ^{13}C recoupling B1-field into **pl11**, set **p1** as determined in step 4.
3. Set **l0=15** (should be, but need not be, a multiple of 5 for SPC5 or of 7 for PC7, SC14). This determines the DQ-build-up time (DQ generation). The reconversion time is usually also controlled by **l0**, it may however be written such as to be independently controlled by a different loop counter. For glycine, about 5 msec will be the optimum.
4. Set the decoupling program **cpdprg1** to **cwlg**. Set **pl13** such as to yield the desired decoupler RF field during the DQ generation/ reconversion, or set it to 120 if the spin rate suffices to omit decoupling. Set **cnst20** = corresponding decoupling RF field.
5. Optimize **pl11** for maximum signal intensity.
6. Optimize **l0** for optimum signal intensity. In a multi-site spectrum the optima may differ for different spin pairs.

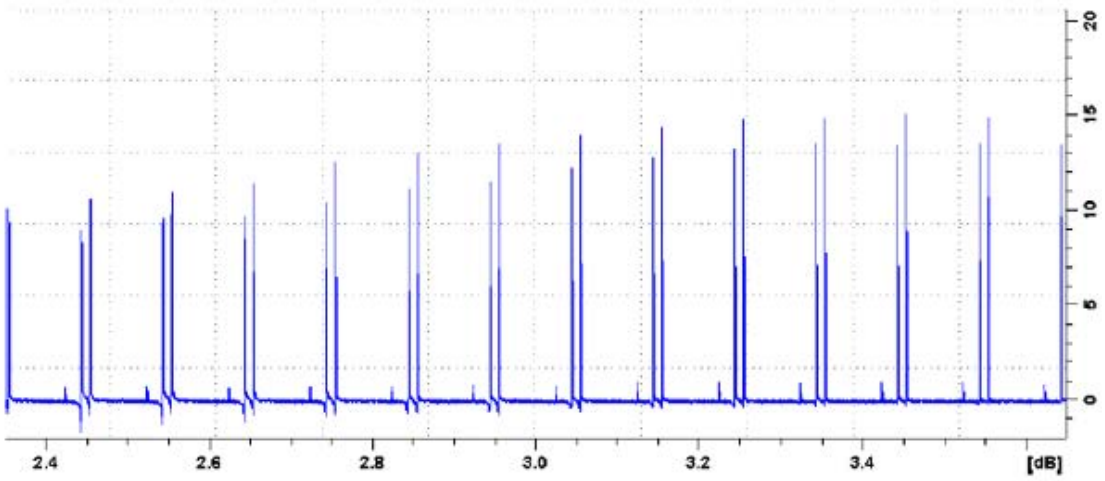


Figure 14.2: Optimization of the RF Power Level for DQ Generation/Reconversion on Glycine

In principle, both peaks must grow together as one approaches the RF=7*MASR condition, but the resonances are differently influenced by non-ideal off-HH conditions. The glycine α-peak is usually hard to get off HH, so it is frequently too small. Optimize the LG decoupling condition on the glycine α-peak (step 4 in section *Setup of the 2D SQ-DQ Correlation Experiment [▶ 167]*).

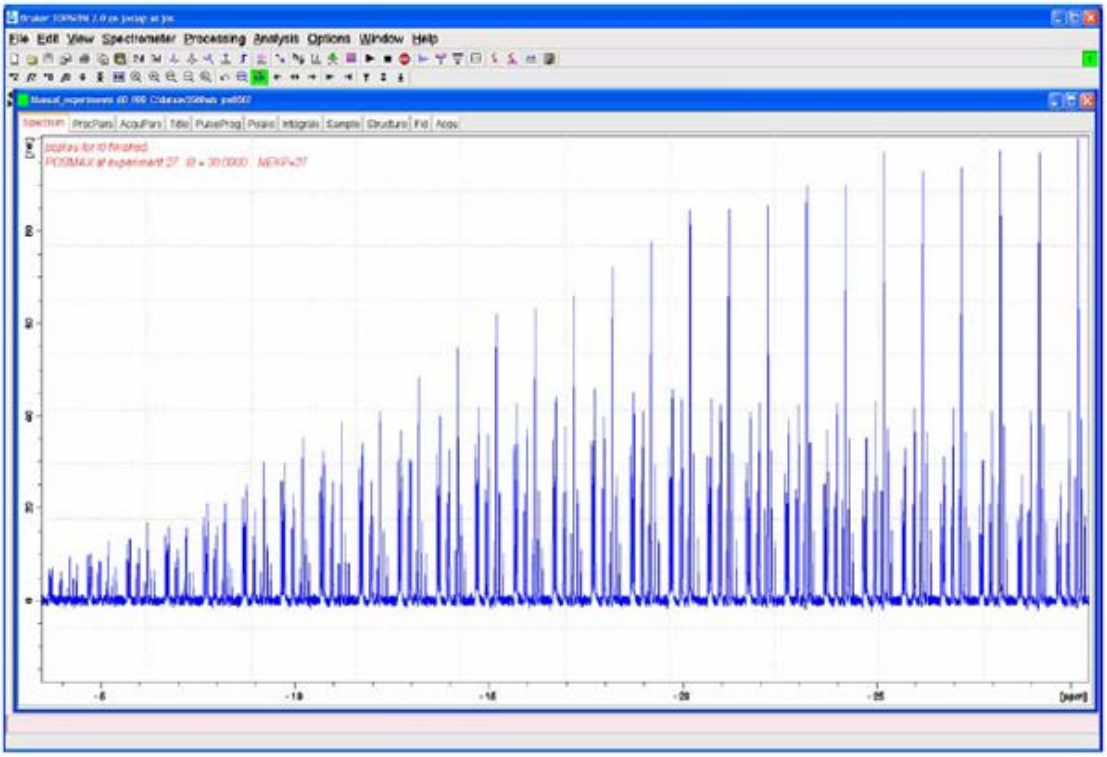


Figure 14.3: Variation of DQ Generation/Reconversion Time on a Uniformly 13C Labeled Peptide (fMLF)

Both times were incremented in units of 2 rotation periods. One can clearly see the different maxima for the Ca, the alphatic carbons and the mobile CH_3-groups. Spinning speed was 13 kHz.

1. Optimize the **cwlg** decoupling if needed by variation of **cnst20** in increments of 5000 and check whether a different offset condition helps improving the signal intensity.
2. Run one experiment and compare with the direct CP experiment to measure the DQ recoupling yield.

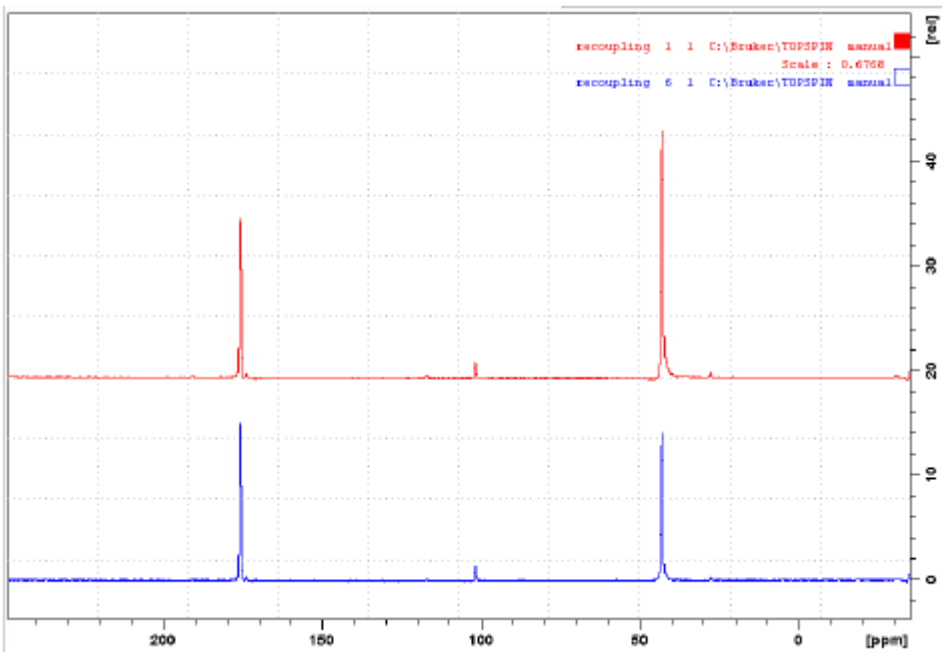


Figure 14.4: PC7 Recoupling Efficiency at a Spinning Speed of 13 kHz

The figure above illustrates PC7 recoupling efficiency at a spinning speed of 13 kHz (about 100 kHz RF field), using a 2.5 mm probe. LG decoupling at 125 kHz was used during DQ generation/reconversion. Quite a noticeable loss on the glycine α-peak due to insufficient HH suppression is visible. Efficiency is 67% on the carboxyl peak (AVIII 700 SB).

14.2.3 Setup of the 2D SQ-DQ Correlation Experiment

1. Running such a correlation experiment on glycine makes little sense, so insert a sample with more ^{13}C sites (fully labelled tyrosine-HCl, or histidine or any other suitable labelled sample). Optimize **l0** for the best compromise in signal intensities.
2. Generate a new data set, set the mode to 2D using the **123** button in **eda**. Load the pulse program **spc5cp2d**.
3. Make sure ^{13}C is selected as **nuc1** in the F1 dimension, set **FnMode** =STATES-TPPI.
4. Set the spectral window along F1. It is desirable to synchronize sampling along F1 with the rotor spin rate in order to eliminate spinning sidebands (fold back onto center band). This may however lead to peak fold over, since achievable spin rates are usually smaller than the spread of DQ-frequencies along F1. This does however not necessarily mean that the spectra are crowded and uninterpretable, because frequently the folding does not lead to cross peak overlap. Of course, the synchronization to the spin rate need not be used, **cnst31** can be set equal to the sweep width along F2 which will normally produce spectra free of folding, but of course, spinning sidebands along F1 will occur and signal intensity will be spread over a larger number of cross peaks. An intermediate sampling rate along F1 can be achieved by incrementing the evolution period synchronized to the phase shifted blocks of the sequence (one PC7-block being $2\tau/7$, τ=rotor period). This will also not generate sidebands along F1, but provide a larger sweep width and less fold over (M. Hong 1999). Fold over can often be tolerated, **xfshear rotate** may be used to shift the spectrum suitably along F1.
5. Set the acquisition time along F1 to about 10 msec for a start. Lines along the double quantum dimension may be narrower than along the single quantum dimension, so a compromise between experiment time and digital resolution along F1 must be found.
6. Start the experiment.

14.3 Data Acquisition

Sample: Fully ^{13}C labelled tyrosine-HCl, or a suitable fully labelled small peptide

Spinning speed: 5 – 20 kHz, depends on experimental requirements, see the table in *Setup [▶ 163]*.

Experiment time: 1-4 hours.

Parameter	Value	Comments
Pulse program	spc5cp2d spc5cp2dlsw sc14cp2d r14cp2d pc7cp2d	See the table in *Setup [▶ 163]* for hints which sequence to prefer. Rule of thumb: high field: fast spinning, sc14. Low field, slow spinning, pc7. Intermediate: spc5. Note: sc14 usually has low DQ yield (35%), but that may not matter.
NUC1	^{13}C	Nucleus on f1 channel.
O1P	100 ppm	^{13}C offset.
NUC2	^{1}H	Nucleus on f2 channel.
O2P	2-3 ppm	^{1}H offset.
PL1	for > 50 kHz vRF	Power level for f1 channel CP and p1.
PL11	dep. on masr	Power level for f1 channel recoupling power.
PL12	as specified	Power level decoupling f2 channel and excitation.
PL13	≈pl12, optimize, or 120, fast spinning	Power level decoupling f2 channel during cw or cwlg decoupling.
P3		Excitation pulse f2 channel.
PCPD2		Decoupler pulse length f2 channel (^{1}H) TPPM.
P15		Contact pulse – first contact.
D1		Recycle delay.
CNST20		Spin nutation frequency at PL13 for cwlg decoupling.
L0	for 0.5-10 msec („mix“ in ased)	Use multiples of 5, 7, or 16 (spc5, pc7, sc14) for full phase cycle.
SPNAM0		Ramp for 1st CP step; e.g. ramp: 80 – 100%
SP0		Power level for proton contact pulse.
CPDPRG2	SPINAL64	SPINAL64 decoupling.
CPDPRG1	cwlg	To avoid HH contacts during DQ-generation, reconversion.
NS	4-32	Number of scans (see pulse program phase cycle).
F2 direct ^{13}C		Left column.
TD	1024 or 2048	Number of complex points.
SW		Sweep width direct dimension, adjust to experimental requirements.
F1 indirect ^{13}C		Right column.
TD	128 - 512	Number of experiments in indirect dimension.

SW	see para. 17 above	Sweep width indirect dimension.
ND0	1	STATES-TPPI, not required. In TS 2.1.

Table 14.2: Acquisition parameters for DQ-SQ correlation experiments using symmetry based recoupling sequences

14.4 Spectral Processing

Processing Parameters

Parameter	Value	Comment
F1 acquisition ^{13}C		Left column.
SI	2-4 k	Number of points and zero fill.
WDW	QSINE	Sine bell squared.
SSB	2-5	Shifted sine bell.
PH_mod	pk	Phase correction if needed.
F2 indirect ^{13}C		Right column.
SI	256-1024	Zero fill.
MC2	STATES-TPPI	
WDW	QSINE	Sine bell squared.
SSB	2	90º shifted sine bell.

Table 14.3: Processing parameters for DQ-SQ correlation experiments using symmetry based recoupling sequences

Processing parameters for DQ-SQ correlation experiments using symmetry based recoupling sequences. The AU program `xfshear` may be used with option **rotate** and argument(+/- δ ppm) to shift the spectrum suitably along F1. Setting 1 **sr** = 2***sr+o1** will set the referencing along F1 correctly (just type **sr**, and in f1 enter the value of sr for F2 and add *2+o1).

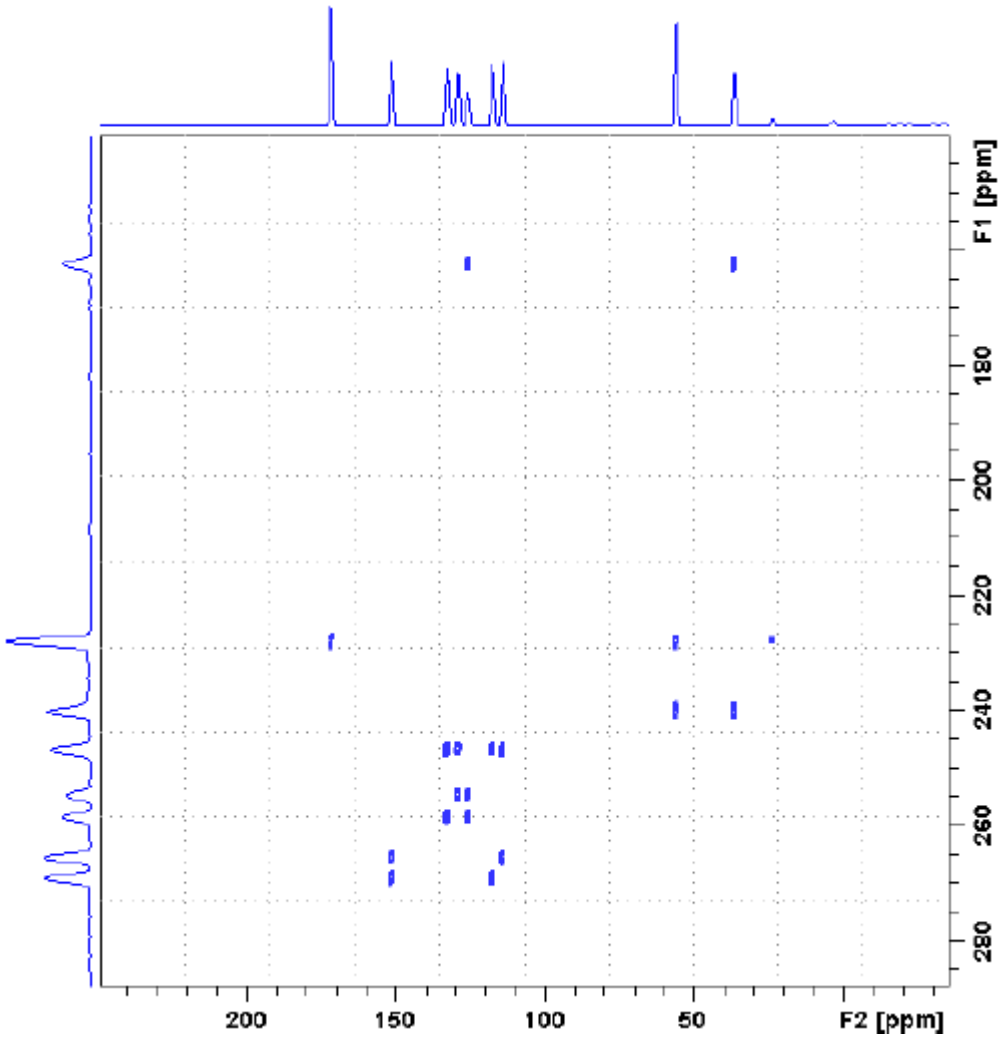


Figure 14.5: SC14 2d SQ-DQ Correlation on Tyrosine-HCl

SC14 2d SQ-DQ correlation on tyrosine-HCl, 56 rotor periods mixing at 26 kHz, 2.5 mm probe, AV III 700 SB. With the sampling window along F1= spin rate, only the α-β-correlation is folded.

14.5 13C-13C Single Quantum Correlation with DQ Mixing

Symmetry based DQ recoupling sequences may also be used as mixing periods in SQ-SQ correlation experiments. The experiment resembles the PDSD or RFDR experiments (see *Proton Driven Spin Diffusion (PDSD) [▶ 129]* and *RFDR [▶ 123]*) as a NOESY-type correlation will be generated. Similarly, the MELODRAMA (see Bennett et al., Dusold et al.) sequence with $v_{RF} = 5*v_{rotor}$ may be used here.

14.6 Data Acquisition

Parameter	Value	Comments
Pulse program	pc7cp2dnoe	Any sequence may be used, make sure to use the correct timing.
NUC1	^{13}C	Nucleus on f1 channel.
O1P	100 ppm	^{13}C offset.
NUC2	^{1}H	Nucleus on f2 channel.
O2P	2-3 ppm	^{1}H offset.
PL1	for > 50 kHz vRF	Power level for f1 channel CP and p1.
PL11	dep. on masr	Power level for f1 channel recoupling power.
PL12	as specified	Power level decoupling f2 channel and excitation.
PL13	≈pl12, optimize, or 120, fast spinning	Power level decoupling f2 channel during cw or cwlg decoupling.
P3		Excitation pulse f2 channel.
PCPD2		Decoupler pulse length f2 channel (1H) TPPM.
P15		Contact pulse.
D1		Recycle delay.
CNST20		Spin nutation frequency at PL13 for cwlg decoupling.
L0	for 0.5-10 msec („mix“ in ased)	Use multiples of 5,7, or 16 (spc5, pc7, sc14)for full phase cycle.
SPNAM0		Ramp for 1st CP step; e.g. ramp: 80 – 100%.
SP0		Power level for proton contact pulse.
CPDPRG2	SPINAL64	SPINAL64 decoupling.
CPDPRG1	cwlg	To avoid HH contacts during DQ-generation, reconversion.
NS	4-32	Number of scans (see pulse program phase cycle).
F2 direct 13C		Left column.
TD	1024 or 2048	Number of complex points.
SW		Sweep width direct dimension, adjust to experimental requirements.
F1 indirect ^{13}C		Right column.
TD	128 - 512	Number of experiments in indirect dimension.
SW	usually = sw (F2)	Sweep width indirect dimension.
ND0	1	STATES-TPPI, not required in TopSpin 2.1.

Table 14.4: Data Acquisition

14.7 Spectral Processing

Processing parameters

The processing parameters are the same as in section *Data Acquisition [▷ 171]*.

Parameter	Value	Comment
F1 acquisition ^{13}C		Left column.
SI	2-4 k	Number of points and zero fill.
WDW	QSINE	Sine bell squared.
SSB	2-5	Shifted sine bell.
PH_mod	pk	Phase correction if needed.
F2 indirect ^{13}C		Right column.
SI	256-1024	Zero fill.
MC2	STATES-TPPI	
WDW	QSINE	Sine bell squared.
SSB	2-5	90° shifted sine bell.

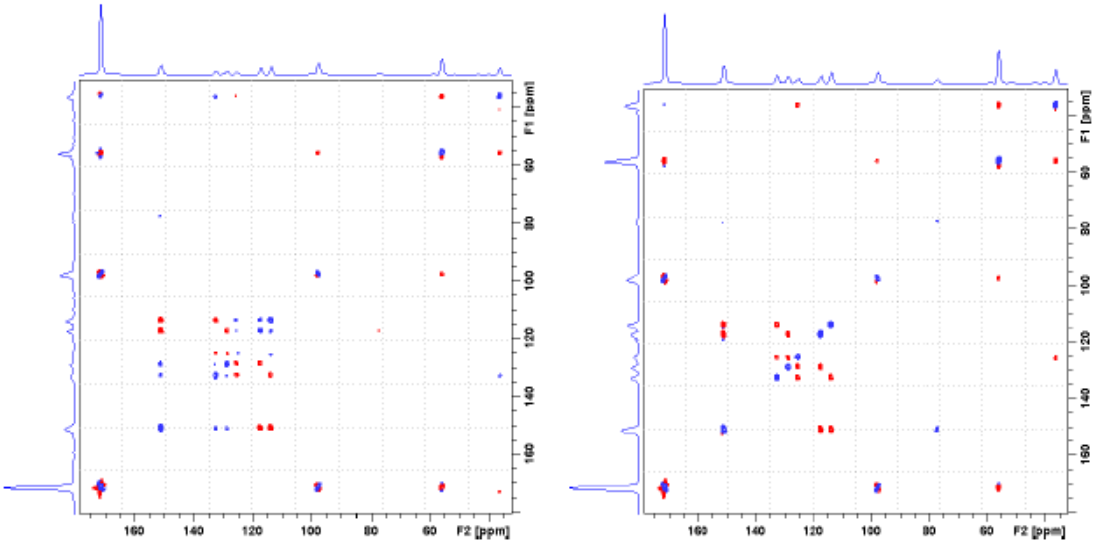


Figure 14.6: PC7 2d SQ-SQ Correlation on Tyrosine-HCl

PC7 2d SQ-SQ correlation on tyrosine-HCl, 56 rotor periods mixing at 13 kHz, 2.5 mm probe, AV III 700 SB. Left: 84 rotor periods DQ mixing, right:56 rotor periods mixing. The 84 periods mixing time show relayed correlations (positive, blue) which are absent at 56 periods mixing (except for the SSB cross peaks). Direct correlations are negative (red).

15 PISEMA

Polarization Inversion Spin Exchange at the Magic Angle is an experiment that correlates the chemical shift of a spin 1/2 X nucleus with the heteronuclear dipolar coupling to another spin 1/2 nucleus. Most of the applications so far reported have been in the field of structural biology, therefore, the X nucleus is normally ^{13}C or ^{15}N and the other hetero-nucleus ^{1}H. The experiment provides orientation information on the vector connecting the ^{13}C or ^{15}N and the ^{1}H nucleus. The achievable high resolution of the CS, as well as, the dipole coupling makes the experiment well suited for 3D NMR experiments on aligned systems or single crystals.

Unlike normal FSLG experiments, where the dipolar and CS interactions are scaled by $\cos(\theta_m) = 0.577$, the scaling factor for the heteronuclear dipolar interaction is $\sin(\theta_m) = 0.816$, because the coupling takes place in the transverse plane of the rotating frame, the spin locked state. The projection is from the tilted frame (locked ^{1}H spin system) to the transverse plane of the rotating frame system (spin locked ^{15}N spin system), i.e., $\sin(\theta_m) = 0.816$ the scaling of the heteronuclear dipolar coupling strength.

Through the combination of spin exchange (dipolar flip flop term) and the homonuclear decoupling using FSLG, PISEMA achieves a line width that is an order of magnitude better than its predecessor, the separated local field experiment.

The central line in the dipolar dimension can be caused among other things by a proton frequency offset introducing a constant term in the time domain signal. That offset frequency also makes the splitting larger. See additional test procedures in the paper about „*Experimental aspects of multidimensional solid state correlation spectroscopy.*“, listed in the references below.

PISEMA is not very sensitive to the exact Hartmann-Hahn condition. A mismatch has little effect on the dipolar coupling. The scaling factor in the indirect dimension depends of the ^{1}H resonance offset and a wrong ^{1}H carrier frequency can cause a wrong scaling factor and some intensity loss and, as mentioned above, a zero frequency contribution. Diligent adjustment of the LG condition and the RF-carrier is critical for accurate measurement of the dipolar coupling as the splitting increases quadratic with increasing (proton) frequency offset.

Simulations of the spin dynamics show that the heteronuclear term in the Hamiltonian leads to a complicated spectrum for small heteronuclear dipolar couplings (usually introduced by remote protons), see Z. Gan's paper for more information.

References

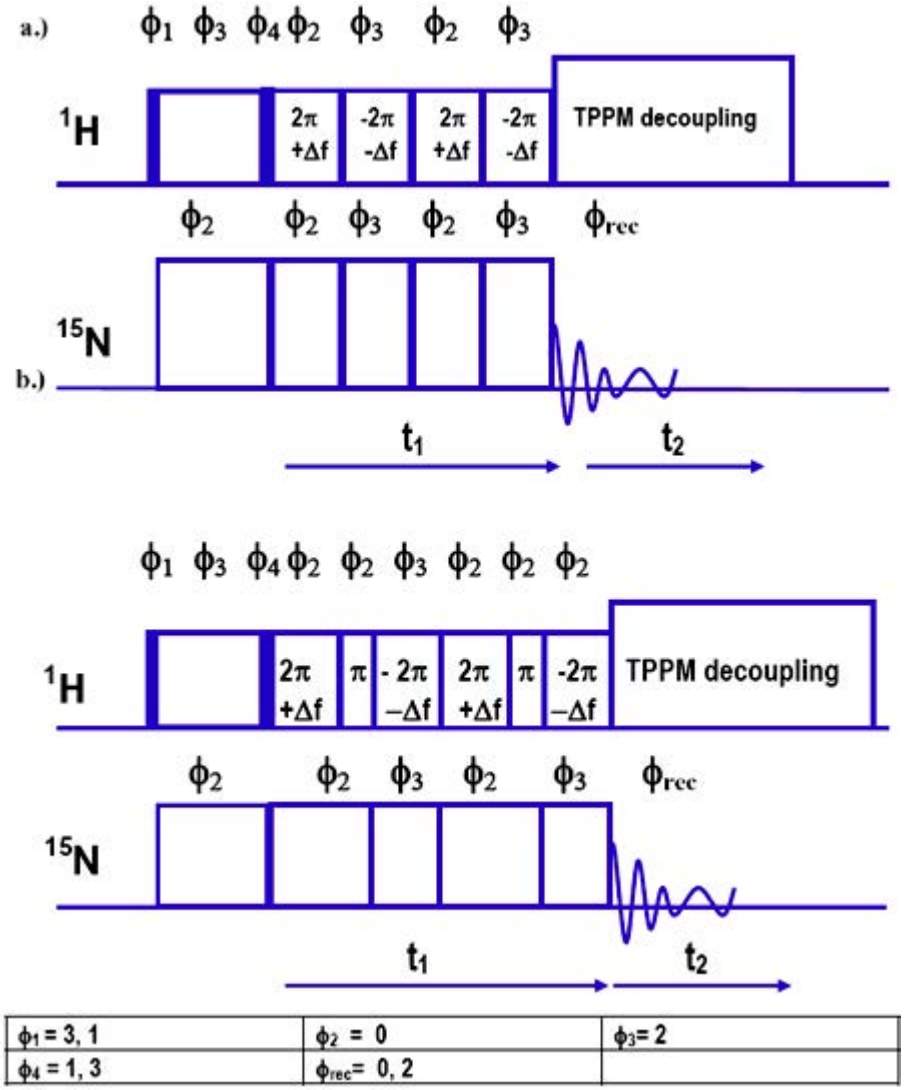
C.H. Wu, A. Ramamoorthy, and S.J. Opella, *High-Resolution Heteronuclear Dipolar Solid State NMR Spectroscopy*, J. Magn. Reson. A 109, 270-272 (1994).

A. Ramamoorthy, C.H. Wu, and S.J. Opella, *Experimental Aspects of Multidimensional Solid State NMR Correlation Spectroscopy*, J. Magn. Reson. 140, 131-140 (1999).

A. Ramamoorthy, and S.J. Opella, *Two-dimensional chemical shift / heteronuclear dipolar coupling spectra obtained with polarization inversion spin exchange at the magic-angle sample spinning* (PISEMAMAS), Solid State NMR 4, 387-392 (1995).

Zhehong Gan, *Spin Dynamics of Polarization Inversion Spin Exchange at the Magic Angle in Multiple Spin Systems*, J. Magn. Reson. 143, 136-143 (2000).

15.1 Pulse Sequence Diagram



$\phi_1 = 3, 1$	$\phi_2 = 0$	$\phi_3 = 2$
$\phi_4 = 1, 3$	$\phi_{rec} = 0, 2$	

Figure 15.1: Pisema Pulse Sequence

In the Pisema pulse sequence above: a.) straight PISEMA, b.) “clean PISEMA” variation for further suppression of phase glitches (Ramamoorthy et al. Solid State NMR 4).

15.2 Setup

Sample: ^{15}N labelled α-glycine powder for power level determination, ^{15}N labelled acetylated glycine or acetylated valine or leucine for running the PISEMA experiment, preferably as single crystal.

Setup time: 0.5 hours on labelled glycine.

Experiment time: 15 hours on a labelled powder sample, 1-2 hours on a good size single crystal.

1. Set up for static ^{15}N CP observation on the α-glycine powder sample, pulse program **cp**. Use a ramp pulse if the HH condition is unknown, with power level settings for an approximate 5 µsec pulse on both channels.
2. Determine the proton 90 degree pulse **p3** at the respective power level, reset the conditions for a square shape and about 50 kHz on both channels for contact. Load the pulse program `cplg`. With **cnst17**=0, **pl2=pl13** and **pl1** all set for 50 kHz, reestablish the HH condition.
3. To adjust the CP condition under a LG-offset, load the pulse program `cplg`. Cnst 17 sets the LG-offset during the contact, the offset frequencies are calculated as cnst 18 and cnst19. Start with cnst17=0.
4. Two possibilities exist to set the FSLG power levels and offset frequencies.
   - Use the appropriate offset frequency for the chosen contact power level of 1H and set **cnst20** accordingly to e.g. 50 kHz, i.e **cnst17 = 50000.0**. This would give an offset frequency of approximately 35 kHz (**cnst19** should show this number in the **ased** display) Then adjust the ^{15}N power level during the FSLG period to best HH match - which is at a power level of appropriately 20*log(sin(54.7))=1.8 dB higher than for the on resonance contact.
   - If that option is not adequate because of power limitations on ^{15}N, one can also leave the on resonance contact levels of ^{15}N and calculate the offset frequency and power level for ^{1}H. That reduces the required proton RF power by about 70% as compared to the power level for a resonant HH match. For the new nutation frequency (B1 field for LG condition):

$$B_{1LG}(^1H) = \sin(\theta_m) * B_{1on\_res}(^1H) = 0.82 * B_{1on\_res}(^1H)$$

The offset frequency for the Lee-Goldburg condition is:

$$f_{LG} = \cos(\theta_m) * B_{1on\_res}(^1H) = 0.578 * B_{1on\_res}(^1H)$$

With the inverse of a 360º pulse:

$$B_{1on\_res} = 1/(\tau_{2\pi})$$

Instead of raising the power level for ^{15}N, the power level for ^{1}H is reduced by about 1.7 dB. Then the new 2π pulse in the tilted frame is:

$$\tau_{2\pi_{LG}} = B_{1on\_res}^{-1} \sin^2(\theta_m) = 0.67 B_{1on\_res}^{-1}$$

In our example of a contact power level of 50 kHz on ^{15}N one would then calculate for **cnst17 = 40 807.0,** giving an offset frequency of 28855 Hz for the LG frequency, which is calculated automatically.

1. In order to verify all calculated power levels and offset frequencies, optimize for the appropriate power level using the pulse program `cplg`.
2. Change to the desired PISEMA sample. Since a powder contains all possible spin pair orientations, the measurement of an oriented sample is recommended, because it is not only much faster, but also allows the performance to be judged much better. N-acetly valine can be rather easily grown to sizeable single crystals and has all properties for a good setup sample: decent proton T1, well defined sites and therefore narrow

resonances. For a good crystal, the residual proton line width will reflect the quality of the setup. Beware of crystal twinning. The orientation of the sample should be selected carefully.

3. If the sample does tune and match very differently than the setup sample, check the HH conditions briefly and verify that the parameters found are valid. Correct power levels or pulse parameters if needed. This is especially important for saline lipid water mixtures.
4. Depending on the decision taken in (4), set
   - Either PL13 = PL2 (case 4a) and set PL11 to the optimized value, higher power, i.e. a value of about 1.8 dB below pl1;
   - Or set PL11 = PL1 (case 4b) and PL13 to the obtained value in the cplg experiment, i.e. about 1.8 dB higher than PL2, which is less RF power.
   - Depending on the orientation of the single crystal, O1 and O2 need to be re-optimized since the peak positions will change. Select an orientation that shows the most low-field ^{15}N peak positions, because that will usually correspond to the biggest ^{1}H-^{15}N dipolar couplings.
5. Re-optimize the power levels for the HH-contact.
6. Create a new experiment and set up a 2D data set, using the **1,2,3** icon. Load the `pisema` pulse program. Go into the **eda** window.
7. Make sure the correct nucleus is selected in the F1 dimension.
8. In order to set the t1 increment, go into the **ased** window, click the pulse symbol, and choose **I3** to be 1, 2 or 3. This sets the t1 increment and the parameter **in0** is updated. Use the calculated value and set **inf1** in the **eda** window correspondingly. Since dipolar couplings between 1H and 15N can be up to 15 kHz, the spectral width (including the experiment scaling factor) should be min. around 20 kHz.

Parameter	Value	Comments
PULPROG	Pisema, pisemaclean	Pulse programs.
NUC1	^{15}N	
SW		Reasonable SW in F2.
O1P	90 – 160 ppm	For ^{15}N labeled acetylated glycine.
NUC2	^{1}H	
O2P	to be optimized	For ^{15}N labeled acetylated glycine.
PL1		For ^{15}N contact.
PL11		Or ^{15}N evolution.
PL2		For ^{1}H contact and excitation.
PL12		For ^{1}H heteronuclear decoupling during t2.
PL13		For ^{1}H Evolution under FLSG condition.
P3		^{1}H excitation pulse.
P15		^{15}N-^{1}H Contact pulse.
P6		^{1}H LG 294 degree pulse.
D3	1.4 µs	For frequency & phase setting D*X only.
cnst20		^{1}H spin nutation frequency achieved with PL13.
cnst21	0	Offset from o2 in Hz.
cnst22		+ LG frequency in Hz calculated.
cnst23		- LG frequency in Hz calculated.

L3	1 – 3	Loop counter for appropriate t1 increment.
F2 acquisition 15N	**********	Left column.
AQ_MOD	qsim	
TD	512	No of points.
DW		Dwell time in t2.
F1 indirect 1H	**********	Right column.
TD	64	Number of points.
IN_F1	L3*2*p5*SF or L3*2.5*p5*SF	Scaling factor for PISEMA 0.82 = sin Gy(54.7 deg.) calculated by pulse program.

Table 15.1: Acquisition Parameters

15.3 Processing

1. Process the direct dimension with **xf2**.
2. Accommodate for the **cos** modulated signal by setting the imaginary part to zero using the au program `zeroim` by typing **zeroim** on the command line.
3. Process the indirect dimension with the command **xf1**.
4. For more automated processing one can write a short macro using the command **edmac** and the file name *pisemaft* for examples: write the following commands using the text editor:
   xf2
   zeroim
   xf1
5. Save and close the **edmac** editor.
6. In the future you can then do the processing by simply typing **pisemaft** on the command line or even creating your own icon in TopSpin for this purpose.

Parameter	Value	Comment
F2 acquisition ^{1}H	**********	Left column.
SI	1k	Number of complex points in direct dimension.
WDW	no	Apodization in t2.
F1 indirect ^{15}N	**********	Right column.
SI	128	Number of complex points in indirect dimension.
MC2	QF	

Table 15.2: Processing Parameters for the Pisema Experiment

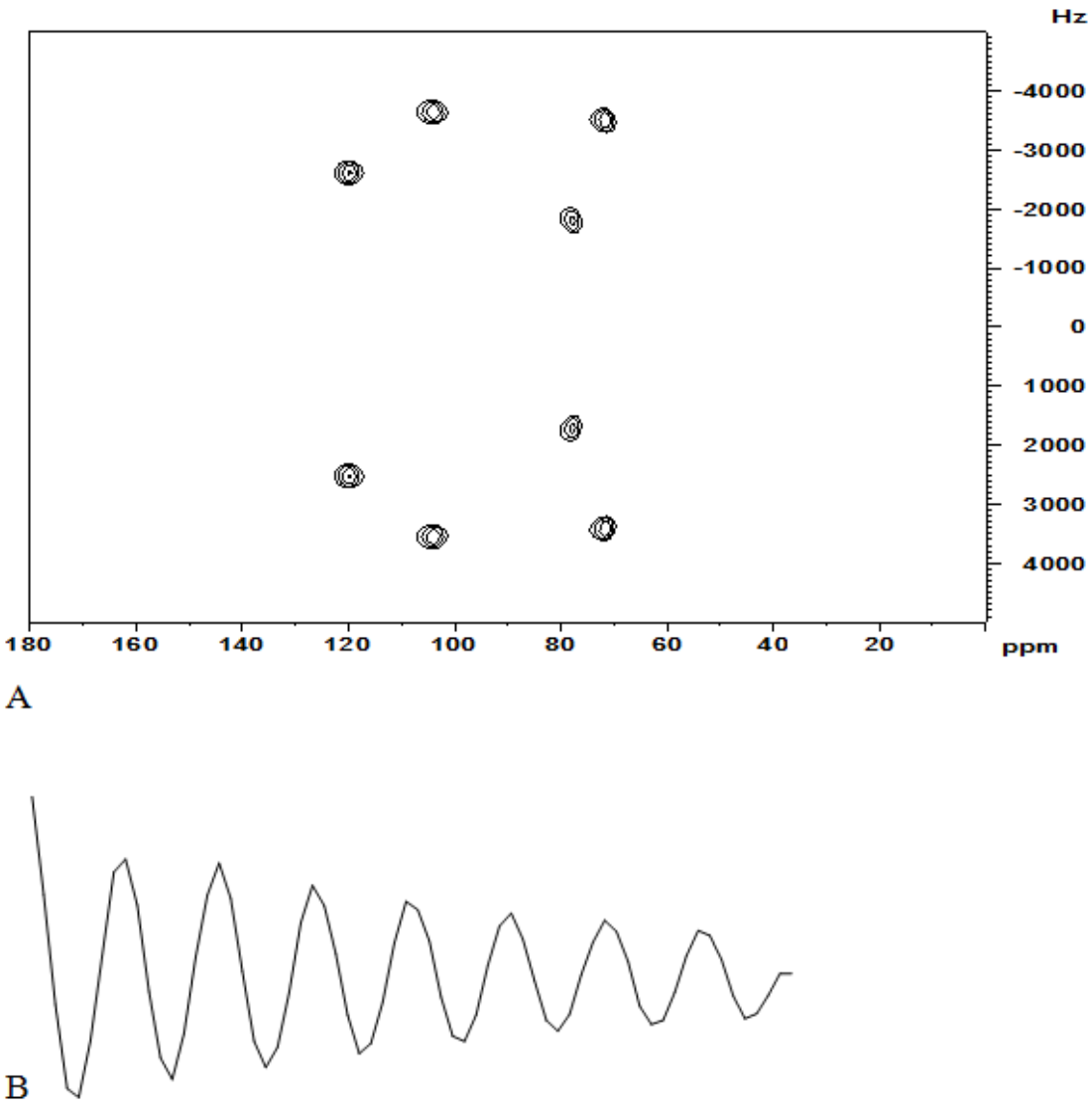


Figure 15.2: PISEMA Spectrum of 15N Labeled Acetylated Valine and FID in t1 over 3.008 ms 64 Data Points

In the figure above, „A“ is a PISEMA spectrum of ^{15}N labeled acetylated valine; and „B“ is a FID in t1 over 3.008 ms 64 data points.

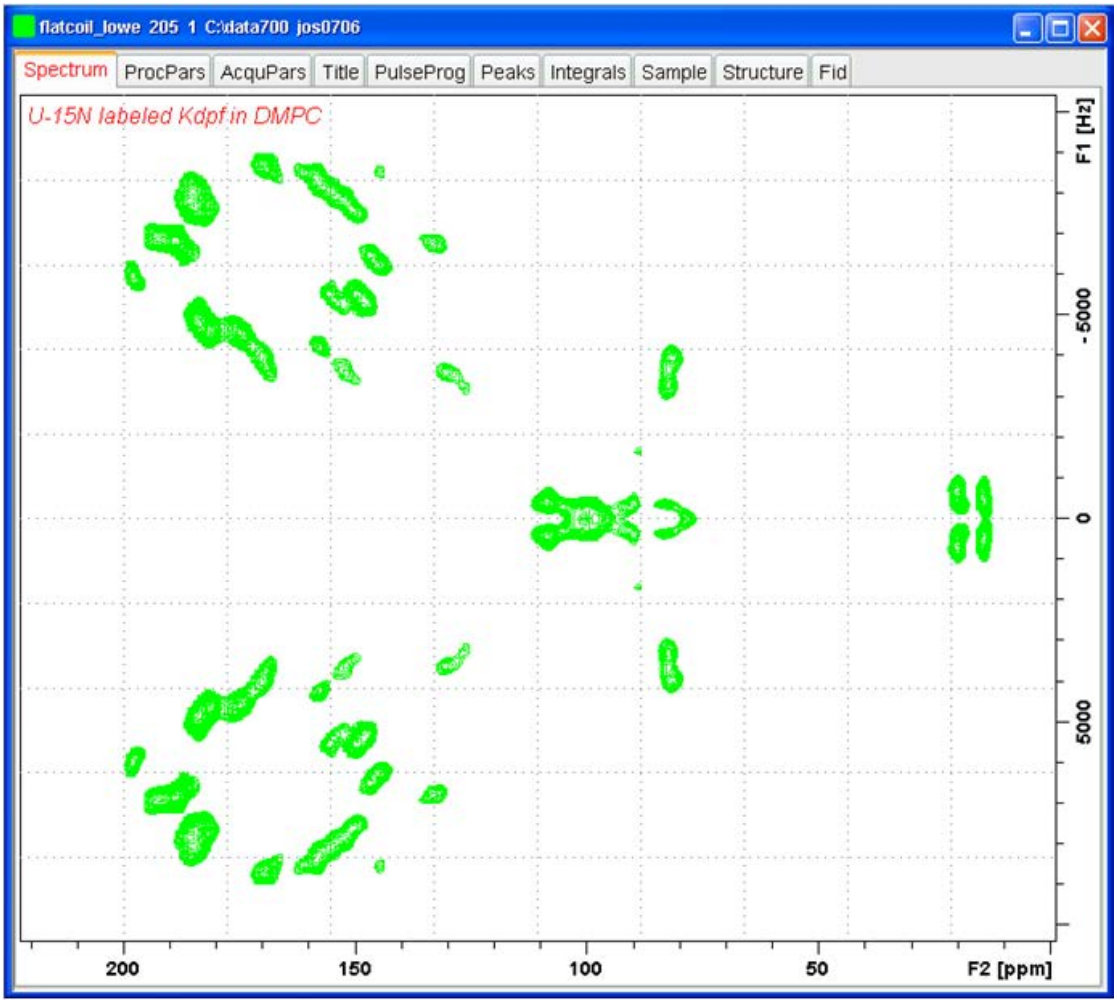


Figure 15.3: PISEMA Spectrum of 15N Labeled Kdpf Transmembrane Protein

The figure above is a PISEMA spectrum of ^{15}N labeled Kdpf transmembrane protein aligned in DMPC (courtesy of NHMFL, Dr. T. Cross) membrane between glass plates using an EFREE 700 MHz probe.

16 Relaxation Measurements

In NMR experiments, one is generally concerned with measuring resonance frequencies, and relating these to the local molecular environment. To do this the state of the system of spins in the sample must be changed from equilibrium. At equilibrium, the net magnetization due to the spins is aligned along the magnetic field axis. By applying a radio frequency pulse the net magnetization is tilted away from the field axis, and the resulting precessing magnetization generates the observed signal. The pulse has disturbed the system from equilibrium, and over time the system will return to its equilibrium state. This process is called relaxation.

This chapter describes experiments used for measuring relaxation rates in solid-state NMR. A basic description of relaxation is provided in order to define terms and introduce the techniques involved, but discussion of the significance and use of relaxation data is outside the scope of this manual. Many text books provide more detail on the theory of relaxation: the classic is Abragam:

- A. Abragam, *Principles of nuclear magnetism*, Oxford: Clarendon Press, (1961)

But simpler descriptions can be found in the books of Slichter and Levitt:

- C.P. Slichter, *Principles of magnetic resonance*, Springer (1996, 3rd ed.)
- M.H. Levitt, *Spin dynamics: Basics of nuclear magnetic resonance*, Wiley (2001)

Some discussion of T1r relaxation, including effects of dipolar coupling to proton spins, can be found in:

- D.L. VanderHart and A.N. Garroway, *13C NMR rotating frame relaxation in a solid with strongly coupled protons: polyethylene*, J. Chem. Phys., 71:2773-2787, 1979

Details of the X T1 experiment with CP are in:

- D.A. Torchia, *The measurement of proton-enhanced 13C T1 values by a method which suppresses artifacts*, J. Magn. Reson., 30:613-616, 1978

The TopSpin software includes a tool for processing the data obtained in relaxation measurements, and this will be demonstrated for the different types of relaxation experiment.

16.1 Describing Relaxation

Relaxation of the net magnetization can be described in terms of two processes. After a pulse, the state of the system differs from the equilibrium in two ways:

- The z-magnetization is not equal to the equilibrium value.
- The net magnetization in the transverse plane is non-zero.

The return of the z-magnetization to equilibrium is termed *longitudinal relaxation*, or *spin-lattice relaxation*, and the return of transverse magnetization to zero is termed *transverse* or *spin-spin relaxation*. Both the transverse magnetization and the difference between the current and equilibrium z-magnetization decay exponentially, with time constants denoted T_1 for longitudinal relaxation and T_2 for transverse relaxation. Relaxation also occurs while radio frequency pulses are being applied to the system. Normally this is ignored, but in the case of spin-locking pulses it is important. During cross-polarization, the magnetization on the dilute spins is increased by transfer from another nucleus, but it will also decay, since the radio frequency field (weak compared to the static field B_0) is insufficient to maintain the resulting transverse magnetization. If the pulse on the excitation nucleus is stopped, and only that on the detection nucleus continued, the transverse magnetization will decay exponentially, with a time constant denoted T1ρ. This rate of decay will be strongly affected by the amplitude of the spin-locking pulse.

Both of these processes occur via spin energy level transitions. It turns out that the spontaneous transition rate is very low, and thus relaxation is dominated by stimulated transitions. Such transitions are stimulated by local magnetic fields, which fluctuate due to local molecular motion, and the transition rates depend on the strength, and details of the fluctuations, of these local fields. Since the fluctuations are random, the rate of fluctuation is defined by the correlation time of the motion. For efficient relaxation via a particular energy level transition, fields fluctuating with an inverse correlation time close to the frequency of the transition are required. *Longitudinal relaxation* occurs via transitions on a single spin, and thus requires fields fluctuating with inverse correlation times near to the Larmor frequency. *Transverse relaxation* occurs also via flip-flop transitions of pairs of spins, which have energies close to zero, and so local fields fluctuating very slowly will cause transverse relaxation. $T_{1\rho}$ relaxation involves transitions at the nutation frequency of the spin-locking pulse, which can be selected by the experimenter. Measurement of these relaxation rates can therefore provide information about local motions on a range of time scales.

16.2 T1 Relaxation Measurements

Longitudinal relaxation can be measured using a number of methods – which method is appropriate depends on the sample involved. Here the experiments are demonstrated on glycine, which has a very simple spectrum and will give results using all the methods discussed. In general the only setup required is to calibrate pulses for the nucleus under observation, and to have some idea of the relaxation time constants involved.

16.2.1 Experimental Methods

The inversion-recovery method is the originally proposed method for measuring T_1 values. The experiment proceeds as follows: firstly, the magnetization is inverted by a 180° pulse. Then, there is a delay during which the magnetization relaxes, and a 90° pulse converts the remaining longitudinal magnetization to transverse magnetization, and an FID is recorded. The intensity of a particular signal in the resulting spectrum depends on the initial intensity, the relaxation delay, and the relaxation time constant T_1 as follows:

$$S(t) = S_E + (S(0) - S_E)\exp(-t / T_1)$$

Where *t* is the relaxation delay, S_E is the maximum signal seen when t is infinite, S*(0)* is the signal measure with no relaxation delay, and T_1 is the relaxation time constant for the spins giving rise to that signal. Measurement of S*(t)* for a number of relaxation delays allows determination of T_1.

The disadvantage of the inversion recovery experiment is that the delay between scans needs to be somewhat longer than the longest T_1 of the slowest relaxing spins in the sample. If cross-polarization from protons is possible, the initial inversion pulse can be replaced by a cross-polarization step followed by a 90° pulse on the nucleus to be observed. Then, the required delay between scans **d1** becomes that for relaxation of the protons. In most cases, the proton T_1 is moderate so inversion recovery (Torchia method) is the method of choice.

If the T_1 relaxation time is extremely long, the saturation-recovery experiment is preferred. Here, the transitions are saturated by a rapid sequence of hard pulses, such that no signal remains. There is then a variable delay, during which relaxation occurs, and then a 90° read-out pulse. If the relaxation delay is very short, no signal is seen, and at long relaxation times the maximum signal is seen. The advantage is that the saturation time required does not need to be many times the longest T_1 value. The state of the system at the start of the experiment is forced by the saturation pulses, so a long recycle delay is not required.

16.2.2 The CP Inversion Recovery Experiment

Sample: Glycine

Spinning speed: 10 kHz

Experiment time: 20 minutes

Before starting the experiment, the spectrometer should be set up as described in the chapter *Basic Setup Procedures [▷ 49]*, including measurement of the carbon pulse lengths, and the CP spectrum of glycine should be acquired for reference. Since relaxation times are necessarily temperature dependent, control of the sample temperature is desirable. The data shown here were all acquired at an approximate temperature of 20° C. The form of the pulse program is shown in the following figure.

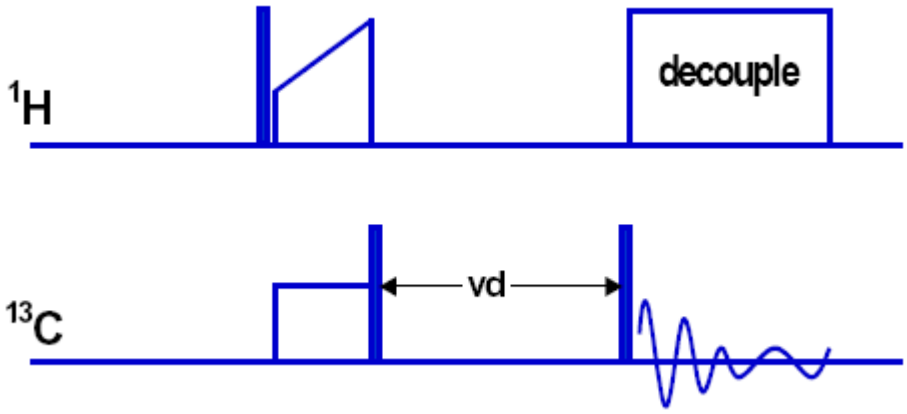


Figure 16.1: The CPX T1 Pulse Sequence

Starting from the glycine spectrum, create a new data set, set parameters according to the table below, and acquire a 1D spectrum. The relaxation delay after inversion is controlled by a variable delay list – this can be created using **edlist**, and the name of the list set as the parameter **vdlist**.

Parameter	Value	Comments
Pulprog	cpxt1	
Vdlist	See text	Relaxation delays after inversion pulse. Short value – to set spectrum phase correctly.
d1	3s	Needs only to be 3x proton T_1.
pl1	X HH contact power	Standard CP setting.
pl11	power for 90 degree pulses	Usually pl11<pl1 for short pulses.
p1	Measured 90° X pulse length at pl11	Adequate for required excitation bandwidth.
Ns	2	Should be enough to see a reasonable spectrum.

Table 16.1: Parameters for the 1D CP Inversion Recovery Experiment

This pulse program uses the method of Torchia, in which the phase of the contact pulse, and the receiver, is inverted in alternate scans. In the first scan, the first 90° pulse creates –z magnetization, and in the second scan it creates +z. The phase cycling of the receiver means that the difference between the two scans is recorded. For short relaxation delays, neither relaxes significantly, and so the maximum signal is recorded. At longer relaxation delays, both the +z magnetization (which is larger than the equilibrium value as it is created by CP), and the –z magnetization relax, and the recorded signal decays exponentially as a function of the relaxation delay. At long times both have relaxed back to equilibrium, and the two scans yield a zero signal.

The resulting spectrum should be phased to give positive peaks – given the very short recovery delay, no appreciable relaxation will have occurred. Set parameters for the 2D acquisition, as in the figure above. Since this is a pseudo-2D experiment, the only relevant parameter in F1 is the number of points, which should be the number of entries in the vd list. The most important setting is the range of relaxation delays set in the vd list. Ideally, the list should run from times short enough for no appreciable relaxation to occur, up to a few times the longest T_1 value. Of course, the accurate relaxation time constants are not known in advance, but order of magnitude estimates can be obtained by running the 2D experiment with a small number of relaxation delays, and a small number of scans per slice. The relaxation delays should be approximately equally spaced in log(delay), in order that decays with all time constants in the range are equally well characterized. Data can always be improved either by increasing the number of relaxation delays sampled, or by averaging more FIDs at each relaxation delay. For the glycine sample, a suitable list of times would be:

100 ms, 220 ms, 450 ms, 1 second, 2.2 seconds, 4.5 seconds, 10 seconds, 22 seconds, 45 seconds.

Parameter	Value	Comments
Parmode	2D	
Vdlist	See text	
td(f1)	Number of entries in vd list	
FnMODE	QF	This is not a real 2D experiment.
NS	4, for the glycine sample	Sample dependent – need to see a reasonable spectrum, but must be an even number.

Table 16.2: Parameters for 2D Inversion Recovery Experiment

16.2.3 Data Processing

Once the pseudo-2D data has been recorded, the processing parameters must be set and checked before it can be evaluated using the T_1/T_2 relaxation tool. The table below lists the relevant parameters. No processing is done in the indirect dimension (the relaxation dimension), but the size must still be set to a power of two for TopSpin to create a processed data file. The size should be next power of two larger than the number of relaxation delays used. The zero points appended are ignored by the relaxation analysis. In principle the line shape in the frequency dimension does not affect the analysis, so exponential multiplication with **lb** of the order of the observed line width can be applied to improve the signal-to-noise ratio.

Parameter	Value	Comment
F2 – acquisition dimension		
SI	=TD	Zero fill.
LB		Matched to line width.
WDW	EM	
Ft_mod	FQC	
ABSF1	1000	Limits for baseline correction.
ABSF2	-1000	Should cover entire F2 width.
F1 – relaxation dimension		
SI	Smallest power of 2 greater than TD(F1)	Must be 2n, but any zeros will be ignored.

Table 16.3: Processing Parameters for CP T1 Relaxation Experiment

Once the parameters are set, process the data with **xf2**, to execute a Fourier transform in the f2 dimension. The phase can be adjusted from within the relaxation analysis tool, but baseline correction should be carried out with **abs2**. Start the relaxation analysis guide with the command **t1guide**. The sequence of icons guides you through the analysis as follows:

Extract slice: The first spectrum row should be selected for phase correction, as this contains the maximum signal. The spectrum should then be phased to give positive peaks.

Define ranges: Here you must define integral regions containing the peaks of interest. The fitting routine can either use the integral of the signal, or the intensity, in which case the maximum signal in each integral region is used. Regions can be defined via the cursor, or between specified limits via a dialog box. The integral regions need to be saved to a special file, by clicking the disk icon towards the left of the integral window (not the standard save integrals button on the right), and selecting **export regions to relaxation module and.ret**.

Relaxation window: Here the intensity or area values from the first integral region are displayed. The icons at the top of this window allow you to move between the integral regions, exclude points from the calculation, display the data on a variety of axes, and start the fit for the displayed region or all regions. The figure below shows the decay of the a-carbon signal of glycine as a function of relaxation delay, along with the fit and calculated relaxation parameters. Note that any peaks with integrals or intensities too close to zero will be omitted from the analysis by the software – if you see less points in the relaxation window than were actually recorded, this may be because they have insufficient intensity.

Fitting function: Here the parameters of the fitting calculation are set. The general parameters should be determined automatically, but ensure that the limits for baseline correction are set to cover the whole spectrum. The fitting function depends on the experiment, but in this case the signals decay exponentially, so the function **expdec** should be selected. The list file name should be 'vd', this will take the specified **vdlist** from the data set. Note that when the experiment is run, the selected **vdlist** is written into the acquisition data directory as the file "vdlist", so it is always available, even if the source list is edited. The fitting program can calculate multi exponential fits, but data with very good signal-to-noise is required for this to be accurate. Unless there is obvious overlap of peaks, the assumption is usually that each peak corresponds to a single nuclear site, and thus a single T_1 value.

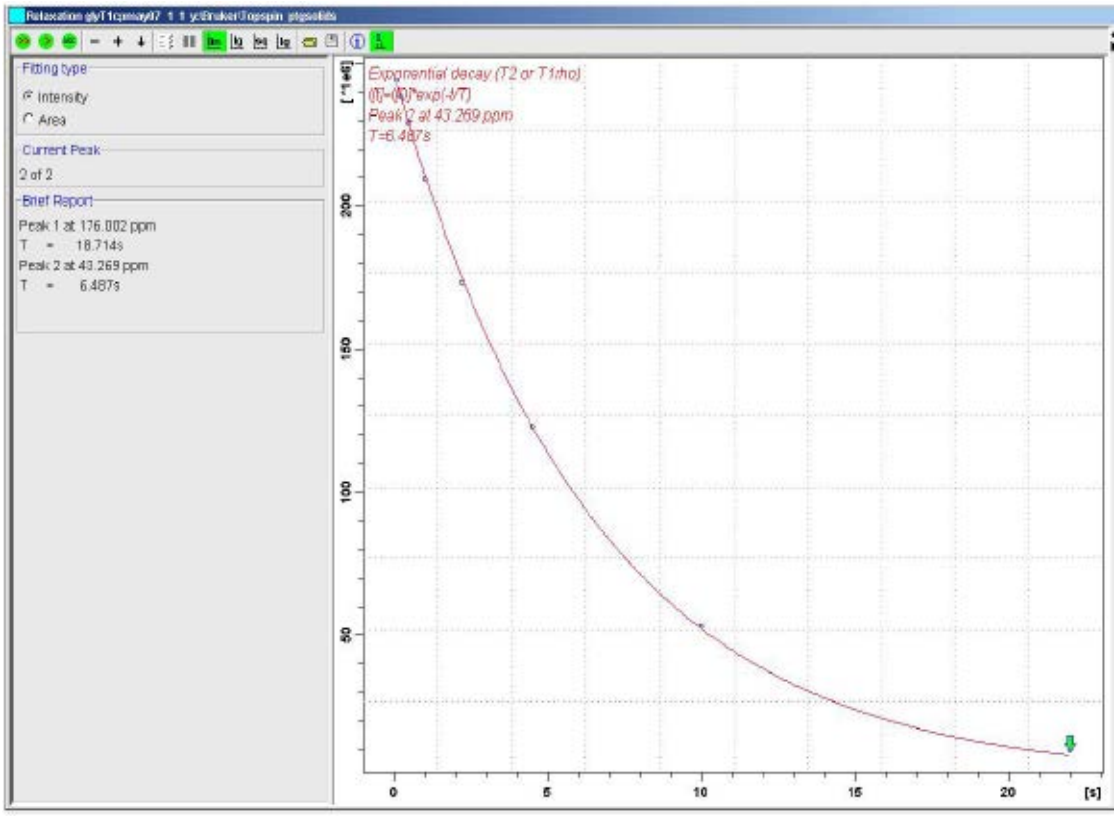


Figure 16.2: Relaxation of Alpha-carbon Signal in Glycine

Start calculation: This will perform the fitting procedure for all regions. The calculated function is displayed as a red line on the same axes as the data points. The plus and minus icons can be used to move through the different regions. If you wish to change whether the fit is based on the integral or the intensity, select the appropriate radio button, and repeat the fit using the icons immediately above. The **>>** icon will fit all the peaks, the **>** icon will fit just the current one.

Display report: This displays a text report of the results of the fit, including the details of the fit function, and the calculated values of the parameters in the function. The experimental and calculated data points are also displayed. Note that the experimental data is normalized such that the most intense point has a value of 1. This report file is also saved in the processed data directory when the fit is calculated. If fitting of a single peak is performed, only this result is written to the report. If the **fit all peaks** option is used, all results will be stored.

The results for glycine at 500 MHz and room temperature should be approximately 18.5 seconds and 6.4 seconds for the carbonyl and alpha-carbon signals respectively. At other field strengths the numbers will be somewhat different. If the signals are really undergoing mono exponential relaxation, the curve should be a good fit to the measured data.

16.2.4 The Saturation Recovery Experiment

For samples where cross-polarization is not possible, the inversion recovery experiment would be very time consuming, as the recycle delay **d1** would need to be approximately 3x the longest T1 value. For glycine at room temperature, this would mean a delay of about 60s per scan, in addition to the variable relaxation delay. The saturation-recovery experiment removes the need for long **d1** by forcing the system into saturation at the beginning of each scan.

Sample: Glycine

Spinning speed: 10 kHz

Time: 20 minutes

Experiment setup

Start with the standard carbon CP parameters, and set pulprog to **satrect1**.

Set **zgoptns** to –Ddec to turn on proton decoupling. If decoupling is not required on a real sample, this can be left blank to turn off the decoupling.

Set **p1** and **pl1** to the measured carbon 90-degree pulse parameters (as used in the CP T1 experiment, or see chapter *Basic Setup Procedures [▷ 49]*). Set **d1** to a relatively long value for the preparation experiments.

Set the number of pulses in the saturation train, **l20,** to zero, and acquire a spectrum. This will give an idea of the amount of signal, and thus how many scans need to be acquired for each relaxation delay.

Create a new data set with **iexpno**, and set the saturation parameters: **l20** = 5-100 and **d20** = 1-50 ms respectively.

Acquire a spectrum, and verify that saturation is complete – there should be no signal at all.

Setting up the 2D experiment

Set the parameters for the 2D acquisition as detailed in the table below.

For the variable recovery delay, the same values can be used as for the inversion recovery experiment.

The recycle delay **d1** can be very short, but take care not to exceed the duty cycle limits. High-power pulsing should not exceed 5% of the total scan time. In the case of decoupling, the acquisition time comprises most of the pulsing, so **d1** should be >20x **aq** – 1 second is reasonable in this case. If the experiment is run without decoupling, then the saturation period is the only significant period of high-power pulsing, and **d1** can be shorter.

Acquire the 2D spectrum with **zg**.

Parameter	Value	Comments
Parmode	2D	
Vdlist	See text	
td(f1)	Number of entries in vd list	
FnMODE	QF	This is not a real 2D experiment.
NS	16, for the glycine sample	More scans needed than for the CP experiment, due to reduced signal.

Table 16.4: Parameters for the Saturation Recovery Experiment

Data Processing

The saturation-recovery data should be processed in the same way as the inversion-recovery data above (see *Table 16..3 [▷ 185]*). The only differences are that the fitting function should be **satrec** rather than **expdec**, and the slice selected for processing should be the last one (signal is maximum at long recovery times). The calculated relaxation time constants should be the same as those obtained by inversion-recovery.

16.2.5 T1p Relaxation Measurements

Rotating-frame relaxation measurements, under a spin-locking RF field, can be used to probe motions shorter domestically than T_1 measurements, with inverse correlation times of the order of the spin-locking RF field strength.

To measure T_{1r} relaxation after CP, a variable length spin-locking pulse is applied to the X nucleus. The remaining X magnetization decays exponentially to zero, as a function of spin-lock time. The parameters of cross-polarization can also be determined from variable-contact-time CP experiments (the function cpt1rho is provided in the relaxation analysis tool for this purpose), but here only simple T_{1r} measurements will be discussed.

It should be noted that the relaxation in a T_{1r} experiment might result from processes other than true T_{1r} relaxation. For example, in glycine, the carbon spins are dipolar coupled to protons, and there is a possible fast relaxation pathway via the protons, which is not T_{1r} relaxation. This is inhibited by having a high spin-lock field strength, but at large field strengths care must be taken over the length of the spin-lock pulse. If apparently non exponential decay is observed, this may result from such alternative relaxation processes.

Experiment setup

Sample: Glycine

Spinning speed: 10 kHz

Time: 20 minutes

Start from standard CP parameters. The only additional calibration required is the carbon RF field strength of the spin-lock pulse. This can be set independently of the field strength for the cross-polarization. In principle, the strength of this field can be set to any value (within probe limits) to probe motions on a range of time scales. However, only at relatively large field strengths is true T_{1r} relaxation the only significant relaxation pathway.

Set **pulprog** to **cp90** and measure the required power level for a 70 kHz RF field (3.57 µs 90 degree pulse).

Make a new data set with **iexpno**, change **pulprog** to **cpxt1rho** and set this measured power as **pl11**.

Set up a variable pulse list for the incrementation of the spin-lock, with the command **edlist vp**. Check that this list is set as the parameter **vplist**. Remember that this is a high-power pulse, so the duration should not be too long. For the glycine sample, a possible set of times would be: 1 ms, 2 ms, 5 ms, 10 ms, 15 ms, 20 ms, 25 ms, 30 ms, 40 ms, 50 ms. This will not allow the signals to decay completely, so is not ideal, but should not place undue stress on the probe. Often a compromise must be reached between recording an ideal decay curve and avoiding the risk of probe damage.

Change **parmode** to 2D and set other 2D parameters as for the other relaxation experiments.

Acquire spectrum with **zg**.

Data Processing

The data can be processed in the same way as the other relaxation experiments. The slice with the shortest spin-lock time contains the most signal, so this slice should be used for processing. The fitting function should be set to **expdec**, and **vplist** should be selected as the list file name, in the fitting function dialog.

At 500 MHz, with a 60 kHz spin-lock field, the $T_{1\rho}$ values should be approximately 400 ms and 48 ms for the carbonyl and alpha carbons respectively. The data for the alpha carbon does not give a perfect fit to a single exponential, but this may result from the relatively low spin-lock field allowing non-$T_{1\rho}$ relaxation.

16.3 Indirect Relaxation Measurements

If proton relaxation measurements are desired, the considerable broadening of the proton resonances seen at even high spinning speeds can make resolution of individual components impossible. In such cases, indirect observation of proton relaxation by X-nucleus observation can be used. A typical example would be attempting to observe the proton relaxation of two components of a mixture or multi phase material. In general, the proton spins within a single molecule are sufficiently stronger coupled by the homonuclear dipolar coupling so that different relaxation is not seen for the different sites. If the experiments are set up with short contact times, the individual carbon signals will be derived only from directly bonded protons, and thus any differences in proton relaxation within a molecule can be isolated.

Such indirect observation can be implemented conveniently for both T_1 and $T_{1\rho}$ relaxation. For T_1, a proton saturation-recovery step is inserted prior to the cross-polarization step in a standard CP sequence. The proton magnetization occurs immediately prior to CP, and thus the observed carbon signal, depends on the extent of recovery after the saturation, so the carbon signal as a function of recovery delay gives the proton T_1 value. For $T_{1\rho}$ a variable length proton only spin-lock pulse is applied after the 90-degree pulse in the CP experiment. The proton magnetization after this pulse, and thus the carbon signal after CP, depends on the proton $T_{1\rho}$ relaxation.

16.3.1 Indirect Proton T1 Measurements

Sample: Glycine

Spinning speed: 10 kHz

Time: 20 minutes

Start with the standard CP parameters, as for the X T_1 measurement with CP, and set **pulprog** to *cph+1*. Set the saturation loop **l20** to zero, and acquire a spectrum, to check signal intensity. Signal to noise should be comparable with the standard CP experiment, so a number of scans similar to that used for the carbon T_1 experiment should suffice.

Saturation parameters can be set like the carbon saturation were previously set: **l20** = 5-100 and **d20** = 1-50 ms. Acquire a spectrum with these parameters and verify that there is again no signal.

Make a new data set with **iexpno** and set parameters for 2D acquisition, as for the previous experiments. D1 can be short, with the same proviso about duty cycle as the X saturation-recovery experiment. A reasonable set of delays for the vdlist would be: 10 ms, 22 ms, 45 ms, 100 ms, 220 ms, 450 ms, 1 second, 2.2 seconds, 4.5 seconds, 10 seconds.

Data processing

The data should be processed in the same way as for the X saturation recovery experiment. Both the carbonyl and alpha-carbon peaks derive their carbon polarization from the same proton spins, and so analysis of the two peaks should give the same result. If you have a sample containing some gamma-glycine (gives peaks at slightly lower shifts than the more common alpha-glycine form), this should show different T_1 values for the two sets of peaks.

At 500 MHz, the proton relaxation time should be approximately 520 ms at room temperature.

17 Basic MQ-MAS

17.1 Introduction

The MQ-MAS experiment for half integer quadrupole nuclei is a 2D experiment to separate anisotropic interactions from isotropic interactions. In the NMR of half integer quadrupole nuclei the dominant anisotropic broadening of the central +1/2 <-> –1/2 transition (CT), and symmetric multiple-quantum (MQ) transitions, is the 2nd order quadrupole interaction which can only partially be averaged by MAS. The satellite transitions (ST, e.g. the ±3/2 <-> ±1/2 transitions) however, are broadened by a 1st order interaction, which is several orders of magnitude larger than the 2nd order broadening. Under MAS the 1st order interaction of the ST can be averaged but since the spinning cannot be fast compared to the first order broadening (of the order of MHz), a large manifold of spinning side bands remains. The 2nd order broadening of the CT can only be narrowed by a factor of 3 to 4 by MAS, so a signal is observed that still reflects this 2nd order broadening, which may be of the order of kHz. Line shapes resulting from nuclei in different environments are thus likely to be unresolved in a simple 1D spectrum.

The 2D MQ-MAS experiment exploits the fact that the 2nd order broadening of the symmetric MQ transitions (e.g. +3/2 «-3/2 in a spin 3/2), is related to the 2nd order broadening of the CT by a simple ratio. A 2D spectrum is recorded which correlates e.g. a +3/2 «–3/2 3Q coherence involving the satellite transitions and the +1/2 «–1/2 single quantum coherence of the central transition. This spectrum shows a ridge line shape for each site, with slope given by the ratio of the second order broadening of the two transitions (-7/9 in the case of the 3Q transition). A projection of the 2D spectrum perpendicular to this slope yields an isotropic spectrum free from quadrupolar broadening.

17.2 Pulse sequences

The figures in this section show two of the basic sequences, a 3-pulse and a 4-pulse sequence with z-filter. Both sequences start with an excitation pulse **p1** that creates 3Q coherence which is allowed to evolve during the evolution period **d0**. In the 3-pulse sequence the subsequent conversion pulse **p2** flips magnetization back along the z-axis, which after a short delay **d4** (to allow dephasing of undesired coherency) is read out with a weak CT selective 90° pulse **p3**. In the 4-pulse sequence, however, the conversion pulse **p2** changes 3Q coherency to 1Q coherency which then passes through a Z-filter of two CT selective 90° pulses in a **p3-d4-p3** sequence.

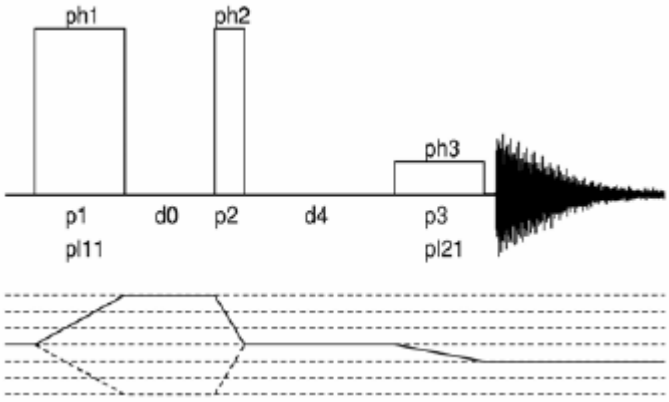


Figure 17.1: A 3-Pulse Basic Sequence with Z-Filter

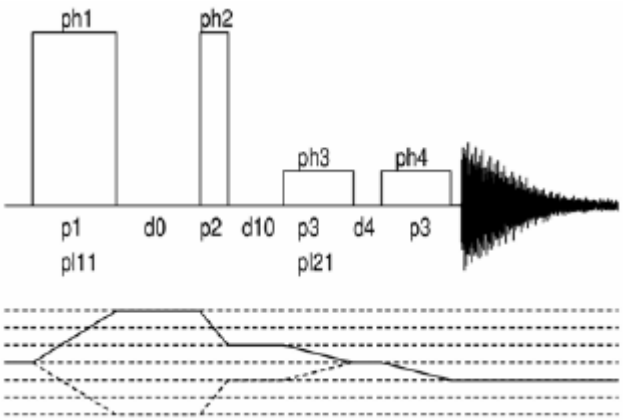
Three pulse sequence and coherence transfer pathway for the 3Q MAS experiment with z-filter (`mp3qzqf.av`). The ratio for pulses **p1** and **p2** is approximately 3. The corresponding power level **pl11** should be set to achieve at least 150 kHz RF field amplitude. **p3** should be some tens of µs, corresponding to an RF field amplitude of a few kHz. Delays **d0** and **d4** are the incremented delay for **t1** evolution and 20 µs for z-filter, respectively. Phase lists are as follows, for phase sensitive detection in F1 the phase of the first pulse must be incremented by 30° in States or States-TPPI mode:

ph1 = 0

ph2 = 0 0 60 60 120 120 180 180 240 240 300 300

ph3 = 0 180

receiver = + – – +.



A 4-Pulse Basic Sequence with Z-Filter.

Four pulse sequence and coherence transfer pathway for the 3Q MAS experiment with z-filter (`mp3qzfil.av`). The pulse **p1** is the same, and **p2** is usually somewhat shorter than in the three pulse sequence. The corresponding power level **pl11** should be set to achieve at least 150 kHz RF field amplitude. The pulse **p3** should be some tens of µs, corresponding to an RF field amplitude of a few kHz. Delays **d0** and **d4** are the incremented delay for t1 evolution and 20 µs for z-filter, respectively. Delay d10 initially is 0 and can be incremented proportional to **d0** (**in10** = **in0** *7/9), if the observe nucleus has spin I=3/2. Phase lists are as follows, for phase sensitive detection in F1 the phase of the first pulse must be incremented by 30° in States or States-TPPI mode:

```
ph1 = 0 60 120 180 240 300
ph2 = 0*24 90*24 180*24 270*24
ph3 = 0
ph4 = 0*6 90*6 180*6 270*6
receiver = {0 180}*3 {90 270}*3 {180 0}*3 {270 90}*3 {180 0}*3 {270
90}*3{0 180}*3 {90 270}*3.
```

The sequence with more pulses has slightly inferior sensitivity; however, it is the basic sequence to improve sensitivity by FAM or DFS. The 3-pulse sequence itself can be used directly to enhance sensitivity by soft-pulse added mixing (pulse program `mp3qspam.av`). In the chapter *MQ-MAS: Sensitivity Enhancement [▷ 207]* some of the sensitivity enhancement techniques will be described.

Note that pulse programs suitable for AV and AVII spectrometers have the `extension.av`, pulse programs for the AVIII have no extension.

17.3 Data Acquisition

Before the 2D experiment on your sample of interest can be started, two set-up steps must be done as described in detail in the next section. All set-up steps should be done on a sample with:

- A known MAS spectrum.
- With sufficiently good sensitivity to facilitate the set-up.
- A 2nd order quadrupole interaction in the order of the one expected for your sample of interest.

In the first step a low power selective pulse must be calibrated in a single pulse experiment. With this the MQ-MAS experiment can be optimized using the 2D pulse sequence for t_1=0.

17.3.1 Setting Up the Experiment

Sample: There are a large number of crystalline compounds which can be used to set up the experiment. Please refer to the table below to select a suitable sample. For the general procedure described here the spin I of the nucleus is not important, and pulse widths obtained will depend on the spin I and the Larmor frequency. You can use any arbitrary sample showing a considerable broadening by the 2nd order quadrupole interaction to adjust the experiment, however, reasonable 1D MAS spectra should be obtained quickly for sensitivity reasons.

The set-up must be done in two steps; in the first step a central transition selective pulse that merely excites the central transition must be calibrated. This pulse must be weak enough so that only this transition is affected and it must be short enough so that the central transitions of all sites in the spectral range are excited. As an example, the sinc shape excitation profile of a 20 µs pulse has its zero-crossings at 1/20 µs = ± 25 kHz which means that the central transition signals must not extend beyond this range, otherwise severe line shape distortions will be observed. On the other hand the corresponding RF field amplitude of a 20 µs 90° pulse will be 1/(80 µs*(I+1/2))=12.5 kHz/(I+1/2). This means that $w_{RF} << w_Q$ as a prerequisite for a CT selective pulse is most likely to be fulfilled. For the calibration of this pulse a power level around 30 dB with 500 W and 1 kW amplifiers and around 20 dB with 300 W amplifiers should be expected. The pulse program `zg` (which uses **p1** and **pl1**) or `zgsel.av` (which uses **P3** and **PL21**) can be used.

Nucleus	Spin	Spectrometer Frequency*1)	d1 [s]*3)	Sample	Comments
17O	5/2	67.78	2	NaPO3	> 10% enriched
11B	3/2	160.42	>5	H3BO3	
23Na	3/2	132.29	10	Na2HPO4*2)	
27Al	5/2	130.32	5	YAG	
27Al	5/2	130.32	0.5	Al2O3	
27Al	5/2	130.32	0.5	VPI-5	
27Al	5/2	130.32	0.5	AlPO4-14	
11B	3/2	160.46	5	H3BO3	
87Rb	3/2	163.61	0.5	RbNO3	
93Nb	9/2	122.25	1	LiNbO3	

*1) *In MHz at 11.7 T (i.e. 500.13 MHz proton frequency).*

*2) Alternatively Na_2HPO_4 * $2H_2O$ can be used. For anhydrous Na_2HPO_4 the sample should be dried at 70° C for a couple of hours before packing the rotor in order to eliminate crystal water completely.

*3) Recycle delays at 11.7 T, longer delays may be required at higher fields.

Table 17.1: Some Useful Samples for Half-integer Spin Nuclei

The figure below shows a comparison of a spectrum excited by a short non-selective pulse with a spectrum that has been obtained by a weak selective pulse. Note that in the latter the spinning sidebands from the satellite transition are no longer visible which is used as an indication that it is not excited.

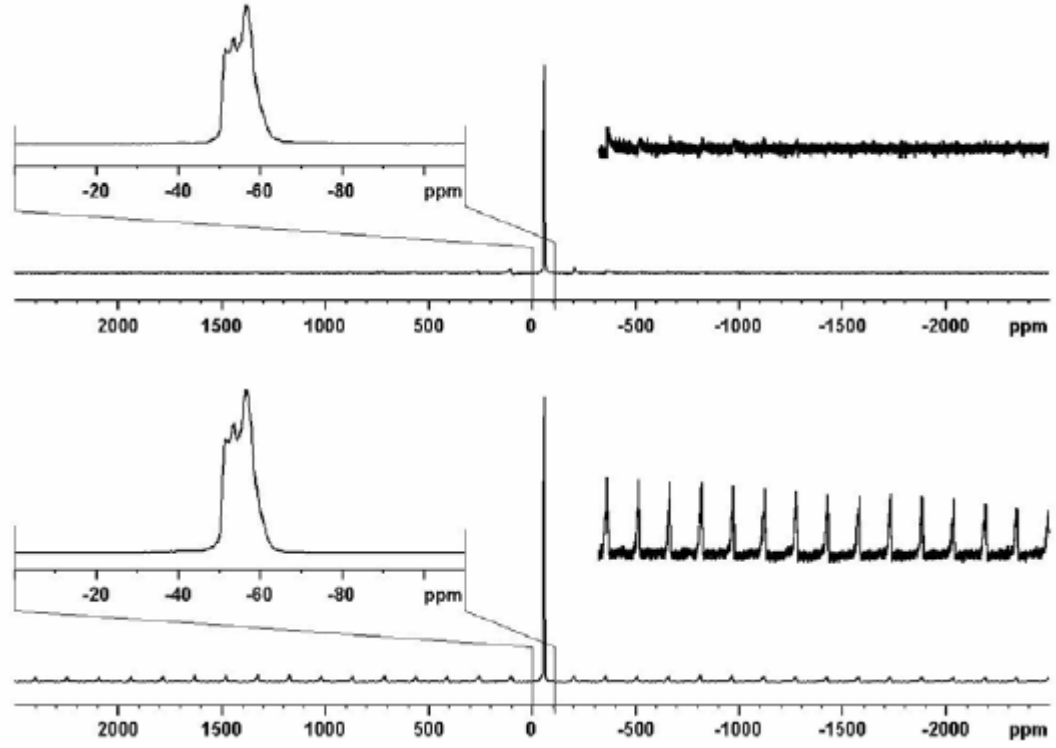


Figure 17.2: Comparison of 87Rb MAS spectra of RbNO3 excited with selective and non-selective pulses

The lower trace is a spectrum excited with a 1 µs non-selective pulse corresponding to a small flip angle. Above is a spectrum excited with a 20 µs selective 90° pulse. Note that in the latter no spinning side bands from the satellite transition are observed. Spectra are taken on AV500WB at a Larmor frequency of 163.6 MHz with a 2.5 mm CP-MAS probe spinning at 25 kHz.

The figure below shows the nutation profiles of a non-selective and a selective pulse, respectively. Note that for the selective pulse a fairly precise 180° pulse of a length of $2*_{\tau 90°}$ can be determined whereas for a non-selective pulse this is not the case.

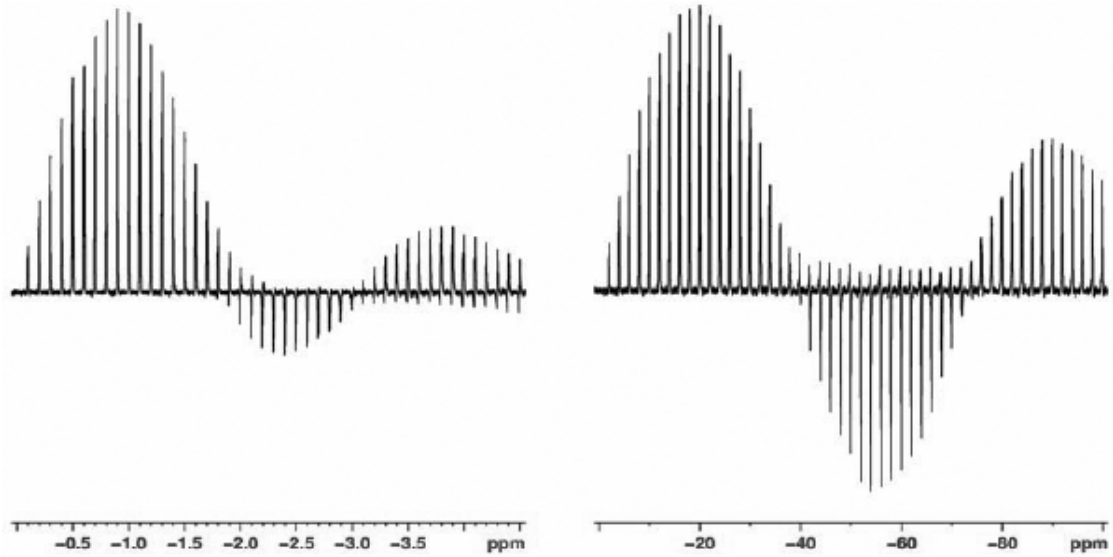


Figure 17.3: Nutation profiles of selective and non-selective pulses

The left diagram shows signal intensity of 87Rb resonances in RbNO3 as a function of a non-selective pulse at approx. 150 W RF power, the right diagram shows the signal intensity as function of a selective pulse at less than approx. 0.5 W. Spectra are taken on an AV500WB at a Larmor frequency of 163.6 MHz with a 2.5 mm CP-MAS probe spinning at 25 kHz. Note the different scaling of the x-axis, which is displayed as “ppm”, but corresponds to the used pulse lengths in µs (apart from the sign).

Once the central transition selective 90° pulse is calibrated the parameters can be copied to a new data set with **iexpno**, and the MQ-MAS pulse program can be loaded. Available pulse programs are **mp3qzqf** and **mp3qzfil**. The first is a 3-pulse sequence, the second a 4-pulse sequence. The sequence with fewer pulses will be slightly more sensitive, while the 4-pulse sequence can be used as an initial set-up for experiments with sensitivity enhancement methods like DFS or FAM (see the chapter *MQ-MAS: Sensitivity Enhancement [▶ 207]* which describes sensitivity enhancement methods).

The table below gives typical values for the pulses and powers that should be close to the final values confirmed by the optimization procedure. Parameters like **O1**, **TD**, **SWH**, **RG**, should already be set in the standard 1D spectrum. For 4 mm probes these pulse lengths are about the limit of what can be achieved, for 2.5 mm probes somewhat shorter pulses can be obtained. For I = 3/2 and I = 5/2 nuclei the ratio of p1/p2 ≈ 3.

For **pl11** an initial value that corresponds roughly to 300 W can be used. Optimization will be done on the first increment of the 2D sequence, i.e. **d0** = 1 µs. Two strategies for the optimization procedure can be followed; either the pulse lengths **p1** and **p2**, or the power level **pl11** can be optimized for maximum signal amplitude. However, the latter can be disadvantageous because a power level above the probe limit might be applied, in order to clearly determine the optimum power. In the case of 300 W amplifiers the maximum signal amplitude may not be obtained even at full power, with the chosen pulse lengths.

Parameter	Value	Comments
Pulprog	mp3qzqf.av or mp3qzfil.av	Pulse program.
NS	12*n (zqf) 96*n (zfil)	Full phase cycle is important.
D0	1u	Or longer, t1-period.
D1	5 * T_1	Recycle delay, use dummy scans if shorter.
D4	20 µs	Z-filter delay.
P1	3.6 µs	Excitation pulse at pl11.
P2	1.2 µs	Conversion pulse at pl11.
P3	20 µs	90° selective pulse at pl21 taken from previous pulse calibration.
PL1	=120 dB	Not used.
PL11	start with ≈ 300 W	Power level for excitation and conversion pulses.
PL21		Power level for selective pulse, approx. pl11+30 dB taken from previous pulse calibration.

Table 17.2: Initial Parameters for Setup

Hence, it is always better to optimize the pulse lengths **p1** and **p2**. In this case **p2** should be optimized before **p1** because the signal intensity is much more sensitive to this pulse length. A suitable set-up for the parameter optimization procedure **popt** is shown in following figure.

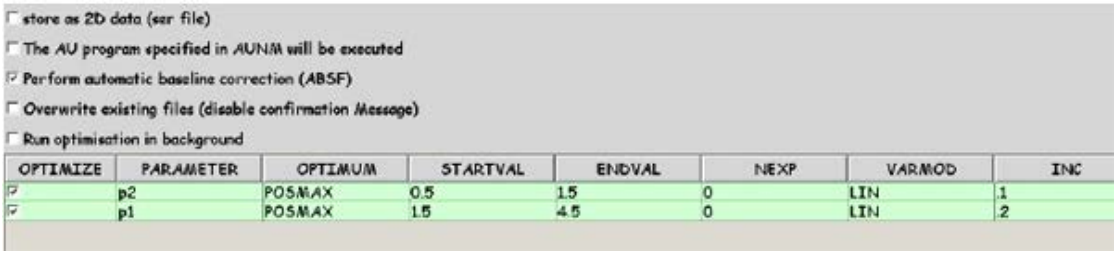


Figure 17.4: Example for popt Set-up for Optimization of p1 and p2.

In the first step **p2** is optimized to which the experiment is the more sensitive. In the second step p1 is optimized using the optimum value found for **p2** in the first step.

For more details about using the **popt** procedure to optimize a series of parameters please refer to the manual. The following figure shows the signal amplitudes as functions of pulse lengths **p2** and **p1**.

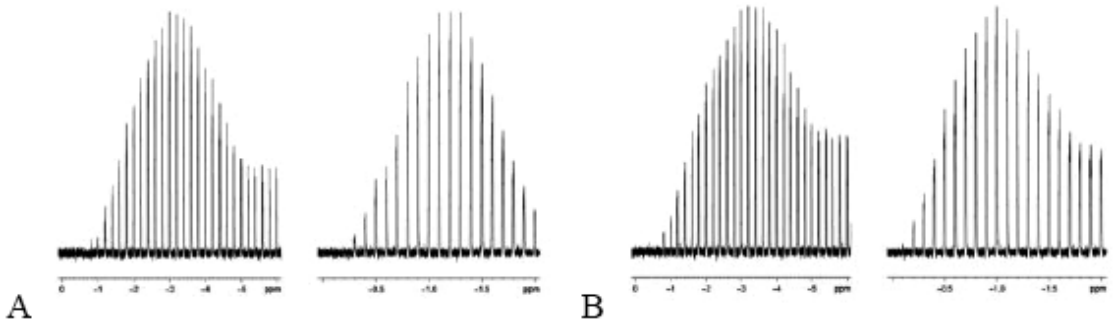


Figure 17.5: Signal Intensities of 87Rb Resonances in RbNO3 as Function of p1 and p2.

Each pair of diagrams in A and B shows the signal intensities as function of the excitation pulse **p1** and the conversion pulse **p2**. In A the 3-pulse sequence and in B the 4-pulse sequence was used. Note that the signal intensity is much more sensitive to the proper length of the conversion pulse. Maximum intensities were 3.0 µs and 1.2 µs in A and 3.2 µs and 1.0 µs in B, respectively. This corresponds to approximate RF field amplitudes of 160 kHz. Spectra are taken on an AV500WB at a Larmor frequency of 163.6 MHz with a 2.5 mm CP-MAS probe spinning at 25 kHz. Note the different scaling of the x-axes, for **p1** they range from 0 to 6 µs, for **p2** from 0 to 2 µs.

17.3.2 Two Dimensional Data Acquisition

Once the pulses are calibrated everything is ready for the 2D data acquisition. Create a new data set and change *parmode* to 2D. In the acquisition parameters for the (new) indirect F1 dimension the following parameters must be set according to the following table.

Parameter	Value	Comments
FnMode	States or States-TPPI	Acquisition mode for 2D.
TD	see text	Number of FID's to be acquired.
SWH	"masr"	Equals spinning frequency for rotor synchronization, from this IN_010 is calculated correctly, if ND_010 is already set.
NUC1		Select the same nucleus as for F2 so that transmitter frequency offset is correctly set (important for referencing).[1]
D10	0	Used in mp3qzfil.av only.
IN10	=in0*7/9	Used in mp3qzfil.av for nuclei with spin I=3/2 only, so that no shearing FT is required.
[1] Note the difference in increment handling in TopSpin 2.1 and higher.		

Table 17.3: F1 Parameters for 2D Acquisition

Some further comments and explanations on the parameters listed above:

FnMode must be States or States-TPPI so that the shearing FT can be performed for processing. If the pulse program `mp3qzfil.av` is used, no shearing is required in case of nuclei with spin I=3/2 if **in10** is set correctly in which case a split-t_1 experiment is performed. The parameter **td** determines the number of FID's to be accumulated in the indirect dimension. This value is determined by the line width and resolution that can be expected in the indirect MQ dimension (F1) and which depend on the properties of the sample. In crystalline material fairly narrow peaks can be expected so a maximum acquisition time in F1 of 2 to 5 ms is expected. In disordered materials where the line width is broader and determined by chemical shift distribution a total acquisition time in F1 of 1 ms may be sufficient. The total acquisition time **aq** in F1 equals (**td**/2)* **in_010**.

For rotor synchronized experiments **in_010** = 1/spinning frequency so will typically be between 100 µs (10 kHz spinning) and 28.5 µs (35 kHz spinning), so only 10 to 40 experiments in amorphous samples but 50 to 200 experiments in crystalline samples might be required. The rotor synchronization means that the spectral range in F1 is limited. Depending on the chemical shift range, spinning frequency, and quadrupole interactions the positions of the peaks may fall outside this range. In such a case care must be taken when interpreting the spectrum. Acquisition with half-rotor synchronization to double the spectral window in F1 may help. However, in this situation one set of spinning sidebands appears and it must be avoided that the spinning side bands of one peak fall on top of other peaks. Some sort of rotor synchronization is always recommended because spinning side bands in the indirect dimension extend over a very wide range, which cannot be truncated by e.g. filtering.

Therefore, rotor synchronization together with States or States-TPPI phase sensitive acquisition helps to fold spinning sidebands from outside back onto center bands or other side bands.

17.4 Data Processing

Processing parameters should be set according to the following table:

Parameter	Value	Comments
F2 (acquisition dimension)		
SI		Usually set to one times zero filling.
WDW	no	Don’t use window function.
PH_mod	pk	Apply phase correction.
BC_mod	no	No DC correction is required after full phase cycle.
ABSF1	1000 ppm	Should be outside the observed spectral width.
ABSF2	-1000 ppm	Should be outside the observed spectral width.
STSR	0	Avoid strip FT.
STSI	0	Avoid strip FT.
TDoff	0	Avoid left shifts or right shifts before FT.
F1 (indirect dimension)		
SI	256	Sufficient in most cases.
WDW	no	Don’t use window function, unless F1 FID is truncated.
PH_mod	pk	Apply phase correction.
BC_mod	no	No DC correction is required after full phase cycle.
ABSF1	1000	Should be outside the observed spectral width.
ABSF2	-1000	Should be outside the observed spectral width.
STSR	0	Avoid strip FT.
STSI	0	Avoid strip FT.
TDoff	0	Avoid left shifts or right shifts before FT.

Table 17.4: Processing Parameters for 2D FT

Data obtained with `mp3qzfil.av` from nuclei with spin I = 3/2 can be processed with **xfb**, if **IN10** has been set appropriately to run a split-t1 experiment. If this is not the case data can be sheared in order to align the anisotropic axis along the F2 axis. This is done with the AU program `xfshear`. The AU program checks the nucleus to determine the spin quantum number, checks the name of the pulse program and decides what type of experiment has been performed. In case the nucleus is unknown to the program, or the pulse program has a name that does not contain a string **nq** nor **nQ** (with n=3, 5, 7, 9), the required information is asked for by the program, in order to calculate the shearing correctly.

Note that using a user designed pulse program that contains e.g. a string **5q** but performs a 3Q experiment (and vice versa) will yield erroneous shearing. When started, the AU program prompts for **Apply ABS2?** and **F1 shift in ppm**. It is advisable to calculate a baseline correction after F2 Fourier transform. Note that the range defined by ABSF1 and ABSF2 is used for this. You should make sure that the limits are at least as large as the spectral width to allow baseline correction of the whole spectrum. The **F1 shift in ppm** allows shifting the spectrum (including its axis) in the vertical direction for cases where peaks are folded due to a limited spectral window in a rotor synchronized experiment. For the first processing both prompts are typically returned. At the end of the processing the AU program corrects the apparent spectrometer frequency of the indirect dimension by a factor |R-p|, where R is defined in equation [1] and p is the order of the experiment (e.g. 3 for 3QMAS):

$$R = \frac{m(18I(I+1) - 8.5m^2 - 5)}{18I(I+1) - 3.5}$$

This ratio is calculated from the spin quantum number I of the nucleus and the magnetic spin quantum number m, which is determined by the experiment, e.g. 3/2 in case of a 3Q experiment of an order p=3. The program stores the “F1 shift” that was calculated and will prompt for it when data are processed next time. If the same F1 shift should be applied as before the AU program can be called with the option **lastf1**. Before giving some further explanations about the experiment, shows the 2D ^{87}Rb 3QMAS spectrum of $RbNO_3$.

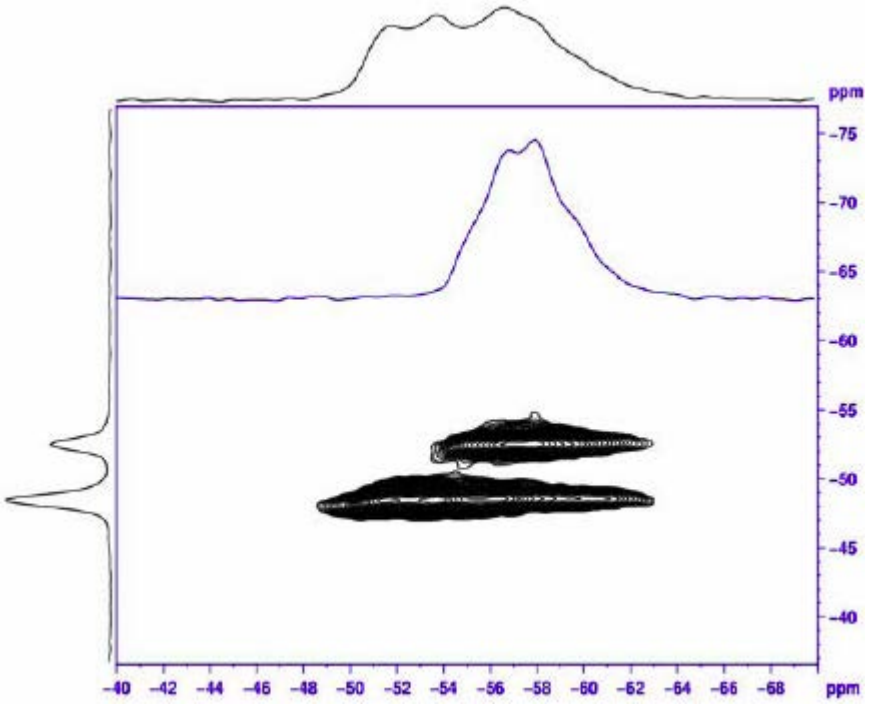


Figure 17.6: 2D 87Rb 3QMAS Spectrum of RbNO3

The top and left projections are the summations over the signal ranges. The spectrum included in the 2D map is a cross section through the resolved peak resonating at approximately 53 ppm. Note that at 11.7 T two of the three sites cannot be resolved in the 2D spectrum. The spectral range shown in F1 corresponds to the spinning frequency. Spectra are taken on an AV500WB at a Larmor frequency of 163.6 MHz with a 2.5 mm CP-MAS probe, spinning at 25 kHz.

Since the quadrupole parameters are usually unknown before performing the experiment the positions of the peaks in the indirect dimension cannot be predicted. Therefore, it may happen that a peak is positioned at the border of the spectral range in the F1 dimension or even folded. When using `xfshear` the prompt **F1 shift in ppm** can be used to shift the spectrum including its axis upfield (negative value) or low field (positive value) accordingly. For data which do not need a shearing transformation, the ppm axis in F1 can be correctly calibrated by running the AU program `xfshear` with the option **rotate**. It will calibrate the F1 axis and perform the 2D FT. The figure below compares the same 2D 3Q MAS spectrum processed with no shift and an additional shift of 5 ppm, respectively. We see that without the additional shift, the uppermost peak is at the border of the spectral range and the projection shows that the edge of this peak reenters into the spectral range from the opposite side. In summary the AU program `xfshear` can be called with the following options:

lastf1: Use the F1 shift value from last processing.

abs: Do abs2 after F2 Fourier transform of data**.**

noabs: Don't do abs2 after F2 Fourier transform of data**.**

rotate: Don't calculate shearing, only use F1 shift to rotate spectrum along F1 axis**.**

ratio: Use different value for ratio R, value can either be entered or passed.

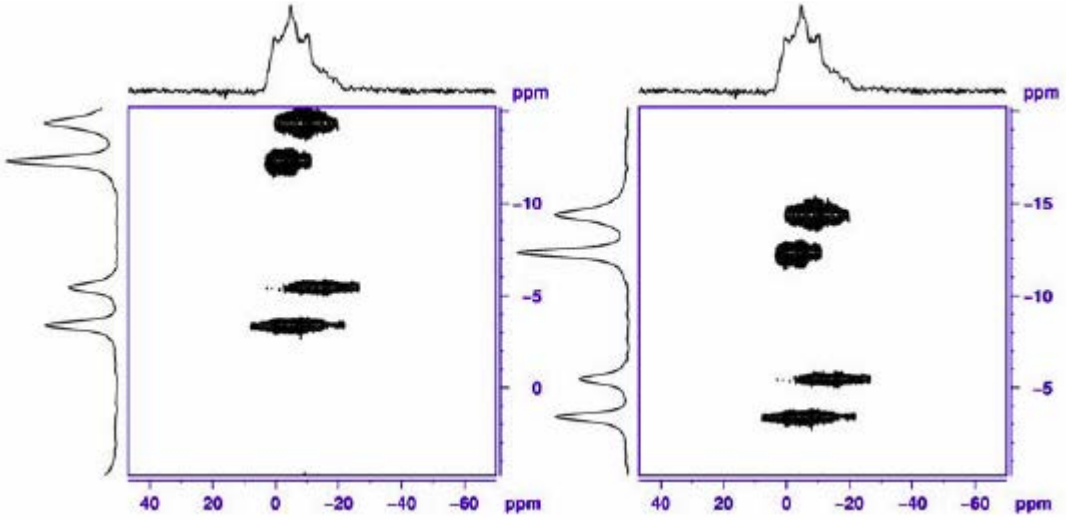


Figure 17.7: Comparison of Differently Processed 2D 23Na 3Q MAS Spectra of Na4P2O7

The left spectrum was processed with an additional F1 shift of 0 ppm, the right spectrum with +5 ppm. Spectra are taken on an AV500WB at a Larmor frequency of 132.3 MHz with a 4 mm CP-MAS probe spinning at 10 kHz. Note that the F1 range equals the spinning frequency of 10 kHz in both cases.

17.5 Obtaining Information from Spectra

The referencing procedure in **xfshear** defines the axis in the MQ dimension such that:

$$\delta_{MQ} = \delta_{iso} - \frac{10}{17}\delta_{qis}$$

The value of δqis is given by:

$$\delta_{qis} = -\frac{3(4I(I+1)-3)}{(4I(2I-1))^2} * \frac{Q_{cc}^2}{\omega_0^2}\left(1+\frac{\eta^2}{3}\right) * 10^5$$

In the above equation ***I*** is the spin quantum number, $\boldsymbol{Q}_{cc}$ the quadrupolar coupling constant, ω_0 the Larmor frequency, and h the asymmetry parameter. This makes δ_{MQ} μ ω_0^{-2}, which causes the MQ positions to be field dependent. An interesting behavior results as one compares spectra at different fields. Plots of the function δ_{MQ} over ω_0^{-2} are shown in the figure below for an arbitrary sample with two sites. If the isotropic chemical shifts of the two sites are identical, then it is obvious that the separation of the two lines increases as the field decreases (plot A). In the opposite case of identical quadrupole couplings separation increases as the field is increased (plot B). In cases where a difference in isotropic chemical shift δ_{iso} exists and the sites have different quadrupole couplings the relative positions depend on which site has the larger quadrupole coupling. The separation of the lines will always increase as the field decreases (plots C and D), but in some cases a crossover of the shift positions may be observed as the field B_0 is altered (plot C).

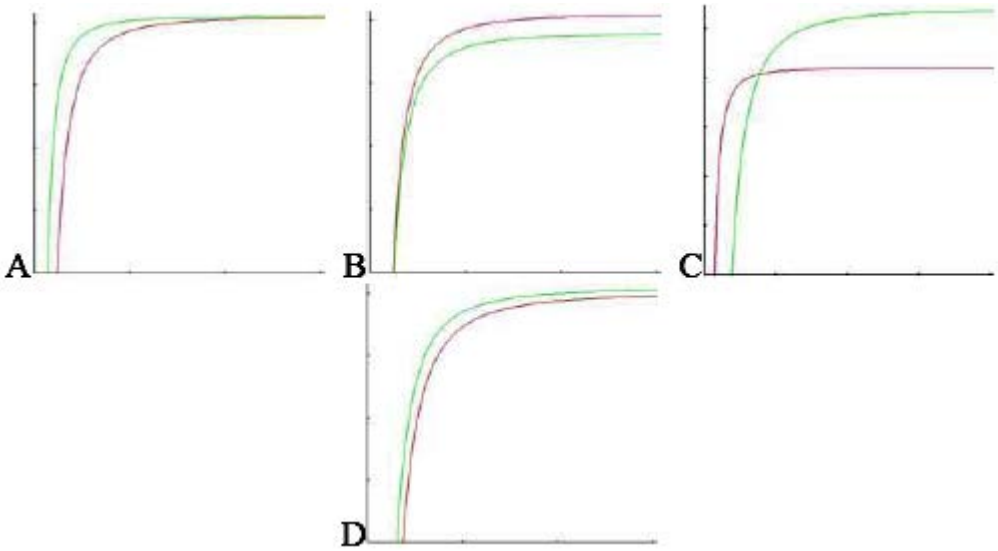


Figure 17.8: Calculated Shift Positions δMQ

Calculated shift positions δ_{MQ} as function of the static magnetic field B0 for two different sites with arbitrary δ_{iso} and δ_{qis}. The x axis in each plot is the static magnetic field B_0 increasing from left to right; the y axis δ_{MQ} increases from bottom to top. Plot A is for identical δ_{iso}, plot B for identical quadrupole coupling and. In plots C and D shift positions for two sites with large and small d_{iso} and large and small δ_{qis} and with large and small δ_{iso} and small and large δ_{qis} are plotted, respectively.

This behavior is independent of the spin quantum number and of the order p of the experiment. Higher quantum order experiments are possible for half integer spin quantum numbers >3/2, however, corresponding pulse programs are not provided in the pulse program library. They can easily be derived from the 3Q pulse program by changing the phase cycle. In the 3-pulse sequence (**mp3qzqf**) e.g. ph2 should be changed for the 5Q experiment to:

```
ph2 = 0 0 36 36 72 72 108 108 144 144 180 180 216 216 252 252 288 288
324 324
```

An 18° phase increment of the phase ph1 of the first pulse is required for States or States-TPPI phase sensitive acquisition. For a full phase cycle a multiple of 20 scans must be used.

For the 4-pulse sequence (**mp3qzfil**) the phases should be changed to:

```
ph1 = 0 36 72 108 144 180 216 252 288 324
ph2 = 0*40 90*40 180*40 270*40
ph3 = 0
ph4 = 0*10 90*10 180*10 270*10
receiver = {0 180}*5 {90 270}*5 {180 0}*5 {270 90}*5 {180 0}*5 {270
90}*5{0 180}*5 {90 270}*5.
```

Again an 18° phase increment of the first pulse for States or States-TPPI phase sensitive detection in F1 is needed. Thus a full phase cycle can be performed with a multiple of 160 scans.

The usefulness of such a 5Q experiment is limited, and there are several drawbacks: First, the sensitivity is much inferior to the 3Q experiment because of the lower transition probability and a less efficient excitation. Secondly, the shift range (in ppm) in the indirect dimension is much smaller when a rotor synchronized experiment is performed. The factors |R-p| are listed in the table below. The shift positions in the MQ dimension in a sheared spectrum are the same for all orders **p** and therefore, no additional information can be expected. However, the observed line widths are slightly reduced in the higher order experiments so in special cases some enhancement of resolution can provide additional information.

Spin I	R(p=3)	|R-p| (p=±3)	|R-p| (p=±5)	|R-p| (p=±7)	|R-p| (p=±9)
3/2	-7/9	3.78	-	-	-
5/2	19/12	1.42	7.08	-	-
7/2	101/45	0.76	3.78	10.58	-
9/2	91/36	0.47	2.36	6.61	14.17

Table 17.5: Values of |R-p| for Various Spins I and Orders p

The spectral width in the MQ dimension of the sheared spectrum is given by spinning speed / |R-p| in a rotor synchronized experiment. A 5Q experiment e.g. gives a 7.08/1.42 = 5 times smaller spectral range in the indirect dimension than a 3Q experiment.

We see that a 5Q experiment has a 5 times smaller range than the 3Q experiment and therefore, folding of peaks will always occur even at fast spinning. For even higher quantum orders the shift ranges are 7 and 30 times smaller for 7Q and 9Q than for the 3Q experiment, respectively. The table below summarizes ppm ranges for the maximum spinning frequencies of 2.5, 3.2, and 4 mm probes, respectively. A Larmor frequency of 100 MHz is assumed. One can see that the ranges become less than the typical chemical shift range for many nuclei. The expression |R-p| acts like a scaling factor that scales the frequency scale directly. Mathematically this is solved in the AU program **xfshear** in such a way that the observe Larmor frequency is multiplied by the factor ÁR-p Á to redefine an apparent Larmor frequency in the MQ dimension.

Spin I and MQ Experiment	15 kHz [4 mm probe]	25 kHz [3.2 mm probe]	35 kHz [2.5 mm probe]
3/2	39.6 ppm	66.0 ppm	92.4 ppm
5/2 3Q	105.6 ppm	176.0 ppm	246.4 ppm
5/2 5Q	26.1 ppm	35.2 ppm	49.3 ppm
7/2 3Q	197.4 ppm	329.0 ppm	460.6 ppm
7/2 5Q	39.5 ppm	65.8 ppm	92.1 ppm
7/2 7Q	14.1 ppm	23.5 ppm	32.9 ppm
9/2 3Q	319.2 ppm	532.0 ppm	744.8 ppm
9/2 5Q	63.2 ppm	160.4 ppm	144.9 ppm
9/2 7Q	45.6 ppm	76.0 ppm	106.4 ppm
9/2 9Q	10.6 ppm	17.7 ppm	24.8 ppm

Table 17.6: Chemical Shift Ranges for all MQ Experiments for All Spins I

Figures are calculated for a Larmor frequency of 100 MHz.

From the isotropic shift and the shift position in the MQ dimension the so-called SOQE parameter can be calculated, d_{qis} being given by equation 2:

$$SOQE = Q_{cc}^2\left(1+\frac{\eta^2}{3}\right) = \delta_{qis} f(I)\frac{\omega_0^2}{10^5}$$

with

$$f(I) = -\frac{(4I(2I-1))^2}{3(4I(I+1)-3)}$$

Where f(I) equals 4, 16.67, 39.2, and 72 for I=3/2, 5/2, 7/2, and 9/2, respectively. So one can see that for a given value of Q_{cc} the second order quadrupole induced upfield shift d_{qis} decreases as the spin I increases. With d_{qis} always being negative this has a direct influence on the appearance of the sheared 2D spectra.

The following figure shows 2D ^{17}O 3QMAS spectra at 11.7 T and 18.8 T where the Larmor frequency of this nucleus is 67.8 and 108.4 MHz, respectively. The sample is sodium metaphosphate $NaPO_3$ in the glassy state. The enrichment of 17O is approx. 30 to 33%. It contains 2 oxygen positions: there are bridging oxygen (P-O-P) and non-bridging oxygen (P-O-Na).

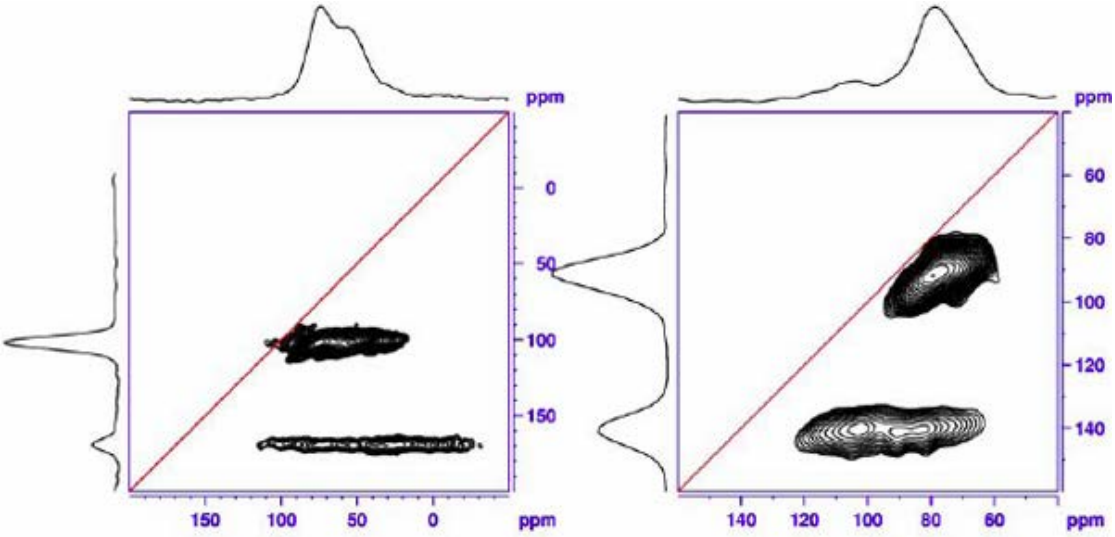


Figure 17.9: 17O MQ-MAS of NaPO3 at 11.7 T (67.8 MHz) on the left and 18.8 T (108.4 MHz) on the right

The red lines in the spectra indicate the isotropic chemical shift axis. Approximate quadrupole parameters of the two sites are $Q_{cc} \approx 7.7$ MHz, $h \approx 0.36$, $d_{iso} \approx 125$ ppm for the lower peak and $Q_{cc} \approx 4.5$ MHz, $h \approx 0.16$, $d_{iso} \approx 85$ ppm for the upper peak (sample courtesy of Alexandrine Flambard, LCPS, Univ. de Lille).

The bridging oxygen gives rise to the lower peaks in the 2D spectra in the figure above, the non-bridging ones give rise to the upper peak. An additional red line is drawn into the spectrum which represents the diagonal, meaning d(F2) = d(F1). One can see that all line positions must be below this diagonal because the negative quadrupole induced shift is scaled down and subtracted from the isotropic shift to give the MQ shift. In the example shown in two sites are visible with distinct differences in their spectroscopic parameters. In the sheared spectra we find the lower peak at 170 ppm (11.7 T) and 140 ppm (18.8 T), respectively, in the 3Q dimension. This peak is dispersed parallel to the F2 axis which means that its line width is mainly due to second order quadrupole broadening. The upper peak at 100 ppm (11.7 T) and approx. 90 ppm (18.8 T) in the 3Q dimension has a much smaller quadrupole coupling which can immediately be recognized from the fact that the peak is much closer to the diagonal. It is a very nice example where the second order broadening, which is still the dominant interaction at 11.7 T, is so much reduced at 18.8 T that the width of the peak is now determined by the distribution of chemical shift. This is expressed in the fact that the peak is dispersed along the diagonal.

The following figure shows the results of the fitting with the solids line shape analysis package included in TopSpin. The spectra used for that have been extracted from rows of the 2D spectrum shown in the previous figure.

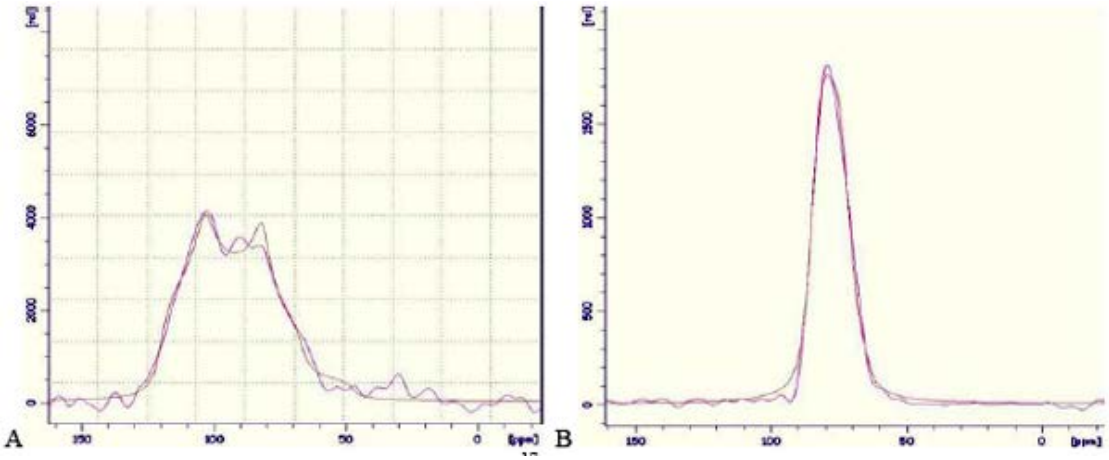


Figure 17.10: Slices and Simulations of the 18.8 T 17O MQ-MAS of NaPO3

Fitted parameters are A) Q_{cc} ≈ 7.7 MHz, h ≈ 0.36, d_{iso} ≈ 125 ppm and B) Q_{cc} ≈ 4.5 MHz, h ≈ 0.16, d_{iso} ≈ 87 ppm for the upper peak (sample courtesy of Alexandrine Flambard, LCPS, Univ. de Lille).

Spectra that are sheared can be evaluated graphically as follows, as shown in the figure below. In addition to the (red) isotropic chemical shift axis, indicated as *axis CS* with the slope $^{\Delta\delta}$ (F2)/$^{\Delta\delta}$ (F1) = 1, there are two more lines drawn. The (blue) axis indicated as *axis Qis* is the quadrupole induced shift axis with the slope $^{\Delta\delta}$ (F2)/$^{\Delta\delta}$ (F1) = -17/10. This axis is identical for all different spins **I** and all orders **p** of the MQ-MAS experiments. This axis can be shifted, retaining the same slope, so that it intersects a spectral line in its center of gravity. Through the intersection point of the Qis axis with the CS axis a third line can be drawn parallel to the F2 axis. This is the dotted black line in the figure below. The shift value that is read from the F1 axis at this position is the isotropic chemical shift of that particular site, and the Qis is then given by the first equation in this section.

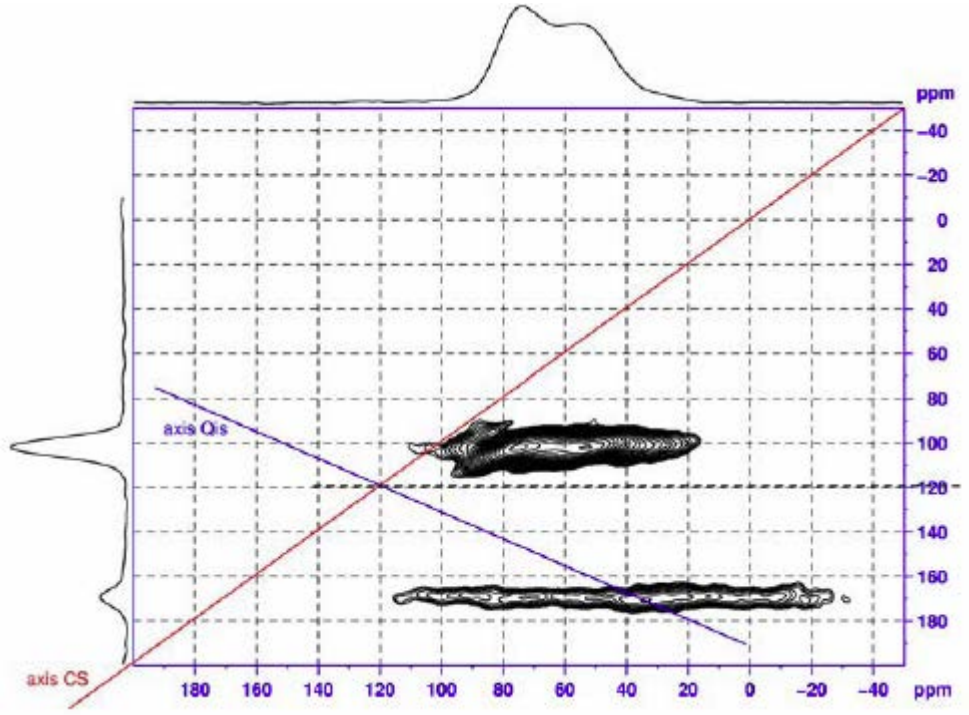


Figure 17.11: Graphical Interpretation of the Spectrum from Figure

In the 11.7 T spectrum this gives quadrupole induced shifts δ_{qis} of ≈ -75 ppm and ≈ -20 ppm for the two sites, respectively. At 18.8 T the δ_{qis} of the lower peak in the 2D spectrum decreases to ≈ -30 ppm, whereas it cannot be determined graphically anymore for the upper peak since the chemical distribution broadens the peak in the F1 dimension more than the theoretical δ_{qis} of ≈ -5 ppm.

18 MQ-MAS: Sensitivity Enhancement

The MQMAS experiment on half integer quadrupole nuclei is an extremely insensitive experiment. This is due to the low efficiencies of both the excitation of 3Q coherence and their conversion to observable magnetization. Several approaches have been taken to enhance the efficiency of the excitation and conversion, mainly focusing on the conversion, as this is the least efficient step. Adiabatic pulses can be used for the conversion instead of a single high power CW pulse, and alternative phase cycling schemes have also yielded improvements. Improving the efficiency of the MQ excitation pulse has been tried but no generally applicable scheme exists so far. Before describing the optimization procedures, some experimental approaches used in combination with these enhancement techniques are introduced.

18.1 Split-t1 Experiments and Shifted Echo Acquisition

The excitation pulse in the MQMAS experiment creates 3Q coherence that can be refocused into an observable SQ echo by the conversion pulse. As the t_1 period is incremented in successive slices of the 2D experiment this echo position relative to the conversion pulse changes as a function of the duration of the actual t_1 delay. If this observable (SQ) magnetization can be refocused again, by a central transition selective 180° pulse, a shifted echo acquisition can be implemented. The delay between the conversion pulse and the 180° pulse or the delay prior to the start of the acquisition must be incremented proportional to t_1. This results in a split-t_1 experiment where the top of the shifted echo appears at a constant position after the final pulse throughout the entire 2D experiment. The position of the echo top depends on the spin I of the observed nucleus. If the delay before the selective 180° pulse is long enough, the signal will have decayed and the full build up and decay of the echo can be recorded. By this method a phase modulated data set is acquired with a full echo that contains twice the intensity of the simple MQMAS experiment (if transverse relaxation can be neglected). At the end of the t_1 period of the split-t_1 experiment, there is no net evolution under the second order quadrupole broadening. This is the case because the evolution of the MQ coherence in the first part of t_1 is cancelled out by the evolution of the SQ coherence in the part of t_1 after the conversion pulse (the lengths of the two periods are related by the ratio of the second order broadening of the MQ and SQ coherence). The resulting 2D spectrum thus requires no shearing transformation to make f2 the isotropic dimension.

Whether this experimental trick is useful for your sample of interest can easily be determined by running a simple Hahn-echo experiment, where the delay τ in the 90°-τ-180° sequence is adjusted so that the FID of the signal is decayed before the 180° pulse is applied. This is shown in the figure below. The FID generated by the initial 90° pulse is not sampled, but after it is decayed it is refocused with a 180° pulse into a so-called shifted echo, meaning that the position of the echo can be shifted by adjusting the delay **d6**.

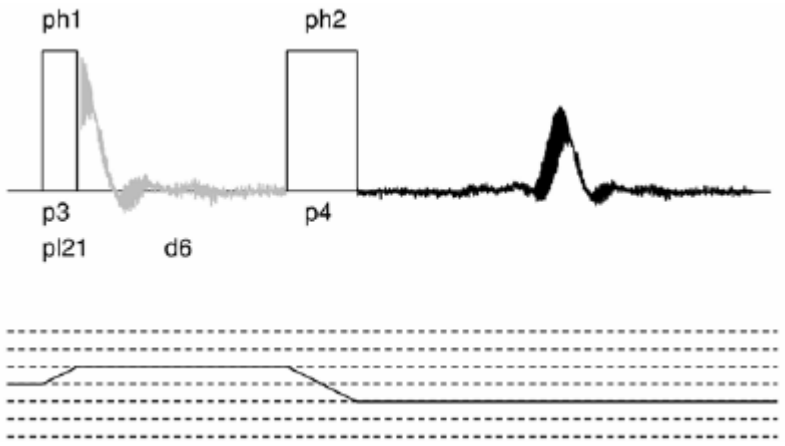


Figure 18.1: Hahn Echo Pulse Sequence and Coherence Transfer Pathway

After the initial 90° pulse **p3** the magnetization dephases and is refocused by the 180° pulse **p4**. If the τ delay **d6** is long enough a full echo can be acquired.

If data acquisition starts immediately after the second pulse the whole echo will be acquired. The integrated intensity of the echo will be almost twice the intensity of the single FID; it is just T_2 relaxation during τ that leads to attenuation. In MAS experiments it is advisable to synchronize the echo with the sample rotation i.e. make τ an integer multiple of rotor periods. For FT of the shifted echo FID there is a slight “inconvenience” as shown in because after a normal FT the signal looks quite unconventional. To obtain the usual spectrum a magnitude calculation can be done on 1D spectra, with the loss of phase information. Alternatively, and in particular in 2D spectra it is possible to apply a large 1st order phase correction **phc1** to compensate for the time delay before the echo top. The value of this is:

$$phc1 = -\frac{d6}{dw} \cdot 180°$$

This value can be entered into the processing parameters and a phase correction pk can be performed. After this the 0th order phase correction still needs to be adjusted interactively. The best method is to phase the spectrum to give minimum signal intensity and add or subtract 90° to the obtained value (click 90 or -90 in the TopSpin phasing interface).

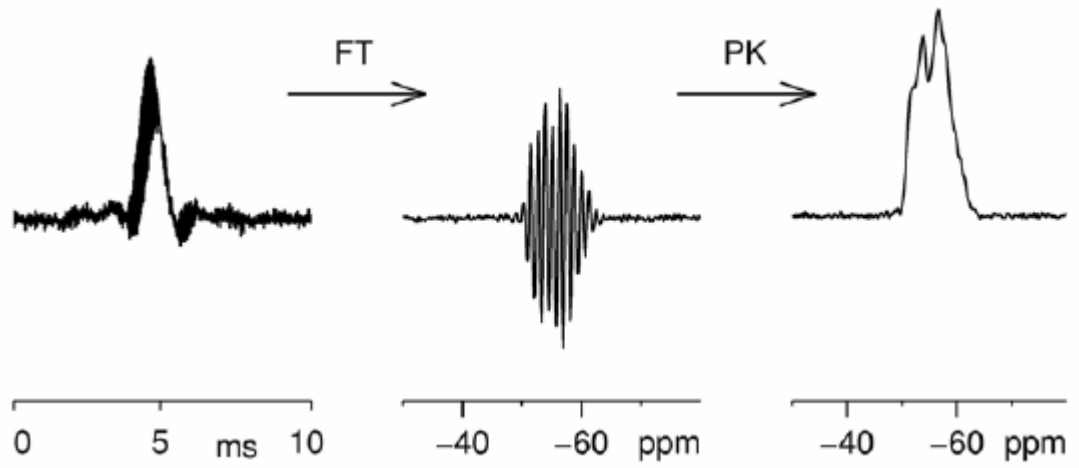


Figure 18.2: Processing of Hahn Echo. Left is the Shifted Echo

The middle shows the spectrum after FT. On the right is the spectrum with the correct first order phase correction.

18.2 Implementation of DFS into MQMAS Experiments

Two pulse sequences are available to implement a double frequency sweep (DFS). The first figure in the next section shows the 4-pulse sequence with z-filter, **mp3qdfsz.av**. The principle of this sequence is already described in the chapter *Basic MQ-MAS [▶ 191]*, with a CW pulse instead of a DFS for conversion. The second figure in the following section shows a 3-pulse sequence with a shifted echo acquisition in a split-t_1 experiment, **mp3qdfs.av**. In both sequences the same sweep is used. Both sequences start with an excitation pulse **p1** that creates 3Q coherence which is allowed to evolve during the evolution period **D0**. The sweep during **P2** is used to change (non-observable) 3Q coherency to observable SQ coherency. In the 4-pulse sequence they pass through a z-filter by a sequence of CT selective 90°-90° pulses **P3-D4-P3**. In the 3-pulse sequence a delay **D6** is introduced so that the obtained signal can first dephase and then be refocused with a 180° CT selective pulse.

18.2.1 Optimization of the Double Frequency Sweep (DFS)

At this point it is assumed that the pulses **p1**, **P3**, and **P4** together with their corresponding power levels **PL11** and **PL21** are already calibrated as described in the chapter *Basic MQ-MAS [▶ 191]*. Starting from this set-up a data set should be created into which either of the two DFS pulse programs is loaded.

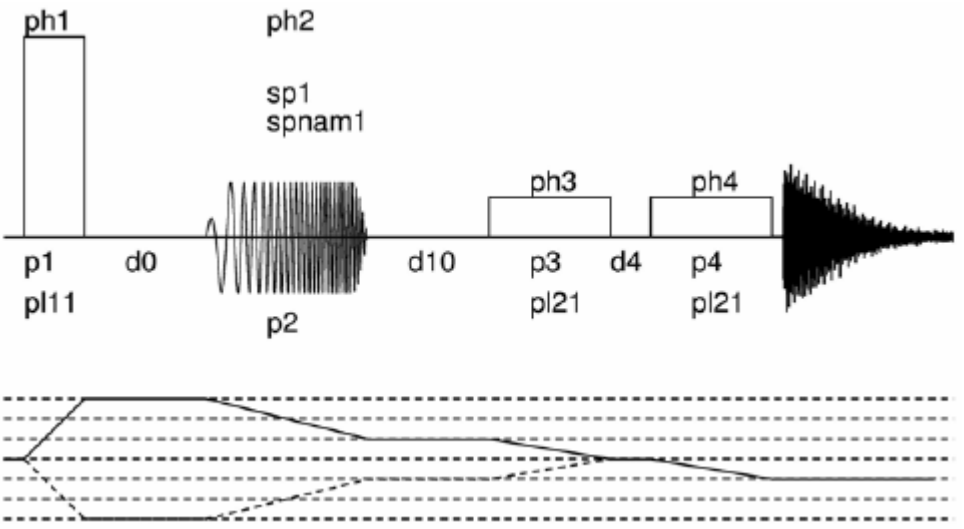


Figure 18.3: Four Pulse Sequence and Coherence Transfer Pathway for the 3Q MAS Experiment

Four pulse sequence and coherence transfer pathway for the 3Q MAS experiment with z-filter (**mp3dfsz.av**) and double frequency sweep (DFS). Excitation pulse p1 and selective pulses P3 are the same as for mp3qzfil.av. Delays **D0** and **D4** are the incremented delay for t1 evolution and 20 µs for z-filter, respectively. Delay **D10** can be incremented for spin I = 3/2 nuclei proportional to **D0**. Power level and duration of the sweep P2 must be optimized. Phase lists are as follows, for phase sensitive detection in F1 the phase of the first pulse must be incremented by 30° in States or States-TPPI mode:

```
ph1 = 0 60 120 180 240 300
ph2 = 0*24 90*24 180*24 270*24
ph3 = 0
ph4 = 0*6 90*6 180*6 270*6
receiver = {0 180}*3 {90 270}*3 {180 0}*3 {270 90}*3
{180 0}*3 {270 90}*3 {0 180}*3 {90 270}*3.
```

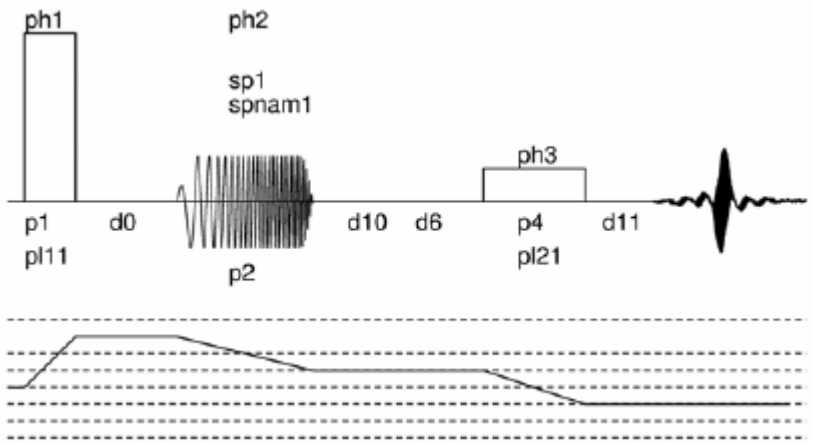


Figure 18.4: Three Pulse Sequence and Coherence Transfer Pathway

Three pulse sequence and coherence transfer pathway for the 3Q MAS experiment with z-filter (**mp3qdfs.av**). Excitation pulse **p1** is the same as for mp3qzfil.av, **P4** is a central transition selective 180° pulse (usually 2*p3). Delays **D0** is the incremented delay for t1 evolution. Delays **D10** or **D11** must be incremented proportional to **D0**. Power level and duration of the sweep **P2** must be optimized. Phase lists are as follows, 2D data are acquired in QF mode, which means that no phase incrementation is required:

```
ph1 = 0 30 60 90 120 150 180 210 240 270 300 330
ph2 = 0
ph3 = 0*12 90*12 180*12 270*12
receiver = {0 270 180 90}*3 {180 90 0 270}*3
```

Initial parameter values are listed in the first figure in section *Split-t1 Experiments and Shifted Echo Acquisition [207]*. A few of these parameters need further explanation:

D6: Is calculated as (1s* **L1**/ **CNST31**)-(**P4**/2)-(**P2**/2). This ensures that the delay from the center of the sweep to the middle of the 180° pulse is an integer multiple of the rotor period. **L1** must be set so that **D6** becomes long enough for a full echo to build up. If we assume the FID has decayed after 3 ms, spinning frequency is 25 kHz, and the 90° pulse is 20 µs, then **L1** should be between 70 and 80 so that **D6** is between 2.77 and 3.17 ms.

P2: Is calculated as 1s/(**CNST31*** **L0**). The duration of the sweep can be adjusted depending on T2 of the sample. The sweep should not be longer than one rotor period, for many samples you may find that a quarter of a rotor period or even less is sufficient.

In **ased** the values of the two parameters **D6** and **P2** are greyed, because they cannot be set anymore. They are calculated in the pulse program from the parameters explained below, which must be set accordingly:

L0: Defines the fraction of a rotor period for the duration of the sweep **P2**, usually between 1 and 8.

L1: Defines **D6** to be an integer number of rotor periods.

The sweep will be defined by further parameters:

CNST1: Start frequency (in kHz) of the sweep; the sweep should start slightly off resonance, usually 30 to 50 kHz from the resonance of the central transition.

CNST2: End frequency (in kHz) of the sweep; the sweep should cover the satellite transition, but this is often broader than the band width of the probe of approximately 1 MHz. Therefore, it does not make sense to have this value bigger than 1000.

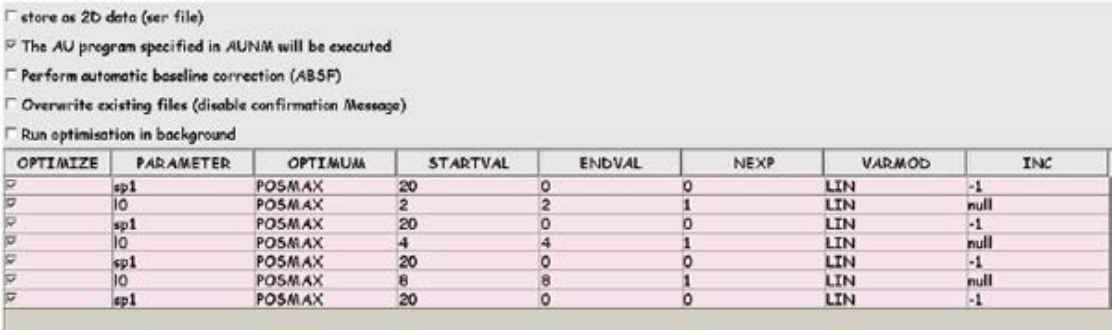
Parameter	Value	Comments
pulprog	mp3qdfs.av, or mp3qdfsz.av	Pulse program.
NS	48*n (dfs) 96*n (dfsz)	Full phase cycle is important.
D0	3 µs	Or longer, t1-period.
D1	5 * T_1	Recycle delay, use dummy scans if shorter. If d1 is too short artefacts in the 2D spectrum may show up.
D4	20 µs	Z-filter delay, mp3qdfsz.av only.
D6	see text	Calculated in pulse program mp3qdfs.av only.
D10	= 4µs	
D11	= 0	Used in mp3qdfs.av only.
P1	≤ 3.6 µs	Excitation pulse at pl11.
P2	see text	Calculated in pulse program.
P3	20 µs	90° selective pulse at pl21 used in mp3qdfsz.av only.
P4	40 µs	180° selective pulse at pl21 used in mp3qdfs.av only.
PL1	=120 dB	Not used.
PL11		Power level for excitation pulse, use value from standard MQMAS optimization.
PL21		Power level for selective pulse, approx. pl11+30 dB taken from previous pulse calibration.
spnam1	dfs	Set by AU program zg_dfs.
sp1	to be optimized	Power level for dfs.
aunm	zg_dfs	AU program to calculate sweep.
L0	see text	Fraction of rotor period for sweep.
L1	see text	Number of rotor cycles for synchronization used in mp3qdfs.av only.
cnst1	see text	Start frequency of sweep (in kHz).
cnst2	see text	End frequency of sweep (in kHz).
cnst3	50	(in ns) timing resolution for sweep.
cnst31	“masr”	Spinning frequency.

Table 18.1: Initial Parameters for the DFS Experiment

CNST3: (in ns) = 50. This is the maximum timing resolution that can be obtained on a shaped pulse, with AV or AVII hardware. However, if **cortab** is defined for that nucleus, 100 ns is the maximum timing resolution possible in a shaped pulse.

CNST31: Magic angle spinning frequency, used for the calculation of the duration of the sweep and the echo delay.

The calculation of the sweep is done via an AU program called **zg_dfs**. It calculates the sweep according to the parameters given above and stores it as a shape file, which is called **dfs**. After the calculation the AU program starts the acquisition. In order to ensure that the correct sweep is always used, it is advisable to enter the name of this AU program into the parameter **aunm**, and start all acquisitions using the command **xaua**. All that needs to be optimized now is the appropriate RF power level for the sweep, defined as parameter **sp1**. As a first guess a value of 3 dB less RF power than for the excitation pulse should be used (i.e. **SP1** = **PL11** + 3 dB). You may use **popt** for the optimization, where **SP1** is decremented by 1 dB up to the same power level used for the excitation pulse (e.g. from 20 dB to 0 dB). Initially a sweep of one whole rotor period (i.e. **L0** = 1) can be used. The optimization of **SP1** can be repeated e.g. for half a rotor period, a quarter of a rotor period and so on. With a shorter sweep you will find that higher RF power will be needed.


store as 2D data (ser file)
The AU program specified in AUNM will be executed
Perform automatic baseline correction (ABSF)
Overwrite existing files (disable confirmation Message)
Run optimisation in background

OPTIMIZE	PARAMETER	OPTIMUM	STARTVAL	ENDVAL	NEXP	VARMOD	INC
☑	sp1	POSMAX	20	0	0	LIN	-1
☑	l0	POSMAX	2	2	1	LIN	null
☑	sp1	POSMAX	20	0	0	LIN	-1
☑	l0	POSMAX	4	4	1	LIN	null
☑	sp1	POSMAX	20	0	0	LIN	-1
☑	l0	POSMAX	8	8	1	LIN	null
☑	sp1	POSMAX	20	0	0	LIN	-1

Figure 18.5: Example for popt to Set-up for Optimization of DFS

Note that the option “The AU program specified in AUNM will be executed” is checked. This ensures that the sweep is recalculated for the variation of **l0** and stored in the shape file dfs. The initial value of **l0** is 1.

The figure above shows a **popt** window with successive optimization of **SP1** for several fractions of rotor periods **l0**. Note that the check mark for “The AU program specified in AUNM will be executed” MUST be set, in order to force recalculation of the shape for each step of the optimization. In this case acquisition is run with the command **xaua**, which ensures that the correct sweep is stored in the shape file **dfs**. In the above example optimization proceeds such that in the first run **SP1** is varied from 20 to 0 in 21 steps. Then **L0** is set to 2 and **SP1** is again varied from 20 to 0. Then **l0** is set to 4 and so on. The next figure shows results of the variation of the RF power level of sweeps with different durations. Experiments have been run at 20 kHz and optimization procedures for 1, 0.5, 0.25, and 0.125 rotor periods corresponding to 50, 25, 12.5, and 6.25 µs respectively have been run. One can see that as the duration of the sweep is reduced the required RF field amplitude is higher. This is true when the spinning frequency is kept constant and the sweep is a smaller fraction of a rotor period and when the sweep is kept at e.g. one rotor period and the spinning frequency is increased (and hence the rotor period decreased). On “real life” samples the differences between signal intensities at one rotor period compared to e.g. ¼ rotor periods will be more pronounced than on a crystalline model compound. Since the spinning frequency is usually determined by the spectrum itself, the only degree of freedom is the amplitude of the sweep. In the example chosen here any of the conditions will provide a good quality spectrum, and the condition with the biggest enhancement at the least power is to be preferred.

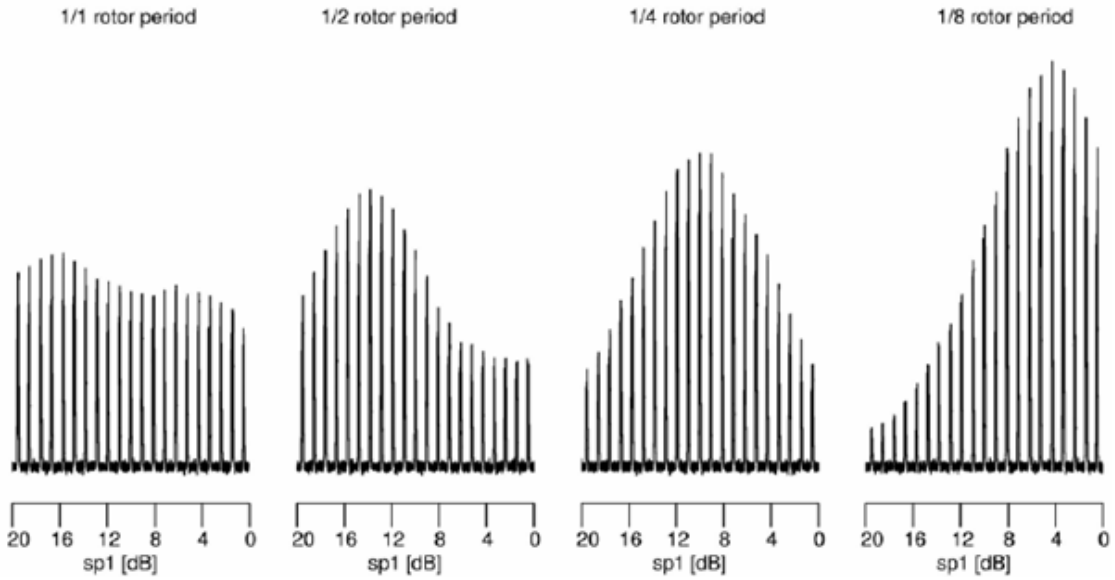


Figure 18.6: Signal Intensities of 87Rb in RbNO3

Signal intensities of ^{87}Rb in $RbNO_3$ as function of duration and RF field amplitude for double frequency sweeps.

18.2.2 2D Data Acquisition

After the parameters for the DFS are adjusted the 2D data acquisition can be prepared. In tables 2 and 3 the important parameters are listed for the two pulse sequences. Parameters are listed separately for F1 and for the pulse program relevant parameters which should be set in **eda** and **ased**, respectively.

The 3-pulse sequence used in **mp3qdfs.av** creates a phase modulated data set and therefore, FnMODE must be QF. However, since a whole echo acquisition is performed a pure absorption mode spectrum can be obtained. Increments for **D10** and **D11** must be set correctly so that a standard two-dimensional FT can be applied.

Using the 4-pulse sequence **mp3qdfsz.av** FnMode must be States or States-TPPI, so that the shearing FT can be performed for processing. However, no shearing is required in case of nuclei with spin I=3/2 where a split-t_1 experiment can be performed, in which case **IN10** must be set correctly.

For both sequences, TD in F1 determines the number of FID's to be accumulated in the indirect dimension. This value is determined by the line width and resolution that can be expected and which depend on the properties of the sample. In crystalline material fairly narrow peaks can be expected so that a maximum acquisition time in F1 of 2 to 5 ms is expected. In disordered material where the line width is broader and determined by distribution a total acquisition time in F1 up to may be 1 ms may be sufficient. The total acquisition time aq in F1 equals (**TD**/2)* **IN_010**. For rotor synchronized experiments **IN_010** = 1/spinning frequency so will typically be between 100 µs (10 kHz spinning) and 28.5 µs (35 kHz spinning), so only 100 to 250 experiments might be required. The rotor synchronization immediately means that the spectral range in F1 is limited. Dependent on chemical shift range, spinning frequency, and quadrupole interactions the positions of the peaks may fall outside this range. In such a case care must be taken when interpreting the spectrum. Acquisition with half-rotor synchronization to double the spectral window in F1 may help. However, in this situation one set of spinning sidebands appears and it must be avoided that the spinning side bands of one peak fall on top of other peaks. Some sort of rotor synchronization is always recommended because spinning side bands in the indirect dimension extend over a very wide range, which cannot be truncated by e.g. filtering.

Therefore, rotor synchronization together with States or States-TPPI phase sensitive acquisition helps to fold spinning sidebands from outside back onto center bands or other side bands.

Parameter	Value	Comments
pulprog	mp3qdfs.av	
F1 parameters:		In eda.
FnMode	QF	Acquisition mode for 2D.
TD	see text	Number of FID's to be acquired.
ND_010	1	There is only one d0 delay in the sequence.
SWH	"masr"	Equals spinning frequency for rotor synchronization, from this in0 is calculated correctly, if ND_010 must be set first.
NUC1		Select the same nucleus as for F2 so that transmitter frequency offset is set the same in both dimensions (essential for referencing).
pulse program parameters:		In ased.
D10	0	
IN10	=IN0*7/9 0	For spin I = 3/2, for all other spin I.
D11	0	
IN11	0 =IN0*19/12 =IN0*101/45 =IN0*91/36	For spin I = 3/2, For spin I = 5/2, For spin I = 7/2, For spin I = 9/2.

Table 18.2: Parameters for 2D Data Acquisition of 3-pulse Shifted Echo Experiment mp3qdfs.av.

Parameter	Value	Comments
pulprog	mp3qdfsz.av	
F1 parameters:		In eda.
FnMode	States or States-TPPI	Acquisition mode for 2D.
TD	see text	Number of FID's to be acquired.
SWH	"masr"	Equals spinning frequency for rotor synchronization, from this in0 is calculated correctly, ND_010 must be set first.
NUC1		Select the same nucleus as for F2 so that transmitter frequency offset is set the same in both dimensions (essential for referencing).
Pulse program parameters:		In ased.
D10	0	
IN10	=IN0*7/9 0	For spin I=3/2, so that no shearing FT is required. For all other spin I.

Table 18.3: Parameters for 2D Data Acquisition of 4-pulse Z-filtered Experiment mp3qdfsz.av.

18.2.3 Data Processing

Processing parameters should be set according to the table below. Data obtained with **mp3qdfs.av** can be processed with **xfb** alone, if **IN10** or **IN11** have been set appropriately to run a split-t1 experiment. Since a whole echo is accumulated FT along F2 from -Δ > t > +Δ is done which necessitates a large 1st order phase correction to compensate for the start of the acquisition before the echo top. This correction can easily be calculated as given in the equation in section *Split-t1 Experiments and Shifted Echo Acquisition [▷ 207]* and should be stored into the parameter **PHC1**. This gives an approximate value, which can be precisely adjusted in the interactive phase correction routine. The phase of the spectrum must be corrected such that there is no signal in the imaginary part, as described above.

Parameter	Value	Comments
F2 (acquisition dimension)		
SI		Usually set to one times zero filling.
WDW	no	Don't use window function.
PH_mod	pk	Apply phase correction.
BC_mod	no	No DC correction is required after full phase cycle.
ABSF1	1000 ppm	Should be outside the observed spectral width.
ABSF2	-1000 ppm	Should be outside the observed spectral width.
STSR	0	Avoid strip FT.
STSI	0	Avoid strip FT.
TDoff	0	Avoid left shifts or right shifts before FT.
F1 (indirect dimension)		
SI	256	Sufficient in most cases.
WDW	no QSINE	Don't use window function. Only use if FID in F1 is truncated.
SSB	2	π/2 shifted squared sine bell.
PH_mod	pk	Apply phase correction.
PHC1		−(d6/dw)*180.
BC_mod	no	No dc correction is required after full phase cycle.
ABSF1	1000	Should be outside the observed spectral width.
ABSF2	-1000	Should be outside the observed spectral width.
STSR	0	Avoid strip FT.
STSI	0	Avoid strip FT.
TDoff	0	Avoid left shifts or right shifts before FT.

Table 18.4: Processing Parameters

Data obtained with **mp3qdfsz.av** can be processed with the AU program **xfshear**, unless in case of nuclei with spin I = 3/2 where a straight 2D FT can be used if a split-t1 experiment has been recorded by setting **IN10** appropriately. The information obtained from the DFS enhanced spectra is the same as from the standard MQMAS experiments. Please refer to the chapter *Basic MQ-MAS [▷ 191]* for further details regarding the shearing transformation and the information obtained from MQMAS spectra.

18.3 Fast Amplitude Modulation - FAM

It must be mentioned at this point that similar approaches have been made where the frequency of irradiation is established by a fast modulation of the amplitude of the pulses. This is realized by a repetitive train of either $[\text{pulse-delay-pulse-delay}]_n$ or $[\text{delay-pulse-pulse-delay}]_n$. Pulses and delays in these trains are of the same length. The phases of the pulses are alternating +x and –x which creates a fast cosine type amplitude modulation. The frequency of this amplitude modulation appears to the spin system as an irradiated frequency.

Two pulse programs are available, **mp3qfamz.av** and **mp3qfam.av**. They correspond to the pulse sequences depicted in the figures in section *Optimization of the Double Frequency Sweep (DFS) [▶ 209]*, but the shaped pulse realizing the DFS is replaced with a sequence **D2-P2-P2-D2** embedded in a loop repeated by loop counter **L2** and with power level **PL14** for the pulses. These sequences are only useful for spin I = 3/2 nuclei. There are also sequences for higher spins, which are not included in the pulse program library. In those cases it is recommended to use DFS. The difference between FAM and DFS can be understood in such a way that FAM establishes the irradiation of a single distinct frequency whereas DFS continuously irradiates (sweeps) over a range of frequencies. These frequencies must lie in the range of the satellite transitions, therefore a single frequency irradiation is sufficient for spin I = 3/2. Higher spins have more satellite transitions and therefore, a correspondingly larger number of irradiation frequencies are required. DFS is here the most convenient solution.

Start values for the parameters determining FAM are listed in the following table:

Parameter	Value	Comments
PL14	PL11 + 3 dB	Less RF power than for a CW pulse is sufficient.
P2	0.8 µs	
D2	= P2	Calculated in the pulse program.
L2	2	

Table 18.5: Parameters for FAM

All other parameters are in full analogy to the other MQMAS pulse programs, in particular for z-filtering and creating the shifted echo. What is left is to find the best conditions for FAM optimizing **PL14**, **P2**, and **L2** consecutively. This can conveniently be performed with the parameter optimization procedure **popt**, where two or three iterations can automatically be performed.

18.4 Soft Pulse Added Mixing - SPAM

A simple and ingenious experimental trick can immediately give a signal enhancement. Starting from the standard 3-pulse sequence the phase cycling of reconversion pulse **P2** and the CT selective 90° pulse **P3** is eliminated. This changes the coherence transfer pathway from 0 → ±3 → 0 → -1 to 0 → +3 → -1, 0, +1 → -1 and 0 → -3 → -1, 0, +1 → -1. It has been shown that this leads to a substantial gain in sensitivity. However, it requires that data are acquired in echo/anti-echo mode in order to store the two different coherence transfer pathways in consecutive FID's in the serial file.

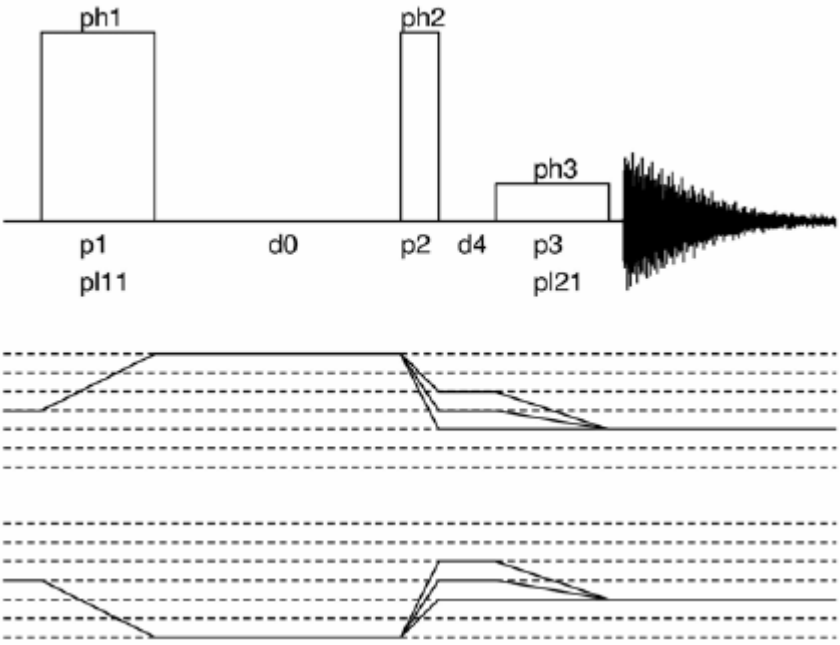


Figure 18.7: Pulse Sequence and Coherence Transfer Pathways for SPAM 3QMAS

It is extremely convenient that the setup of the pulse lengths and power levels can be done with the pulse program **mp3qzqf.av**. The setup procedure is exactly the same as described for this experiment in the chapter *Basic MQ-MAS [▷ 191]*. Before the start of the 2D data acquisition all that needs to be set is the pulse program, **mp3qspam.av**, and a small number of other parameters. These are listed in the following table:

Parameter	Value	Comments
pulprog	mp3qspam.av	
Further F1 parameters:		In eda.
FnMode	Echo/Anti-echo	Acquisition mode for 2D.
Further pulse program parameters:		In ased.
D4	0.5 µs	Not 20 µs like in mp3qzqf.av.
l4	1	Set by the pulse program, internally used counter.
l5	see text	Number of anti-echos to be acquired 0 > L5 > td{F1}/2.
l6	3 1	For spin I = 3/2. For all other spins.

Table 18.6: Further Parameters for 2D Data Acquisition of SPAM MQMAS Experiment mp3qspam.av

FNMODE: Even though this parameter is not evaluated by the pulse program it will be used by the processing AU program **xfshear**.

D4: A very short delay is used here, just to allow for amplitude and phase switching.

I4: This loop counter is internally used for checking if the echo or anti-echo is currently being acquired.

I5: In the acquisition of echo-anti-echo 2D spectra signals from the echo and anti-echo pathways are stored into consecutive FID's in the serial file. In MQMAS experiments these echos and anti-echos behave differently. For t_1 = 0 both signals have their echo top immediately after the selective 90° pulse. As t_1 is incremented the top of the echo appears at later point in time whereas the top of the anti-echo appears at an earlier point in time. It means that the contribution of the anti-echo becomes less and less until finally the signal fades out completely and only noise is sampled. It can be advantageous to terminate the acquisition of this noise in order to increase the overall S/N and save spectrometer time. However, in the processing of echo-anti-echo data two consecutive FID's are linearly combined in the following way:

$$re1 = -im2 - im1 \quad re2 = re2 - re1$$

$$im1 = re2 + re1 \quad im2 = im2 - im1$$

Where *re* and *im* refer to the real and imaginary points of FID's 1 and 2. Hence, acquiring a smaller number of anti-echos than echos leads to the usual truncation effects (wiggles in the spectrum). Furthermore, since both signals contribute to the phase information care must be taken that the pure absorption line shape of the 2D peaks is not obscured. Therefore, in case of doubt it is probably the best idea to set **L5** = **TD**{F1}/2. If less anti-echos are to be accumulated the question is how many anti-echos to acquire - this depends on the sample. In amorphous or disordered materials the FID decays rapidly and so does the anti-echo. In such a case 4 to 8 anti-echos may be sufficient. In the case of crystalline materials it takes many more t1 increments before the anti-echos decay. Hence, the number of anti-echos should be of the order of half the number of echos. It is always better to acquire more anti-echos than are really needed, because then you can be sure that you acquire a 2D spectrum with a reliable 2D absorption line shape. Never risk gaining sensitivity or saving experimental time at the expense of quality of line shapes.

II6: The value of this loop counter is needed to set the phases of the soft pulses correctly and define what is an echo and what is an anti-echo (which are different for spin I = 3/2 and all the other spin quantum numbers).

Processing of these spectra is done in analogy to spectra obtained with **mp3qzqf.av**. However, phase correction in the acquisition dimension F2 cannot be determined on the first FID. Therefore, **xf2** must be applied first and then F2 phase correction can be determined on either the first slice, in case of nuclei with spin I = 3/2, or the second slice for all other nuclei. 2D processing is then done with the AU program **xfshear**. Alternatively **xfshear** can be used first, with a subsequent 2D interactive phase correction.

19 STMAS

The STMAS experiment for half integer quadrupole nuclei is a 2D experiment to separate anisotropic interactions from isotropic interactions. In the NMR of half integer quadrupole nuclei the dominant anisotropic broadening of the central +1/2 ↔ –1/2 transition (CT), and symmetric multiple-quantum (MQ) transitions, is the 2nd order quadrupole interaction which can only partially be averaged by MAS. The satellite transitions (ST, e.g. the ±3/2 ↔ ±1/2 transitions) however, are broadened by a 1st order interaction, which is several orders of magnitude larger than the 2nd order broadening. Under MAS the 1st order interaction of the ST can be averaged but spinning cannot be fast compared to the first order broadening (of the order of MHz), a large manifold of spinning side bands remains. The 2nd order broadening of the CT can only be narrowed by a factor of 3 to 4 so a signal is observed that still reflects this 2nd order broadening.

The 2D STMAS experiment exploits the fact that the 2nd order broadening of the ST transitions (e.g. + 3/2 ↔ +1/2 in a spin 3/2), is related to the 2nd order broadening of the CT by a simple ratio. A 2D spectrum is recorded which correlates a single quantum coherence of the satellite transitions (usually one of the inner transitions, ± 3/2 ↔ ± 1/2), and the + 1/2 ↔ – 1/2 single quantum coherence of the central transition. The resulting 2D spectrum yields an isotropic projection where the 2nd order broadening has disappeared. The information content is in full analogy to the MQMAS experiment.

19.1 Experimental Particularities and Prerequisites

In contrast to the MQMAS experiment the first pulse in the STMAS experiment excites single quantum (SQ) coherency. The signal which is thus generated consists of contributions from both the CT and the ST. In the figure below the contribution of the CT shows up in the cosine curve starting at the (blue) filled rectangle resembling the initial pulse.

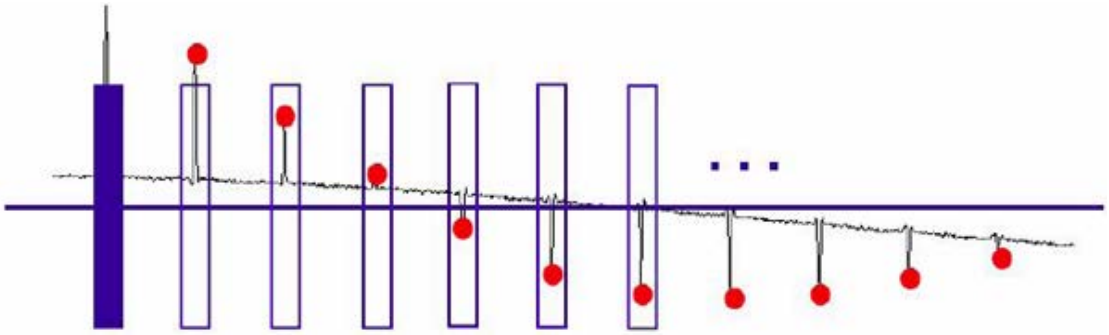


Figure 19.1: Principle of 2D Data Sampling in STMAS Experiments.

The (blue) filled rectangle on the left symbolizes the first pulse, which starts the evolution period t_1. After each revolution of the rotor, rotational echo's show up which are indicated by the (red) filled circles. The open rectangles symbolize the second pulse (one pulse at the end of each individual t_1 increment). They must always occur precisely on top of the rotational echo. For each increment the t_2 data acquisition, which is not shown here, starts after the second pulse.

In the following discussion we will ignore this part completely. It showed up in the original experiments but can be completely suppressed by a double quantum filter. The contribution of the ST "rides" on top of the CT signal like spikelets. Since MAS efficiently averages the 1st order quadrupole interaction of the ST, the corresponding MHz broad signal is now narrowed into a huge number of spinning side bands. These coherency originating from the ST dephase rapidly and refocus into rotational echoes with each rotor cycle. A pulse precisely on

top of such a rotational echo can transfer the SQ coherency from the ST to SQ coherency of the CT, the signal from which can then be acquired under standard MAS conditions. The evolution in the indirect dimension is achieved in such a way that the delay between the two pulses, which is the evolution period t_1, is incremented by integer multiples of the rotor period.

Two extremely important points must be considered for the experimental realization of the STMAS experiment. Firstly, the spinning frequency must be kept absolutely constant. The duration of the rotational echoes in the STMAS experiment is determined by the width of the satellite transition, giving a length of e.g. 1 µs for a satellite transition of 1 MHz width. If the rotor period varies from that specified in the parameters, the calculated delay in the pulse program is incorrect and the pulse misses the echo top, so less or no signal intensity is obtained. The table below summarizes the time deviation that occurs when the spinning frequency fluctuates by ± 1 Hz and ± 10 Hz at various desired spinning frequencies. One can see that when the t1 increment accumulates to as much as 100 rotor periods it is possible to miss an echo completely. For example, if the duration of the rotational echo is 1 µs it will be missed when the deviation is larger, which is the case for a 1 Hz deviation at 10 kHz, but requires a fluctuation of 10 Hz at 30 kHz spinning.

Fluctuation of … Hz at desired spinning frequency	Deviation from precise rotor period after 1 rotor period	Deviation from precise rotor period after 100 rotor periods
10 Hz @ 30 kHz	11 ns	1.1 µs
1 Hz @ 30 kHz	1.1 ns	110 ns
10 Hz @ 20 kHz	25 ns	2.5 µs
1 Hz @ 20 kHz	2.5 ns	250 ns
10 Hz @ 10 kHz	100 ns	10 µs
1 Hz @ 10 kHz	10 ns	1 µs

Table 19.1: Time deviation of the rotor period for spinning frequency variations of ± 1 and ± 10 Hz for various spinning frequencies

Typically the spinning frequency must be stable within ≤1 Hz throughout the entire 2D data acquisition. Secondly, the accuracy of the magic angle setting is extremely important. The sidebands resulting from the first order broadening are narrowed from the full first order line width by a factor of $(3\cos^2\theta-1)$, hence for a deviation of $d\theta$ from the magic angle the broadening is $3\cos\theta\sin\theta d\theta$, which close to the magic angle is $\sqrt{2}d\theta$. The magnitude of the interaction that must be narrowed in the present case is in the order of MHz so even a small deviation causes a severe broadening in the STMAS spectrum. This can easily be understood as the rotational echoes decay much more rapidly when the magic angle is off. Experimentally it has been found that the precision for setting the angle must be ≤ 0.002°. The dependence on the accuracy is so important that the experiment itself must be used to find the most precise magic angle setting. It is obvious that this can only be done on a well-known sample, as we will see later in the chapter. In order to achieve the necessary precision for the adjustability of the angle a special goniometer screw with an adapted gear transmission ratio is provided for the magic angle setting knob as an upgrade for Bruker WB MAS probes. Once the best angle setting is found it is advisable to leave the probe in the magnet. Sample changes, however, on the WB probes doesn't change the setting noticeably.

19.2 Pulse Sequences

The figures in this section show two of the basic sequences, which are 4-pulse sequences with z-filter (**stmasdqfz.av** and **stmasdqfe.av**). Both sequences start with a non-selective excitation pulse p1 that creates SQ coherency on the innermost ST which is allowed to evolve during the evolution period **D0**. Shortly before the end of the t1 period there is a selective 180° pulse **P4**. This provides a double quantum filter by which magnetization of the CT transition is eliminated which will otherwise give a strong diagonal signal from a CT –> CT coherence transfer pathway. The t_1 period is terminated with the second non-selective pulse **P2**.

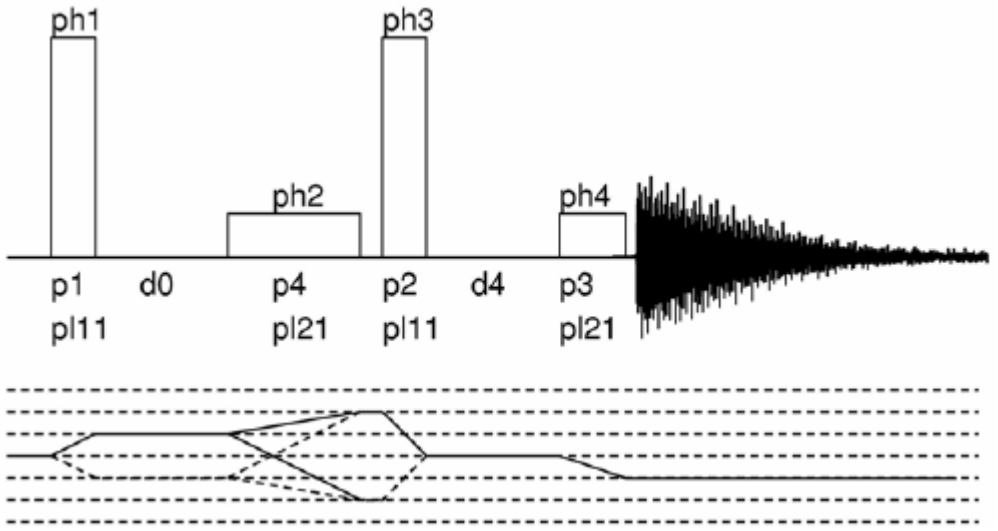


Figure 19.2: Four-pulse sequence and coherence transfer pathway for the double quantum filtered STMAS experiment with z-filter (stmasdqfz.av)

Pulses **p1** and **P2** are non-selective pulses. The corresponding power level **PL11** should be set to achieve around 100 kHz RF field amplitude. **P3** and **P4** are CT selective pulse 90° and 180° pulses of about 20 and 40 µs, respectively, corresponding to an RF field amplitude of a few kHz. Delays **D0** and **D4** are the incremented delay for t1 evolution and 20 µs for z-filter, respectively. Phase lists are as follows, for phase sensitive detection in F1 the phase of the first pulse must be incremented by 90° in States or States-TPPI mode:

```
ph1 = 0 2
ph2 = 0 0 2 2
ph3 = 0 0 0 0 1 1 1 1 2 2 2 2 3 3 3 3
ph4 = 0*8 1*8 2*8 3*8 4*8
receiver = 0 2 2 0 2 0 0 2 0 2 2 0 2 0 0 2 1 3 3 1 3 1 1 3 1 3 3 1 3
1 1 3
2 0 0 2 0 2 2 0 2 0 0 2 0 2 2 0 3 1 1 3 1 3 3 1 3 1 1 3 1 3 3 1.
```

In the **stmasdqfz.av** pulse sequence this pulse flips the magnetization back along the z-axis. After a short z-filter delay **D4** a CT selective 90° pulse **P3** creates transverse magnetization. In the **stmasdqfe.av** pulse sequence the non-selective pulse **P2** converts the ST SQ coherence into CT SQ coherence. This is allowed to evolve for another delay **D6** after which it refocuses into an echo by a CT selective 180° pulse. When either the **D6** before the 180° pulse or **D7** after the 180° pulse is incremented proportionally to the t_1 period a split-t_1 experiment as described in chapters 16 and 17 for some MQMAS experiments will be performed.

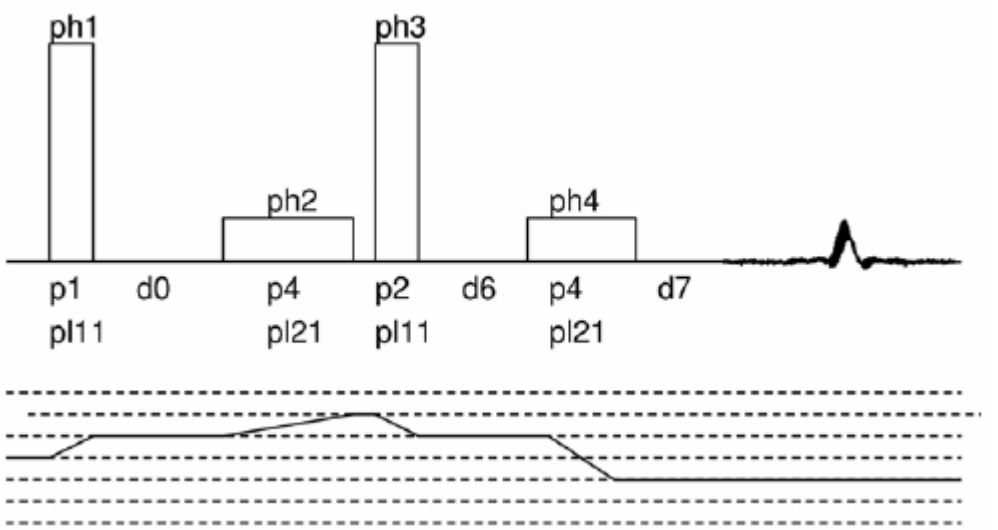


Figure 19.3: Four pulse sequence and coherence transfer pathway

Four pulse sequence and coherence transfer pathway for the double quantum filtered STMAS experiment with shifted echo acquisition (**stmasdqfe.av**). Pulses **p1** and **P2** are non-selective pulses. Corresponding power level **PL11** should be set to achieve around 100 kHz RF field amplitude. **P3** and **P4** are CT selective pulse 90° and 180° pulses of about 20 and 40 µs, respectively, corresponding to an RF field amplitude of a few kHz. Delay **D0** is the incremented delay for t1 evolution, **D6** and **D7** can be incremented proportional to **D0** depending on the spin of the observed nucleus. Phase lists are as follows, incrementation of the phase of the first pulse is not required because a phase modulated data set is acquired with FnMODE being QF:

```
ph1 = 0 180 90 270
ph2 = 0*4 90*4 180*4 270*4
ph3 = 0*16 90*16 180*16 270*16
ph4 = 0
receiver = ph3-ph1-ph2
```

19.3 Experiment Setup

Before the 2D experiment on your sample of interest can be started some setup steps must be done as described in detail below. All setup steps should be done on a sample with:

- A known MAS spectrum,
- With sufficiently good sensitivity to facilitate the set-up, and,
- A 2nd order quadrupole interaction of the order of the one expected for your sample of interest.

In a first step, a low power selective pulse must be calibrated in a single pulse experiment. After this the STMAS experiment can be optimized using the 2D pulse sequence for the first t_1 increment.

19.3.1 Setting Up the Experiment

Sample: There are a large number of crystalline compounds that can be used to set-up the experiment. Please refer to the following table to select a suitable sample. For the general procedure described here the spin I of the nucleus is not important, of course the obtained pulse widths will depend on the spin I, and the Larmor frequency. For the STMAS experiment, in contrast to MQMAS, it is advisable to use a well-known sample for the setup because the accuracy of the magic angle setting is extremely critical.

Nucleus	Spin	Spectrometer frequency*1)	d1 [s]*3)	Sample	Comments
^{17}O	5/2	67.78	2	$NaPO_3$	> 10% enriched
^{11}B	3/2	160.42	>5	H_3BO_3	
^{23}Na	3/2	132.29	10	Na_2HPO_4*2)	
^{27}Al	5/2	130.32	5	YAG	
^{87}Rb	3/2	163.61	0.5	$RbNO_3$	
^{93}Nb	9/2	122.25	1	$LiNbO_3$	
*1) In MHz at 11.7 T (i.e. 500.13 MHz proton frequency). *2) Alternatively Na_2HPO_4 * $2H_2O$ can be used. For anhydrous Na2HPO4 the sample should be dried at 70° C for a couple of hours before packing the rotor in order to eliminate crystal water completely. *3) Recycle delays at 11.7 T, longer delays may be required at higher fields.					

Table 19.2: Some Useful Samples for Some Nuclei with Half Integer Spin

As for MQMAS the setup must be done in two steps; in the first step a central transition selective pulse that merely excites the central transition must be calibrated. This pulse must be weak enough so that only the central transition is affected and it must be short enough so that the central transitions of all sites in the spectral range are excited. These conditions are typically fulfilled by a 20 µs pulse. For the calibration of this pulse a power level around 30 dB with 500 W and 1 kW amplifiers and around 20 dB with 300 W amplifiers should be expected. The pulse program **zg** (which uses **p1** and **pl1**) or **zgsel.av** (which uses **P3** and **PL21**) can be used. For more details please refer to *Relaxation Measurements [▶ 181]*.

Once the central transition selective 90° pulse is calibrated the STMAS pulse program can be loaded. Available pulse programs are **stmasdqfz.av** and **stmasdqfe.av**. Both are double quantum filtered 4-pulse sequences, the first with a z-filter, the second with a shifted echo. If this second sequence is to be used a proper setting of the timing for the shifted echo is required, to allow collection of the full echo signal. This is explained in *Basic MQ-MAS [▶ 191]*, where a shifted echo can be used in DFS enhanced MQMAS experiments.

In the next two tables the starting parameters for the setup of the two sequences are given. Typical values for the pulses are entered so one should see some signal for further optimization. Parameters like **O1**, **TD**, **SWH**, **RG**, should already be set in the standard 1D spectrum. Since the experiment is not as dependent on the pulse lengths or the applied RF field amplitude as MQMAS, pulse lengths between 1 and 2 µs, which can be achieved with every probe with 4 mm or smaller rotor diameter, are sufficient.

Parameter	Value	Comments
pulprog	stmasdqfz.av	Pulse program.
NS	16*n	For set-up the full phase cycle is not so critical.
D0	see text	Calculated in pulse program.
D1	5 * T1	Recycle delay, use dummy scans if shorter.
D4	20 µs	Z-filter delay.
P1	1.5 µs	Excitation pulse at pl11.
P2	1.5 µs	Conversion pulse at pl11.
P3	20 µs	90° selective pulse at pl21 taken from previous pulse calibration.
PL1	=120 dB	Not used
PL11	start with ≈ 150 to 300 W	Power level for excitation and conversion pulses.
PL21		Power level for selective pulse, approx. pl11+30 dB taken from previous pulse calibration.

Table 19.3: Initial Parameters for the Set-up of stmasdfqz.av

For **PL11** an initial value that corresponds to 150 to 300 W can be used. Optimization will be done on the first increment of the 2D sequence, which is calculated within the pulse program according to „D0=(1s* **L0**/ **CNST31**)- **P1**/2- **P4**-0.3µ- **P2**/2“, because it is essential that the centers of the pulses **P1** and **P2** are exactly an integer number of rotor periods apart. In this formula **P1**, **P2**, and **P4** are the RF pulses as listed in table 2, **CNST31** must be set equal to the spinning frequency. This means that the first increment can last between 100 µs (10 kHz spinning) and 28.5 µs (35 kHz spinning). Since **P4** is the 180° selective pulse, which can be as long as 40 µs, **L0** must be set large enough to avoid the situation where the calculated d0 is negative. Optimization of the pulses **P1** and **P2** can be done using **popt** in full analogy to the optimization of the pulses in MQMAS.

Parameter	Value	Comments
pulprog	stmasdqfe.av	Pulse program.
NS	16*n	For set-up the full phase cycle is not so critical.
D0	See text	Calculated in pulse program.
D1	5 * T1	Recycle delay, use dummy scans if shorter.
D6	See text	
D7	See text	
P1	1.5 µs	Excitation pulse at pl11.
P2	1.5 µs	Conversion pulse at pl11.
P3	20 µs	90° selective pulse at pl21 taken from previous pulse calibration.
PL1	=120 dB	Not used.
PL11	Start with ≈ 150 to 300 W	Power level for excitation and conversion pulses.
PL21		Power level for selective pulse, approx. pl11+30 dB taken from previous pulse calibration.

Table 19.4: Initial Parameters for the Set-up of stmasdfqe.av

19.3.2 Two Dimensional Data Acquisition

Once the pulses are calibrated the 2D data acquisition can be used to find the correct and precise magic angle setting. Create a new data set and change parmode to 2D. The acquisition parameters for the (new) indirect F1 dimension must be set according to the following table.

Parameter	Value	Comments
F1 parameters:		In eda.
FnMode	States TPPI, or States QF	2D acquisition mode for stmasdqfz.av. 2D acquisition mode for stmasdqfe.av.
TD	see text	Number of FID's to be acquired.
SWH	"masr"	Equals spinning frequency for rotor synchronization, from this IN0 is calculated correctly, if ND_010 is already set.
NUC1		Select the same nucleus as for F2 so that transmitter frequency offset is correctly set (important for referencing).
Pulse program parameters:		In ased.
D6	0	Used in stmasdqfe.av only.
IN6	=IN0*8/9	Used in stmasdqfe.av for I = 3/2.
D7	0	Used in stmasdqfe.av only.
IN7	=IN0*7/24 =IN0*28/45 =IN0*72/55	Used in stmasdqfe.av for I = 5/2. Used for I = 7/2. Used for I = 9/2.

Table 19.5: F1 Parameters for the 2D Data Acquisition

Similar considerations for the maximum t_1 period, determined by the number of FID's to be acquired and the t_1 increment, can be made as for MQMAS. Because the shift range in ppm is twice as big as in 3QMAS a larger increment can be used to give an equivalent shift range, the increment being calculated from the spinning speed. Since the magic angle is probably not yet perfect 32 to 64 FID's will be sufficient initially. Processing parameters are described in the next section.

In the figure below two 2D plots of the ^{87}Rb STMAS experiment on $RbNO_3$ are compared. The spectrum on the left was obtained after the first execution of the experiment. The spectrum on the right was obtained after several iterations of resetting the angle and rerunning the spectrum. From this it is obvious that the spectrum of the sample for the setup must be known because otherwise it is impossible to judge whether a shoulder or a splitting is due to an incorrectly set angle or another signal from another site in the sample.

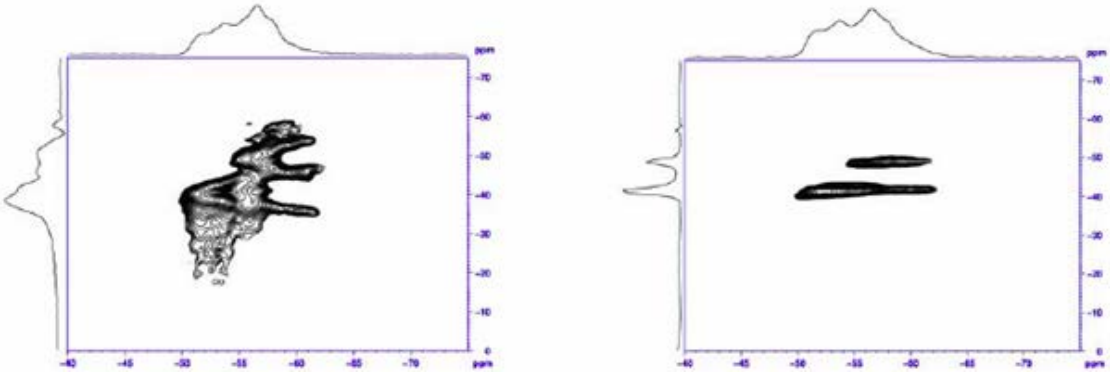


Figure 19.4: 87Rb STMAS Spectra of RbNO3

While the left spectrum has been obtained after adjusting the magic angle with KBr, the right spectrum can be obtained after several iterations of readjusting the angle and rerunning the 2D spectrum.

19.4 Data Processing

Processing parameters should be set according to the table below:

Parameter	Value	Comments
F2 (acquisition dimension)		
SI		Usually set to one times zero filling.
WDW	No	Don’t use window function.
PH_mod	Pk	Apply phase correction.
BC_mod	No	No DC correction is required after full phase cycle.
ABSF1	1000 ppm	Should be outside the observed spectral width.
ABSF2	-1000 ppm	Should be outside the observed spectral width.
STSR	0	Avoid strip FT.
STSI	0	Avoid strip FT.
TDoff	0	Avoid left shifts or right shifts before FT.
F1 (indirect dimension)		
SI	256	Sufficient in most cases.
WDW	no	Don’t use window function, unless F1 FID is truncated.
PH_mod	pk	Apply phase correction.
BC_mod	no	No DC correction is required after full phase cycle.
ABSF1	1000	Should be outside the observed spectral width.
ABSF2	-1000	Should be outside the observed spectral width.
STSR	0	Avoid strip FT.
STSI	0	Avoid strip FT.
TDoff	0	Avoid left shifts or right shifts before FT.

Table 19.6: Processing Parameters for the 2DFT

Data obtained with **stmasdqfe.av** can be processed with **xfb** only, if **IN6** or **IN7** have been set appropriately to run a split-t1 experiment. Data acquired with the pulse program **stmasdqfz.av** should be processed with the AU program **xfshear**. For information about this program please refer to *Relaxation Measurements [181]*. In an analogous way to MQMAS spectra the apparent Larmor frequency in the indirect dimension is recalculated by multiplying the real Larmor frequency with the corresponding value of |R-p|. The values for the different spin quantum numbers are summarized in the following table, for experiments using the inner ST (± 3/2 ↔ ± 1/2). R determines the shearing ratio, i.e. the slope in a non-sheared spectrum, |R-p| is the scaling factor for referencing in the indirect dimension. Using this procedure the shift positions in the indirect dimension are identical (in ppm) to all MQMAS experiments, and the information obtained is the same.

Refer to the chapter *Basic MQ-MAS [191]* for details about the information obtained from such spectra.

Spin I	R	|R-p| (p = ± 1)
3/2	-8/9	1.889
5/2	7/24	0.70833
7/2	28/45	0.3777
9/2	55/72	0.263111

Table 19.7: Values of R and |R-p| for the Various Spin Quantum Numbers Obtained in the STMAS Experiment

20 Double-CP

Double Cross Polarization (DCP) experiments use two consecutive cross polarization steps. Usually, the first step transfers from protons to one type of X-nucleus (to achieve high sensitivity), the second step transfers to a different (Y) nucleus in order to probe the dipolar coupling between X and Y. The sequence of transfers is in principle arbitrary, but usually sensitivity is an issue, so transfer from protons (to generate a large magnetization) and detection on the nucleus of higher sensitivity (to gain signal intensity) is the standard procedure. Detection of the most sensitive nucleus, protons, is also possible, but is difficult if the homonuclear proton-proton dipolar coupling is strong (see *CRAMPS: General [▶ 243]*).

In this chapter, the most popular double CP experiment is described. Here, the first CP step transfers magnetization from protons to ^{15}N. Then, in a second cross polarization step, magnetization is transferred from ^{15}N to ^{13}C; the signal is finally detected on ^{13}C under suitable proton decoupling. The purpose of this experiment is to gain information about the C-N dipolar coupling which in turn provides special information.

Naturally, the C-N, or in general, the X-Y dipolar coupling is much smaller than any dipolar coupling involving protons. For C-N, it is <2.5 kHz. This has some experimental consequences:

1. There is no need to decouple ^{15}N while observing ^{13}C, since the coupling is spun out already at moderate spin rates.
2. The Hartmann-Hahn condition for this cross polarization is extremely sharp and must be adjusted very carefully for every spin rate.
3. The magnetization transfer is substantially slower than from protons, meaning that contact times are usually longer.
4. The transfer occurs (unlike CP from protons) not out of a bath of abundant spins, but behaves (especially at high spin rates) more like a transfer between spin pairs.
5. Labeled samples must be used so that an observable number of coupled spins is present.

Advanced experimental schemes use tangential pulses to provide adiabatic conditions during the cross polarization (S. Hediger et al.) or provide only selective polarization transfer, Specific CP (Baldus et al.).

References

J. Schaefer, T.A. Skokut, E.O. Stejskal R.A. McKay, and J.E. Varner, *Proc. Nat*. Acad. Sci. USA 78, 5978 (1981).

J. Schaefer, E.O. Stejskal, J.R. Garbow, and R.A. McKay, *Quantitative Determination of the Concentrations of 13C-15N Chemical Bonds by Double Cross-Polarization NMR*, J. Magn. Reson.59,150-156 (1984).

M. Baldus, A.T. Petkova, J. Herzfeld, and R.G. Griffin, *Cross Polarization in the tilted frame: assignment and spectral simplification in heteronuclear spin systems*. Mol. Physics 5, 1197-1207 (1998).

20.1 Pulse Sequence Diagram, Double CP (DCP)

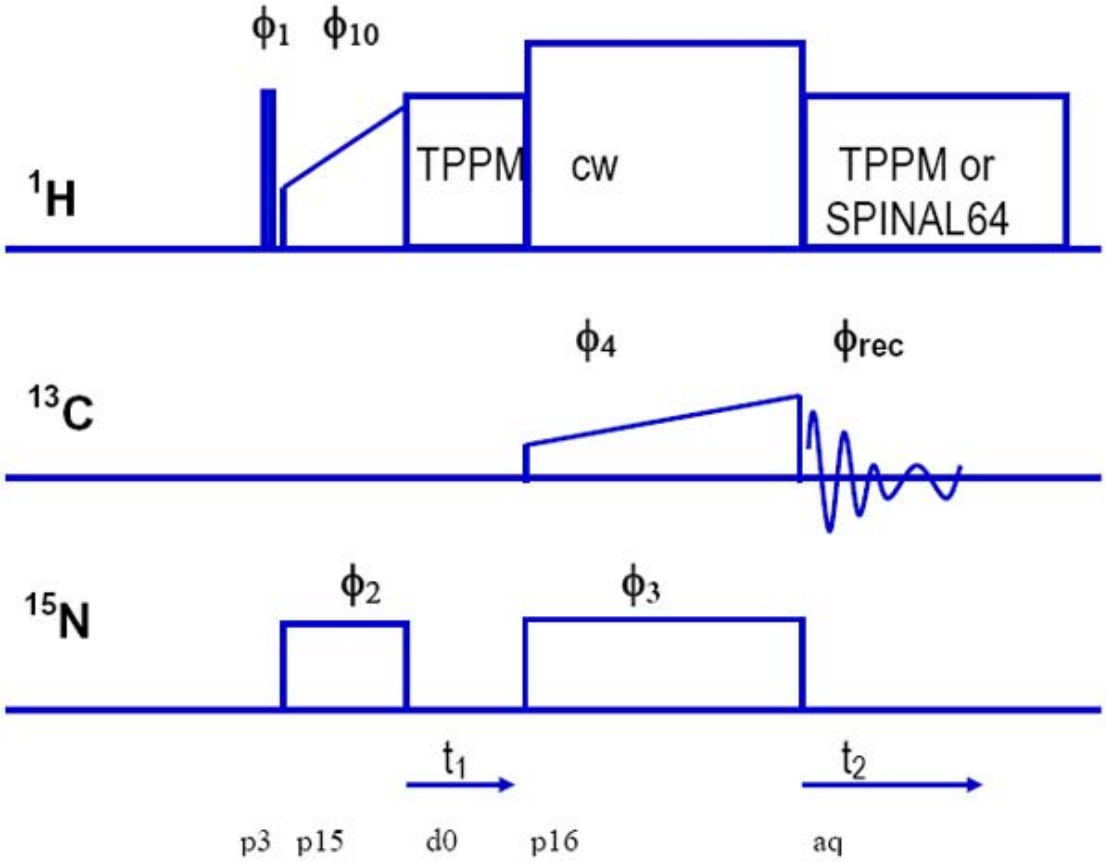


Figure 20.1: Pulse sequence diagram for 1D (t1=0) and 2D double CP experiments.

φ1= 1, 3	φ2= 0 + States-TPPI (t1)	φ4= 0 0 0 0 1 1 1 1 2 2 2 2 3 3 3 3	φrec=0 2 2 0 3 1 1 3 2 0 0 2 1 3 3 1
	φ3= 0 0 0 0 2 2 2 2		

20.2 Double CP Experiment Setup

20.2.1 Double CP 2D Experiment Setup

- Prepare your probe for triple resonance applications H/C/N
- Load a sample of glycine, ^{15}N and C_1, -C_2 (or only C_2) ^{13}C labelled. Make sure the sample is α-glycine, you will get nowhere with γ-glycine, since the proton $T_{1\rho}$ is very short and CP just does not work with high efficiency. Rotate at 11 kHz. The sample may be fully labelled or diluted with natural abundance glycine. A restricted volume rotor is preferred. If a different spin rate is used, a different shape must be generated for the second CP step.
- Check the edasp routing and set up 3 RF channels for C, H and N, such that the lower power amplifier (500 W or less) is used for ^{13}C. (^{15}N may require more than 500 W). Set for ^{13}C observation.
- Make sure the preamplifiers in use are set up for the appropriate frequencies. The following external RF filters are required: proton bandpass, ^{13}C bandpass, and ^{15}N low pass. The channel isolation required between X and Y (here ^{13}C and ^{15}N) is usually sufficient with a bandpass on one of the channels, but a filter to remove the proton

decoupling RF interference is required for X and Y. This means that one of the band pass filters on X or Y may be replaced by a proton reject, X low pass filter. If the channel isolation between X and Y is not adequate, the probe cannot be tuned.

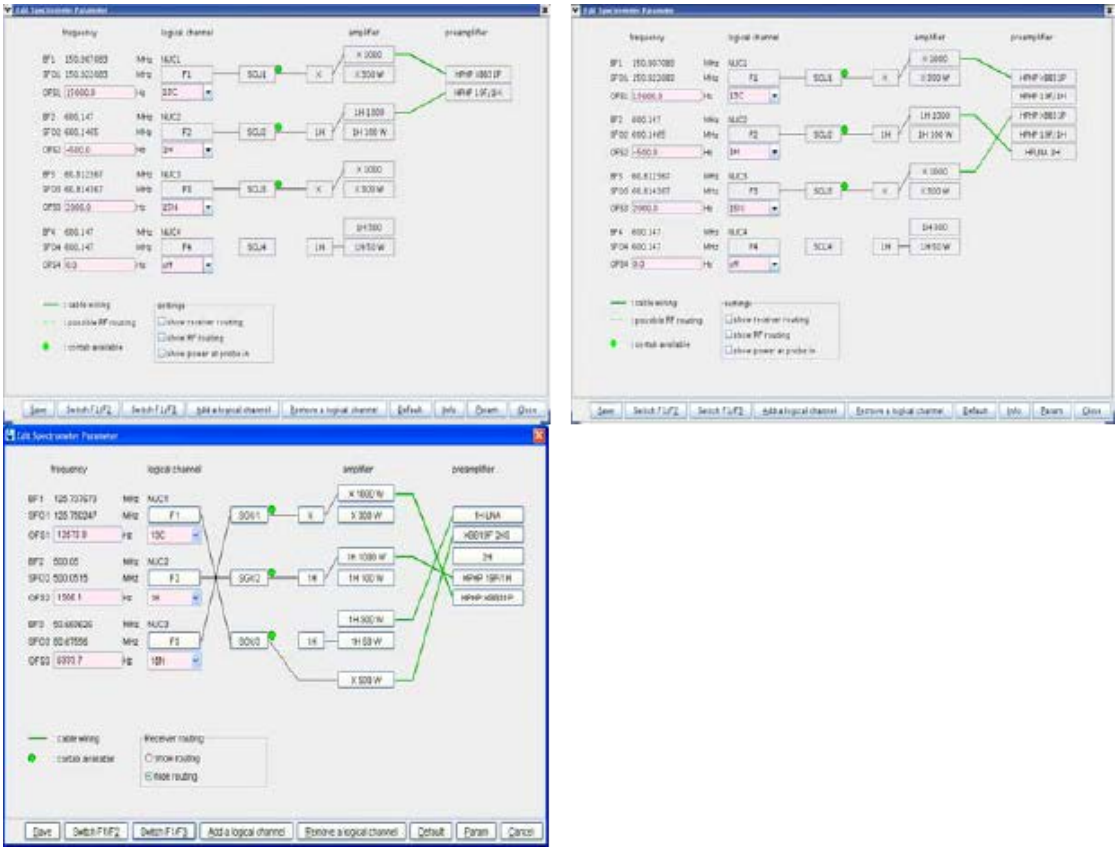


Figure 20.2: The edasp routing tables for H-C-N double CP

Three examples are shown: Setup with only one X-HP-preamplifier (must be recabled for ^{13}C and ^{15}N setup), setup with 2 X-BB HP-preamplifiers and 2 HP transmitter, and setup with one HP transmitter and one 500 W transmitter. The higher frequency nucleus is set for the lower power amplifier.

- Set up for standard ^{13}C CP operation in triple mode. Remember that a double tuned probe has better signal to noise and requires less power on X than a triple probe.
- Optimize decoupling and CP condition, run a reference ^{13}C CP/MAS spectrum of the labelled glycine sample, using 16 scans. This reference spectrum will serve to measure the efficiency of the DCP magnetization transfer.
- To set up the conditions for the N to C transfer, one must define the RF field at which the transfer is to take place, and find the appropriate power levels to achieve these RF fields. In order to minimize losses due to insufficient excitation bandwidths and $T_{1\rho}$ relaxation, the contact should be executed at high power. However, there are limitations in terms of what the probe can take, and there are losses due to unwanted HH contact to the proton spin system. On the other hand, in many samples (bio-samples) the spread of chemical shifts that one wants to cover is not extremely wide or one even wants to execute the transfer selectively (“Specific-CP”). An RF field of 35 kHz is a decent compromise. So, using the **cp90** pulse program, and moving the carrier close to the C_α-peak, determine a power level **pl11** which corresponds to 35 kHz RF (7.14 µsec 90 ° pulse). However, since the sample spins at 11 kHz, the HH condition will require 46 or 24 kHz on one nucleus and 35 kHz on

the other nucleus. If you decide to account for the spin rate on the 13C side, calculate the required power levels for 24 and 46 kHz (= 35 kHz RF field +/- 11 kHz spin rate) RF field using **xau calcpowlev**. Now the ^{13}C channel is set.

20.2.2 15N Channel Setup

- Create a new experiment using **edc**.
- Setup the proper routing by going into **edasp**. Click the switch F1/F3 button to get ^{15}N on channel 1.

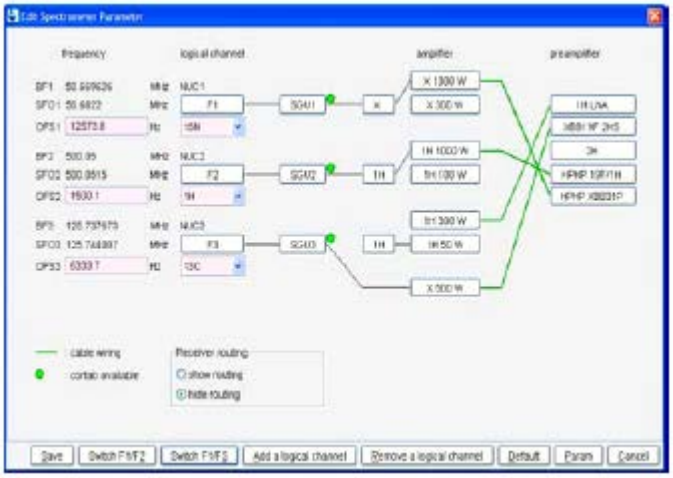


Figure 20.3: Routing table for triple resonance setup change for 15N pulse parameter measurement and CPMAS optimization

- If available, set **pl1** and **sp0** for a proton/^{15}N HH condition in triple mode. On a labelled sample, even the previous settings for ^{13}C should give a signal which allows optimizing the HH condition.
- When the HH condition is optimized, find the power level to achieve a 35 kHz RF field (7.14 µsec 90° pulse, carrier close to the ^{15}N-resonance). It is essential to optimize the first proton to nitrogen HH contact. This is not as trivial as one might think, since the transfer efficiency depends strongly on the timing and RF fields of the HH match. The proton $T_{1\rho}$ of the glycine NH_2 protons (from which the nitrogen is polarized) is fairly short, so the polarization transfer is not efficient. On a fully labelled sample, a maximum enhancement factor of 8.3 is possible (5 protons transfer to one nitrogen). Comparing the cross polarized ^{15}N spectrum to the directly observed spectrum (using **hpdec** and 90 degree pulses at 4 sec repetition) one can measure the enhancement factor rather easily. Without optimization, the enhancement factor may be as low as 5 fold. It should be at least 6.5 fold, more than 7.5 fold is hard to achieve. To achieve a good result, the HH RF-fields should be set as high as possible with a contact time of 4 msec (higher proton RF fields yield a longer proton T1ρ and allow longer contact times). Of course, the RF field is limited by transmitter power and probe breakthrough limits. Note the optimum power levels (**sp0** and **pl1**) and contact time (**p15**).

20.2.3 Setup of the Double CP Experiment

- Read the reference carbon data set and generate a new data set using **edc** or **iexpno**.
- Select the pulse program doubcp. Set the optimum ^{15}N cp parameters as found in the previous step (set **sp0**, **p15** and **pl3** for proton to ^{15}N cp). Set o3 close to the ^{15}N peak position.

- Now we have to select the appropriate parameters for the nitrogen to carbon magnetization transfer. In the standard **doubcp** pulse program, **p16** is used as the second contact time, and the shapes **sp1** and a square pulse at **pl5** are specified for the ^{13}C and ^{15}N contact, so **p16**, **sp1** (^{13}C power level) and **pl5** (^{15}N power level) are the relevant parameters. The C-N contact consists of a square pulse on one channel and a ramp or adiabatic shape on the other channel. Using a square pulse on the ^{15}N channel is preferred, but the sequence can be rewritten to use a ^{13}C square pulse and have the shape on the ^{15}N channel.

Parameter	Value	Comments
Pulse program	doubcp	AVIII, Topspin 2.1 only, else use doubcp, doubcp.av
nuc1	^{13}C	Nucleus on f1 channel
o1p	100 ppm	^{13}C offset
nuc2	^{1}H	Nucleus on f2 channel
o2p	2-4 ppm	^{1}H offset, optimize
nuc3	^{15}N	Nucleus on f3 channel
o3p	≈35 glycine ≈120 histidine ≈65 -130 protein	^{15}N offset depending on sample
sp1		Power level for f1 channel, NC contact pulse
pl3		Power level for ^{15}N channel HN contact
pl5		Power level for ^{15}N channel, NC contact pulse
pl12		Power level decoupling f2 channel and excitation
pl13		power level during second contact, cw dec.
cnst24		offset for cw decoupling during p16
p3		Excitation pulse f2 channel
pcpd2		Decoupler pulse length f2 channel (^{1}H) TPPM
p15	3-5 msec	Contact pulse – first contact
p16	5-12 msec	Contact pulse – second contact f1 – f3 channel
d1	5-10s histidine 4s α-glycine	Recycle delay
spnam0		Ramp for 1st CP step; e.g. ramp: 80 – 100%
sp0		Power level for Ramp HN contact pulse ^{1}H
spnam1	ramp45-55, tcn5500, or square.100	ramp, tangential contact pulse **tcn5500** **on C, or square**
spnam2	square.100 or ramp45-55, tcn5500	square on N, or ramp/tangential pulse
cpdprg2	SPINAL64	SPINAL64 decoupling
ns	2, 4, or 16	Number of scans

Table 20.1: Recommended Parameters for the DCP Setup

- Now we select the shape to use for the C-N contact. To find the HH contact more rapidly (it is a very narrow condition) it is recommended to start with a ramp shape. In order to find a HH condition independent of the type of shape, it is recommended to select shapes which are all centred around 50% amplitude, which allows arbitrary amplitude modulation without changing the HH condition. For a start, generate a ramp shape from 45 to 55% with 100 slices, using shape tool (**stdisp**). Store the ramp as ramp4555.100. Select this ramp as **spnam1**.

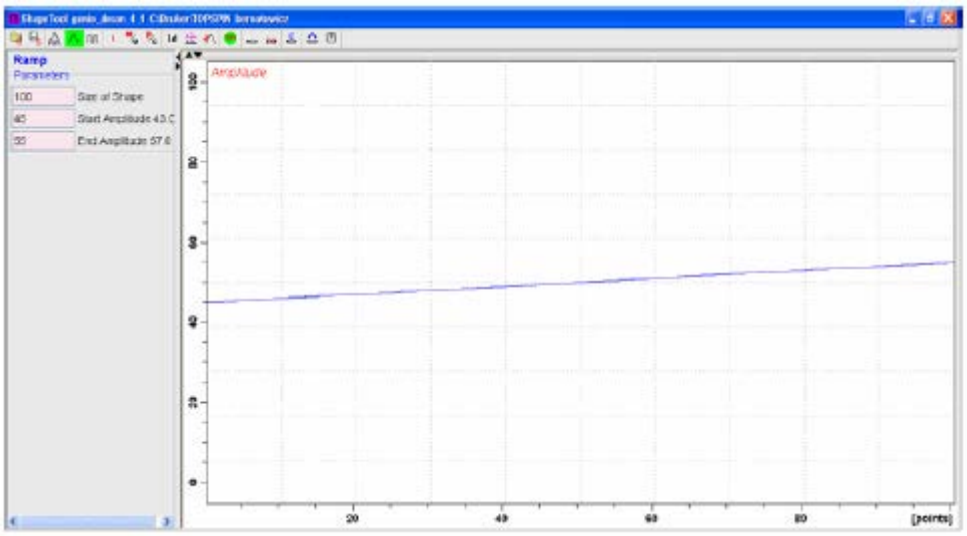


Figure 20.4: Shape Tool display with ramp shape from 45 to 55%.

The amplitude factor is 50%, corresponding to 50% RF field or a power level change of 6 dB, since the amplitude corresponds directly to the pulse voltage.

- Since the shape is centered around 50%, the RF voltage here is down by a factor of 2 (= 6 dB in voltage), the power must be increased by 6 dB to get the same RF field as with a 100% square pulse.
- Usually, the ramp shape is set on ^{13}C (but it can also be used on 15N). Set **pl5** to 46 or 24 kHz RF field on ^{15}N. Set **spnam1** to **ramp4555.100, sp1** to 35 kHz RF field on ^{13}C - 6 dB. Set **pl13=pl12** for a start, set **p16** to 5 msec. Optimize the power level **pl5**. A variation over -1 to +1 dB in steps of 0.2 dB should be ample. In order to be sure, one can optimize **sp1** and **pl5** as an array, **sp1** in steps of 0.5 dB, **pl5** in steps of 0.2 dB. Use a full phase cycle to avoid signal from a direct proton to carbon transfer (which is cancelled by the phase cycle). Optimize **p16** between 5 and 15 msec. See the figure below for an optimization of **pl5** (^{15}N square pulse power).
- With a ramp shape for the N-C transfer, one should get 40-50% DCP efficiency (see figure Double CP Yield below), compared to the reference direct ^{13}C CP spectrum. If this cannot be achieved, even with careful HH matching, the following parameters should be checked:
- Re-optimize **p16**, the optimum should be > 10 msec. If the signal gets worse with longer contact time, there is a loss due to direct ^{13}C-^{1}H contact. Minimize this loss in the following way:
- Never use a pulsed proton decoupling schemes during p16. Frequency shifted Lee-Goldburg decoupling is no alternative, since the signal will broaden and decay with shorter T_2. Use **cw** decoupling during **p16**, and carefully optimize the decoupling power (**pl13**) for maximum signal. A slight offset may be set using **cnst24**.

- Instead of a 45-55% ramp, a tangential amplitude modulation shape can be used. Since this shape provides 100% transfer efficiency on a spin pair system (compared to 50% of a standard rectangle or ramp shape), the DCP efficiency can be increased. With such a shape, one can get 50-70% DCP efficiency. To generate such a shape in **stdisp**, select TanAmpMod as a shape model, select solids notation, select 1000 points, set the spin rate to half the actual spin rate (5500), set the RF field to the actual RF field used on this channel, select 400 for the dipolar coupling, and 50% for the scaling factor. Save the shape as tcn5500 (if not already available), and select this name for **spnam1**. The efficiency should be noticeably better. Reoptimize the first HH contact, decoupling and **p16**. More than 50% should definitely be obtained.

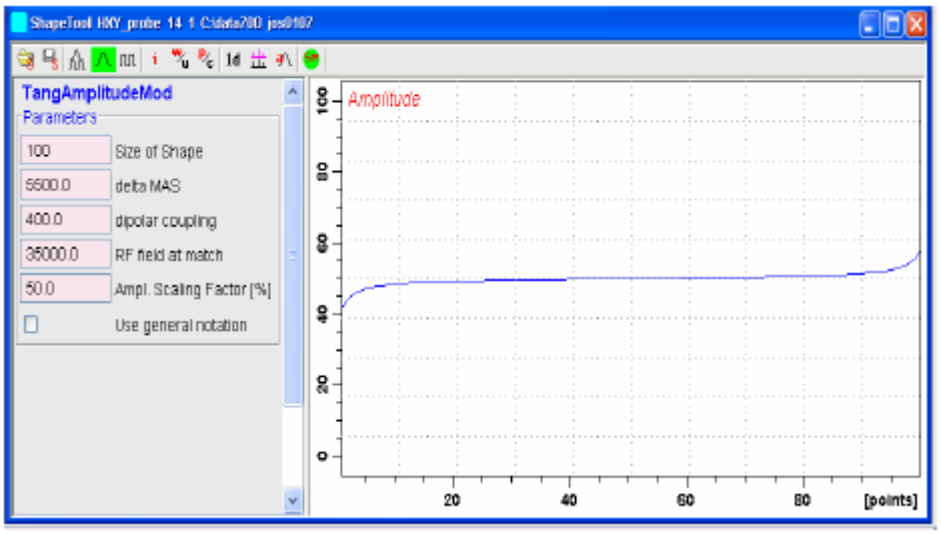


Figure 20.5: Shape Tool display with a tangential shape for adiabatic cross polarization

The amplitude factor is 50%, corresponding to 50% RF field or a power level change of +6 dB (4 fold power in TopSpin 3.0 and later), since the amplitude corresponds directly to pulse voltage.

- Optimize the DCP condition on the rectangular pulse, using 0.1 dB steps over a range of +/-1 dB around the optimum found with the ramp.

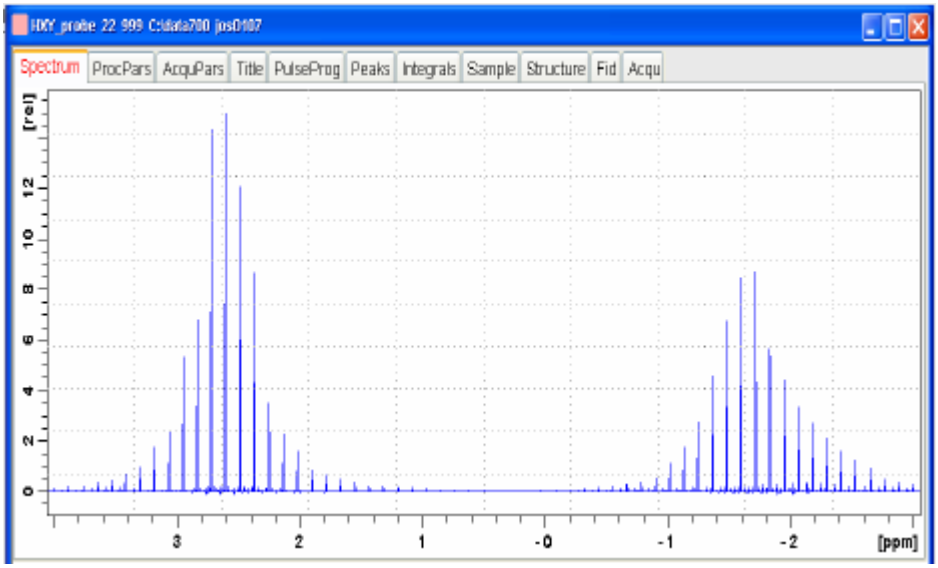


Figure 20.6: Double CP optimization of PL5 in increments of 0.1 dB

Note how narrow the optimum DCP conditions are. However, with diligent preparation, one should be very close to the optimum with the first try.

- Run an experiment with 16 scans and compare the signal amplitude with the signal intensity of the ^{13}C CPMAS experiment with the same number of scans, using dual display. The intensity ratio of the aliphatic resonance of the CPMAS compared to the one obtained with the DCP experiment gives the DCP yield:

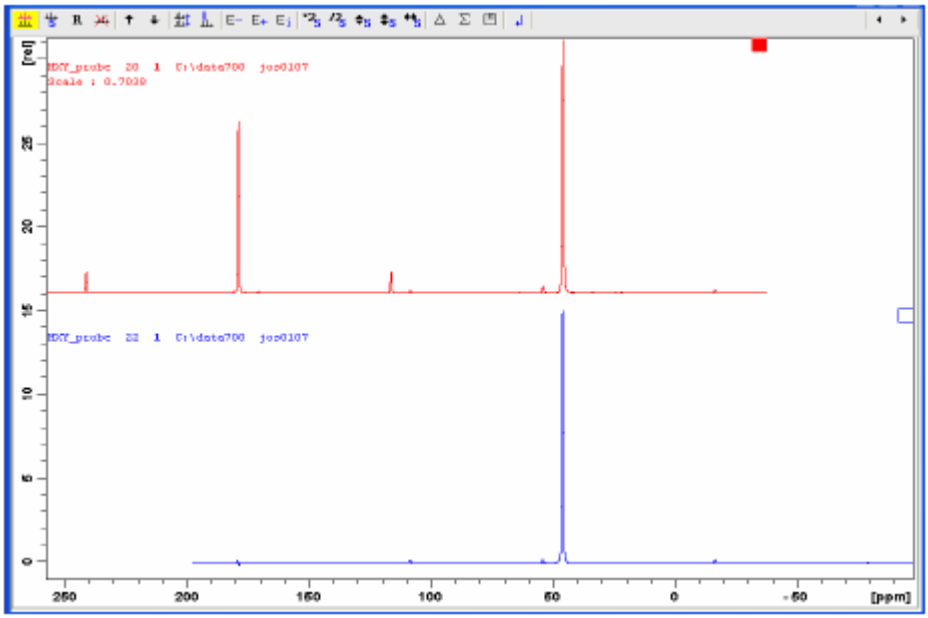


Figure 20.7: Double CP yield, measured by comparing CPMAS and DCP amplitudes of the high field resonance

Note that the C_1 carbon receives very little magnetization under these conditions, the transfer is rather selective.

The setup for DCP can be rather much sped up and simplified by a python program named dcpset.py. This program will ask for the 90° pulse widths and associated power levels and for the spin rate and calculate the appropriate power levels for the HH condition for all pulses except for the proton channel.

Contact *solids@bruker.de* to receive this python program and some instructions for use. This program serves as an example on how the experiment setup can be controlled by the high level script language.

20.2.4 Setup of the 2D Double CP Experiment

- Load a suitable sample, spin it up, set the desired temperature and match and tune the probe. As a simple setup sample, full ^{13}C, ^{15}N-histidine may be used (**d1**=10s, 2-4 scans, **p15**=1 msec, **p16**=3 msec).A labeled oligopeptide or small protein will of course provide a more interesting spectrum. With proteins, good results should only be expected if the preparation is micro-crystalline. In such a case, water, salt and cryo-protectant (glycol, glycerol) will very likely be present. This means that the probe proton channel will be detuned to lower frequency, and tuning may be difficult, if not impossible at high proton frequencies and salt contents. In such cases, E^{free} probes are recommended.
- Run standard 1D cp ^{13}C and ^{15}N experiments; determine the required offsets for all frequencies and the required sampling windows.
- Re-optimize the H-N and N-C HH conditions.
- Generate a new data set and switch to 2D data mode, using the “123”-icon in **eda**.
- In **eda**, set the pulse program to **doubcp**. Set FnMode as desired, usually STATES-TPPI.
- Make sure the correct nucleus (^{15}N) is selected in the F1 dimension.
- Set the sampling windows for both dimensions from the previously acquired 1D spectra.
- Both acquisition times in F2 and F1 should be considered with care, since the decoupler is on at high power during both periods. Especially for biological samples, where the RF heating may be high and the samples are temperature sensitive, it is essential not to use overly long acquisition times and high duty cycles. Remember that the heating effect is generated inside the sample where the temperature increases within milliseconds, whereas cooling requires transfer of the energy to the outside of the spinner, which takes seconds! E^{free} probes eliminate these problems to a large extent.
- The basic double-CP experiment can be extended into many different variations. One example is the double transfer N-C_α-C_β, where the second transfer step is made selective so only α-carbons are polarized from the nitrogen, then magnetization is transferred from the α-carbons to the adjacent β-carbons. This can be done by a simple PDSD or DARR proton spin diffusion step, or by a ^{13}C-^{13}C homonuclear recoupling step (HORROR, DREAM, or other). Likewise, the N-C_α-C_x experiment transfers from the α-carbons to all (X) carbons which are in close enough proximity to the α-carbons. Check with your applications support for appropriate pulse programs.

20.3 2D Data Acquisition

Sample: ^{15}N, ^{13}C-labeled histidine, peptide or protein.

Spinning speed: 10 – 15 kHz, depending on ^{13}C spectral parameters (rotational resonance must be avoided)

Experiment time: 30 minutes to several hours

Acquisition Parameters:

Parameter	Value	Comments
Pulse program	doubcp	Pulse program
nuc1	^{13}C	Nucleus on f1 channel
o1p	100 ppm	^{13}C offset
nuc2	^{1}H	Nucleus on f2 channel
o2p	2-4 ppm	^{1}H offset, optimize
nuc3	^{15}N	Nucleus on f3 channel
o3p	65 – 150 ppm	^{15}N offset depending on sample
pl1		Power level for f1 channel, NC contact pulse
pl3		Power level for ^{15}N channel HN contact
pl5		Power level for ^{15}N channel, NC contact pulse
pl12		Power level decoupling f2 channel and excitation
p3		Excitation pulse f2 channel
pcpd2		Decoupler pulse length f2 channel (^{1}H) TPPM
p15	1-5 msec	First contact, optimize on ^{15}N cp spectrum
p16	3-10 msec	Second contact f1 – f3 channel, optimize on 1D dcp spectrum
d1	5-10 s for histidine	Recycle delay, optimize on 1d
spnam0		Ramp for 1st CP step; e.g. ramp: 80 – 100%
sp0		Power level for Ramp HN contact pulse ^{1}H
spnam1	tcn5500	Tangential or ramp contact pulse
spnam2	square.100	Shape on ^{15}N channel
cpdprg2	SPINAL64	SPINAL64 decoupling
ns	2 or 16	Number of scans
F2 direct ^{13}C		(left column)
td	2k	Number of complex points
sw	≈200 ppm	Sweep width direct dimension, adjust to experimental requirements
F1 indirect ^{15}N		(right column)
td	128-.512	Number of real points
sw	≈100-150 ppm	Sweep width indirect dimension

Table 20.2: Recommended Parameters for the DCP 2D Setup

20.4 Spectral Processing

Processing parameters

Parameter	Value	Comment
F1 acquisition ^{13}C		(left column)
si	2-4k	FT-size
wdw	QSINE	Squared sine bell
ssb	2-5	Shifted square sine bell, >2: res. enhancement
ph_mod	pk	Phase correction if needed
F2 indirect ^{15}N		(right column)
si	512-1024	Zero fill
mc2	STATES-TPPI	
wdw	QSINE	Squared sine bell
ssb	2-5	

Table 20.3: Recommended Processing Parameters for the DCP 2D

20.5 Example Spectra

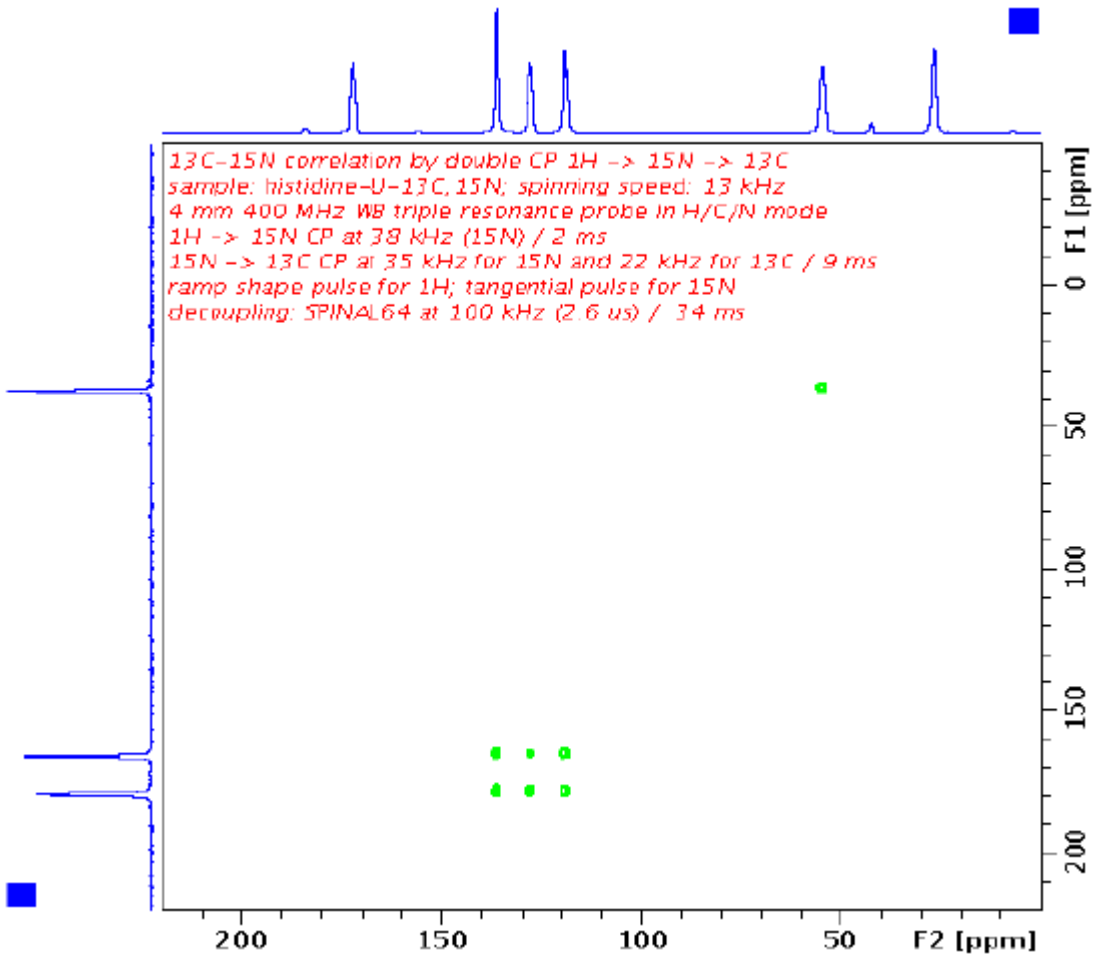


Figure 20.8: C-N correlation via Double CP in histidine (simple setup sample). 4 mm Triple H/C/N Probe.

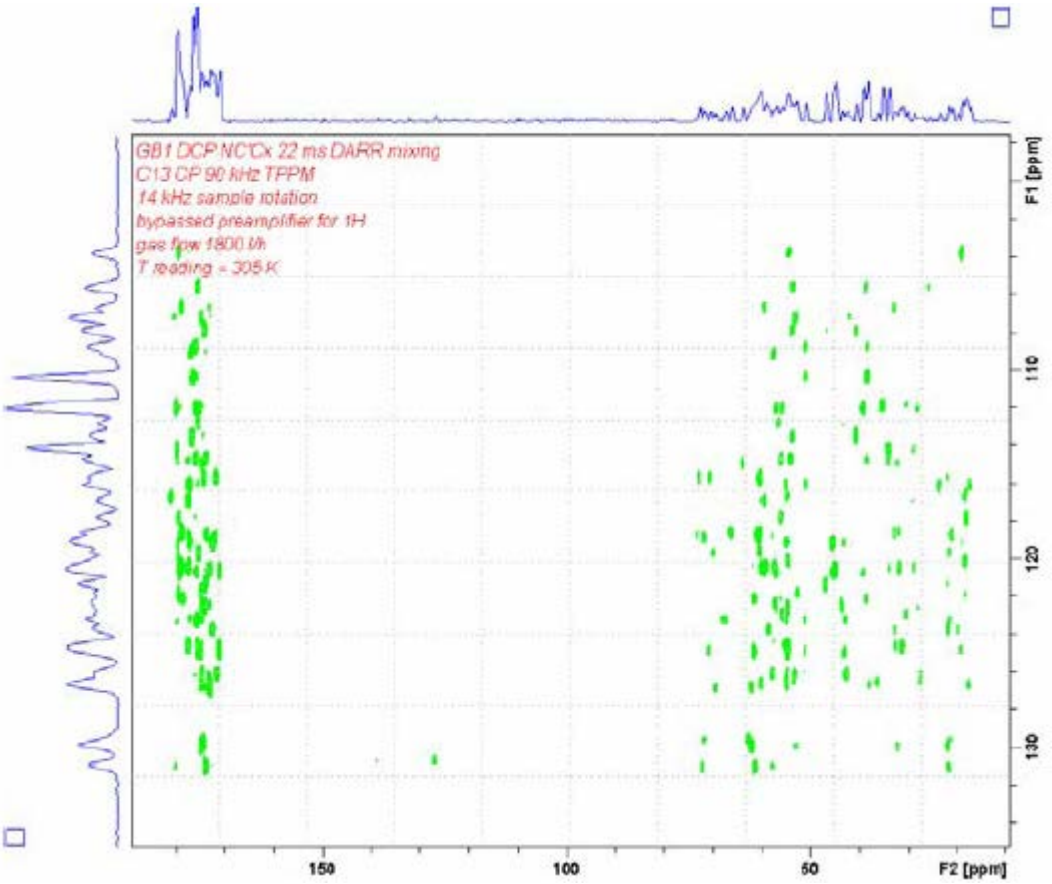


Figure 20.9: NCαCx correlation experiment with 22 ms DARR mixing period

In the period above is an $NC_{\alpha}C_x$ correlation experiment with 22 ms DARR mixing period for C_{α}-C_x spin diffusion on GB1 protein run using an E^{FREE}-Probe.

DARR transfer from C_{α} to C_{β} or C_x generates positive cross peaks, HORROR or DREAM transfer generates negative cross peaks. See the chapter on spin diffusion experiments for more information about DARR or PDSD.

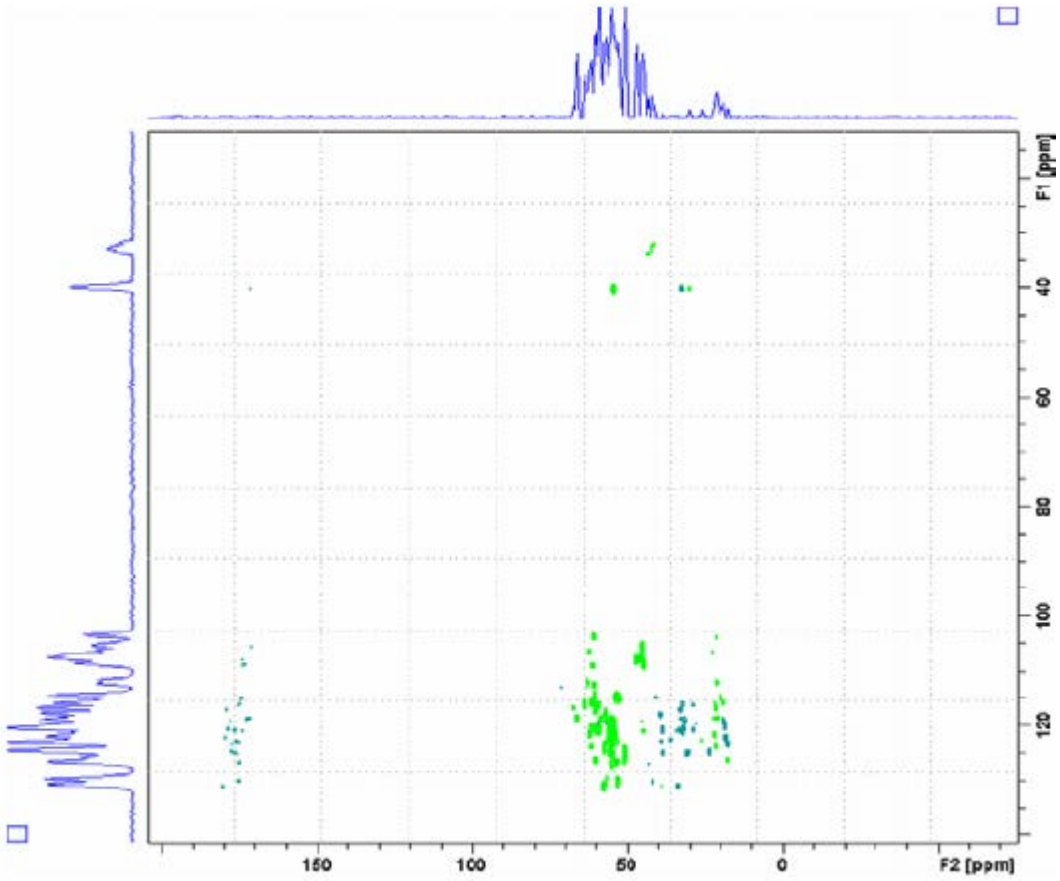


Figure 20.10: NCaCx correlation experiment with 4.2 ms SPC5-DQ missing period

In the figure above is an NC_aC_x correlation experiment with 4.2 ms SPC5-DQ mixing period for C_aC_x spin diffusion on GB1 protein run using an E^{FREE}-Probe at 14 kHz sample rotation and 100 kHz decoupling.

See the chapter *Setup for the Recoupling Experiment [▶ 165]* for the SPC5 setup and for more information about DQ recoupling sequences. Note the inverse phase of the cross peaks generated by the DQ-mixing step.

21 CRAMPS: General

CRAMPS is an acronym standing for **C**ombined **R**otation **A**nd **M**ultiple **P**ulse NMR **S**pectroscopy. Multiple Pulse Spectroscopy had long been thought not to work under spinning around the magic angle, but in fact it does work, as long as the pulse cycle times are substantially shorter than the rotation period.

CRAMPS suppresses homonuclear dipolar interactions between the abundant spins (mostly protons) and chemical shift anisotropy simultaneously through the combination of multiple pulse techniques and magic angle spinning. J-couplings and large heteronuclear dipolar couplings are not suppressed.

Reference

L.M. Ryan, R.E.Taylor, A. J. Patt, and B. C. Gerstein, *An experimental study of resolution of proton chemical shifts in solids: Combined multiple pulse NMR and magic-angle spinning*, J. Chem. Phys. 72 vol.1, (1980).

21.1 Homonuclear Dipolar Interactions

Homonuclear dipolar interactions among spins with a strong magnetic moment and high natural abundance - mainly ^{1}H or ^{19}F, and to a much smaller extent ^{31}P - are usually very large unless averaged by high mobility. Especially in the case of protons, spin exchange is usually rapid compared to routinely achievable rotation periods, meaning that MAS alone cannot suppress the homonuclear dipolar broadening. Even spin rates in the order of 70 kHz, which is no longer a mechanical problem, cannot fully average this interaction in rigid solids. As chemical shift differences among the coupled nuclei become larger, the interaction becomes more heterogeneous and MAS can suppress it more efficiently. This is the reason why fast spinning alone works much better on ^{19}F or ^{31}P than on protons. heteronuclear dipolar coupling, such as between ^{13}C and ^{1}H, can in principle be spun out, but only if the homonuclear coupling between protons is small, or averaged by motion or a suitable pulse sequence. CRAMPS sequences therefore play an important role also in experiments where X-nuclei are observed.

21.2 Multiple Pulse Sequences

Dealing with a heteronuclear dipolar coupling is easy: continuous high power irradiation of one coupling partner will decouple it from the other nucleus, as in the case of ^{13}C observation while decoupling protons. However observing a nucleus while decoupling it from like spins at the same time is obviously not trivial, since the signal cannot be observed under the much higher decoupling RF. Observation of the signal and decoupling pulses must therefore be alternately applied. Suppression of a homonuclear dipolar interaction occurs when the magnetization vector of the coupled spins is tilted into the magic angle. This condition can be achieved either by 4 л/2 pulses of suitable phase and spacing (multiple-pulse methods), or by off-resonance irradiation of suitable offset and RF-field (Lee-Goldburg). To observe the signal, a gap within the pulse sequence must be supplied, which is long enough to observe one or several data points while the magnetization vector points along the magic angle. This condition obviously persists only for a time period short compared to the transverse relaxation of the signal. To observe the time dependence of the signal, the sequence must be repeated and more data points accumulated until the signal has decayed under the influence of residual broadening. Obvious problems of this experiment are the requirement to observe a relatively weak signal shortly after a strong pulse (dead time problem) and the requirement to time the sequence in such a way that the magnetization vector is accurately aligned with the magic angle (requires precise pulse lengths and phases, and it requires RF fields strong

compared to the interaction and shift distribution). Many sequences have been devised after the original WHH-4 (or WaHuHa) sequence which yield better results due to better error compensation (MREV-8, BR-24, C-24, TREV-8, MSHOT). Modern hardware has made the setup and application of these sequences a lot easier since pulse phase and amplitude errors are negligible, higher magnetic fields have led to better chemical shift dispersion and also to shorter dead times. The resolution achieved with long, highly compensated sequences like BR-24 is very good, but their applicability at limited spin rates (because of the need for the cycle time to be short with respect to the rotor period) often presents a problem.

References

S. Hafner and H.W. Spiess, *Multiple-Pulse Line Narrowing under Fast Magic-Angle Spinning*, J. Magn. Reson. A 121, 160-166 (1996) and references therein.

M. Hohwy, J. T. Rasmussen, P. V. Bower, H. J. Jakobsen, and N. C. Nielsen. *1H Chemical Shielding Anisotropies from Polycrystalline Powders Using MSHOT-3 Based CRAMPS*, J. Magn. Res.133 (2), 374 (1998), and references cited therein.

21.3 W-PMLG and DUMBO

W-PMLG and DUMBO are shorter sequences which avoid turning high power pulses rapidly on and off, which is what most multiple pulse sequences do. This avoids undesired phase glitches. Also, they use higher duty cycles during the decoupling period. As a result, the sequences are simpler and shorter, requiring fewer adjustments and allowing higher spin rates.

Both sequences use repetitive shaped pulses with detection in between. PMLG uses the principle of a Frequency Switched Lee Goldburg (FSLG) sequence (continuous irradiation with a net RF field along the magic angle), where the frequency shifts are replaced by a phase modulation. DUMBO basically works like a windowless MREV-type pulse sequence where the individual pulses are replaced by a single pulse with phase modulation.

References

E. Vinogradow, P.K. Madhu, and S. Vega, *High-resolution proton solid-state NMR spectroscopy by phase modulated Lee-Goldburg experiment*, Chem. Phys. Lett. 314, 443-450 (1999).

D. Sakellariou, A. Lesage, P. Hodgkinson and L. Emsley, *Homonuclear dipolar decoupling in solid-state NMR using continuous phase modulation*, *Chem. Phys. Lett. 319, 253 (2000).*

21.4 Quadrature Detection and Chemical Shift Scaling

Under homonuclear decoupling, the magnetization processes in the transverse plane of a tilted rotating frame whose new z-axis is along the direction of the effective field. The projection of this plane into the X-Y plane is therefore an ellipse, and the signal intensities sampled in the two quadrature channels along the x and y direction are different, since these are the major and the minor axes of the projected ellipse. As can be seen from the figure, the X and Y observe direction will see a signal of different amplitude. This means that quadrature images will always be present, if quad detection is used. In the case of single detection, the signal may be smaller or larger depending on the receiver phase. The standard procedure is to use quad detection and suppress the quad images by a suitable phase cycling scheme. This is however not as straightforward as it is with standard excitation/observation. The quad phase cycling must occur in the precession plane, so a prepulse is required to tilt the initial magnetization into the direction of the precession axis. Usually, a combination of 2 pulses is used for initial excitation ($90°_{x,y,-x,-y}$ + ($90°$ -$55°$) $_{adjust}$).

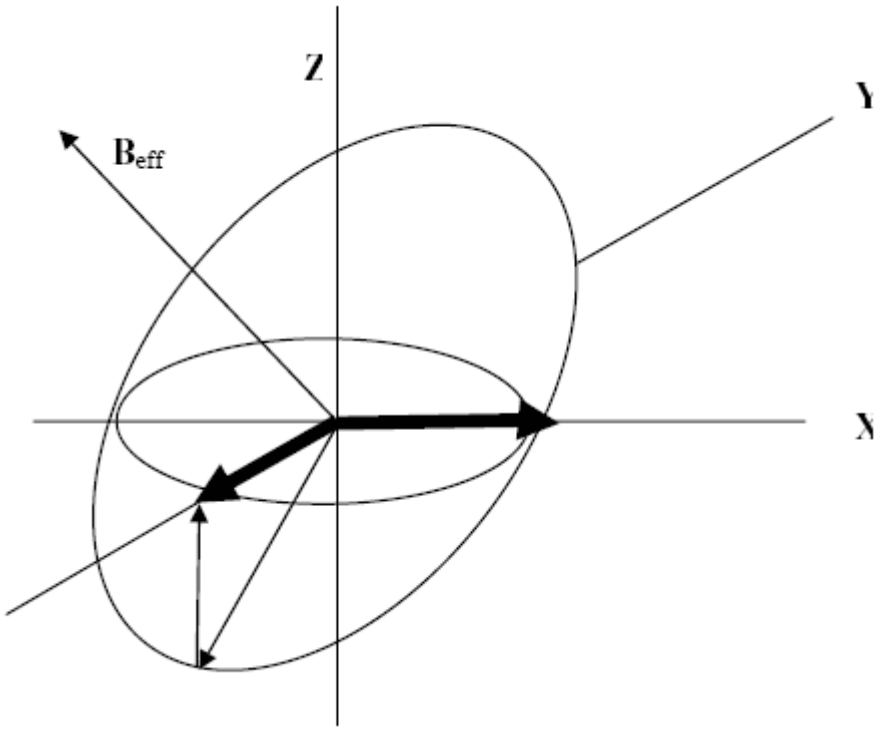


Figure 21.1: Difference in Amplitude of the Quadrature Channels X and Y

The difference in amplitude of the quadrature channels X and Y, caused by the tilted precession plane. Along X, the full amplitude is observed, along Y only the component in the XY-plane is detected.

As the spins process around a tilted effective field and not only around the direction of the external field, the precession frequencies are changed, which means that the observed chemical shifts are changed. As the frequencies are always smaller than in the standard excitation/observation scheme, the chemical shift range appears scaled down. The scaling factor depends on the pulse sequence used. To achieve a spectrum comparable to spectra acquired conventionally, the shift range must be scaled up again by this scaling factor, i.e. the spectral window given by the repetition rate of the pulse sequence must be multiplied by this scaling factor in order to place the resonances correctly. This scaling factor can be calculated from the tilt angle, but is also slightly dependent on the offset and RF-field. Since the correct chemical shifts are usually unknown, one must be aware of the fact that the shifts may not be as precise as they are in high resolution liquids experiments. An example of shift calibration, taking the scaling factor into account, will be given in the practical chapter.

CRAMPS: General

22 CRAMPS 1D

As outlined above, many sequences are available to achieve homonuclear dipolar decoupling. We want to concentrate on those that allow fast spin rates and are easy to set up. The performance of DUMBO and W-PMLG is very similar. The pulse sequence is also very similar on AV3 instruments, just different shapes and different timings are loaded.

22.1 Pulse Sequence Diagram of W-PMLG or DUMBO

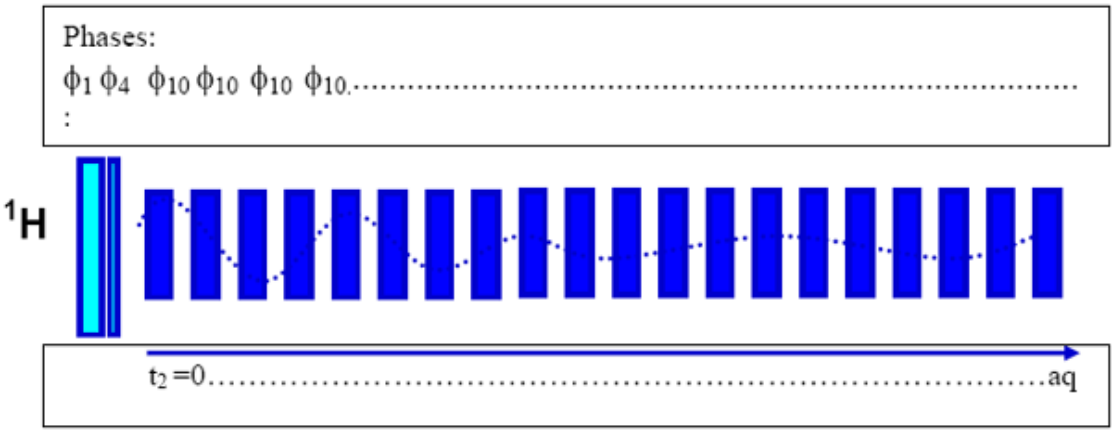


Figure 22.1: Pulse Sequence Diagram

Phases	RF Power Levels	Timing
φ1 = CYCLOPS, 1 2 3 0	pl12=set for around 100 kHz	p1 around 2.5 µsec.
φ4 = 0 +cnst25, adjust	ditto.	p4 about 45 degrees, adjust.
φ10 = 0	sp1: set for 100-130 kHz	WPMLG: p5, 1.2-1.5 µsec or calculated from cnst20=RF field DUMBO: p10 set by xau dumbo.
φ31 = CYCLOPS, 0 1 2 3		

Table 22.1: Phases, RF-Levels, Timings

22.2 Pulse Shapes for W-PMLG and DUMBO

Both shapes are purely phase modulated pulses, their amplitudes are constant throughout. The PMLG shape is a standard shape delivered with the software (**wpmlg1**, **m5m**, **m5p**). DUMBO shapes are generated using the standard AU-program dumbo. Calling dumbo with **xau dumbo** will ask for the slice length of the shape (usually 1 µsec), the number of slices (usually 32), generate the shape and load the name of this shape into the current parameter set, it will also set the length of the shaped pulse, **p10**, to 32 µsec. Note: With magnetic fields higher than 500 MHz, it is recommended to replace the standard 32 µsec timing by 24 µsec timing and increasing the power level accordingly, since this has been found to give better results.

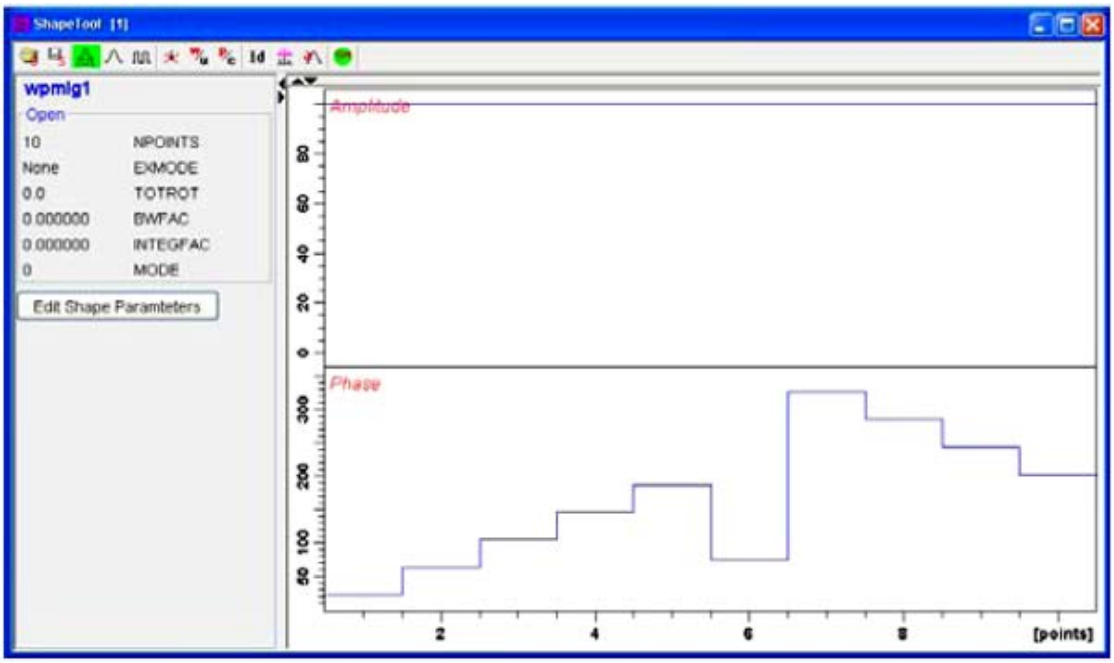


Figure 22.2: PMLG Shape for wpmlg, sp1

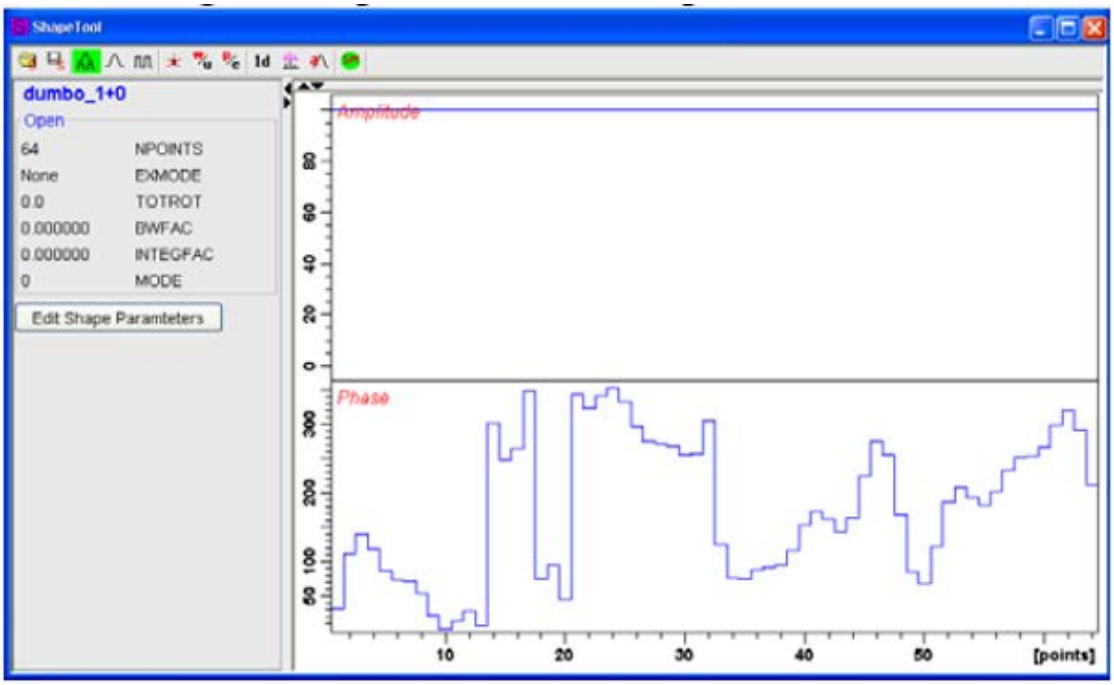


Figure 22.3: Shape for DUMBO, sp1

22.3 Analog and Digital Sampling Modi

AV3 instruments allow different acquisition modi, one which resembles the previous mode of analogue filtering in so far as the down-conversion is done without simultaneous digital filtering, whereas the digital mode always down converts and filters simultaneously. Remember that at a standard sampling rate of 20 MHz (the fixed sampling rate of the DRU) down-conversion must be done to obtain data sets of reasonable sizes. The pulse programs dumboa and wpmlga are written for the pseudo analog mode without digital filtering, **dumbod** and **wpmlgd** are written for the digitally filtered mode.

22.3.1 Analog Mode Sampling

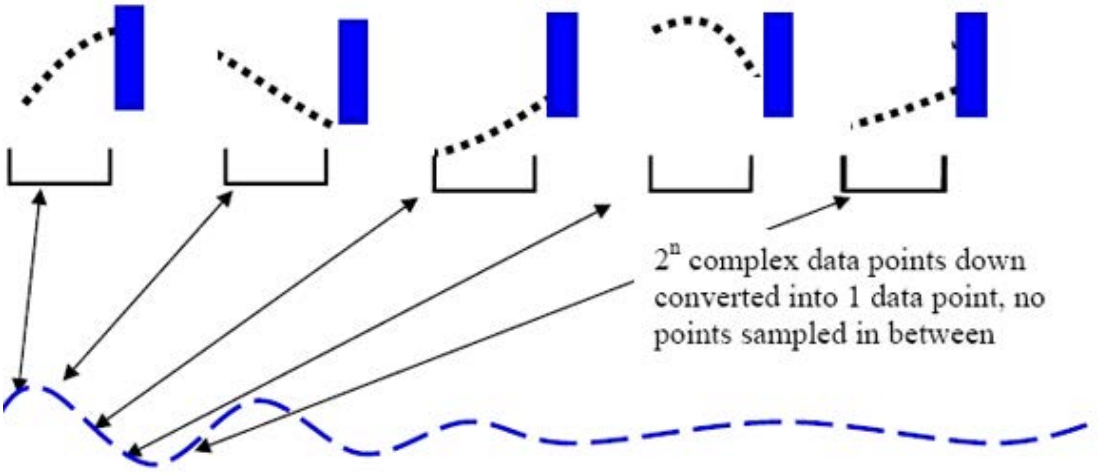


Figure 22.4: Analog Sampling Scheme

22.3.2 Digital Mode Sampling

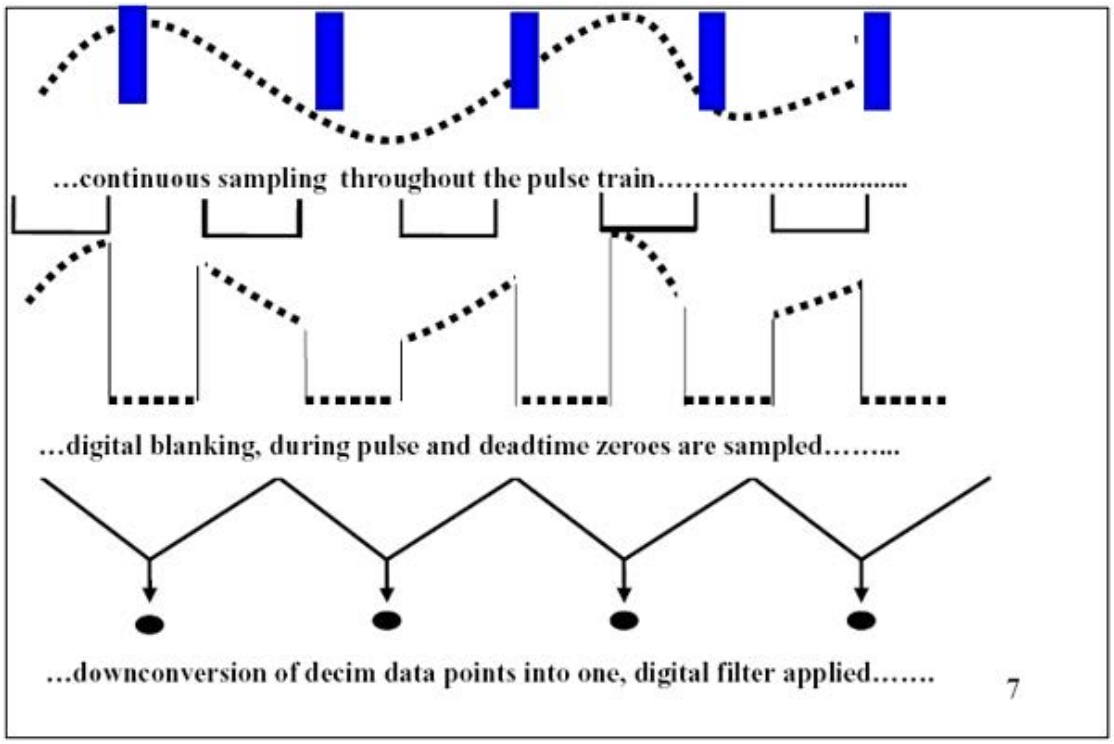


Figure 22.5: Digital Sampling Scheme

22.4 Setup

At frequencies of 400 MHz and higher, double or triple resonance CP-MAS probes may be used on the proton channel, at lower fields a CRAMPS probe is required due to the increased ring down time at lower field. Spinner diameters of 4 mm or smaller are preferred, since we want to spin over 10 kHz. Since only one nucleus is observed, no filters are required and should be avoided. Good impedance matching between probe and transmitter is important in order to optimize the effect of the pulses on the spins. If the RF cable has been too strongly bent or the connectors been twisted, the cable may not have 50 Ohms and the result will always be bad. Likewise, if the preamp is burnt, it is not possible to get good results.

The fewer connectors are between probe and preamp, the better you can expect the 50 Ohm match to be. In **edasp**, set F1 for ^{1}H-observation, select the high power proton amplifier and high power preamplifier. Tune and match as usual. It is assumed that the magic angle is precisely set, which can easily be achieved with KBr on a double resonance probe, or on $BaClO_3$*H_2O, looking at the proton signal, much like one does on the Br^{79} resonance.

Shimming will also be important, since protons are observed, and on some samples, good resolution is expected. Looking at the protons in adamantane, find the power level for a 2.5 µsec 90° pulse. Set the B_0-field or **o1** to be close to resonance (see chapter Basic Setup Procedures for more details). Calibrate the adamantane proton shift to 1.2 ppm. Then load a spinner with α-glycine (precipitate from cold water with acetone and dry, if you are not sure about the composition of your sample). A spinner with 50 µl or less sample volume is preferred since high H_1-homogeneity is desired, although it is by far less important than is commonly stated in the literature.

22.5 Parameter Settings for PMLG and DUMBO

Parameter	Value	Comment
pulprog	wpmlga	Runs on AV 3 instruments only.
pl12 sp1	for 100 kHz RF field dto, set cnst20=100 000	 To be optimized during setup.
spnam1	wpmlg1, m5m or m5p	
p1	2.5 µsec	As for 100 kHz RF field.
p8	1.2µsec	
p14	0.7 µsec	To be optimized.
cnst25	140	To be optimized.
p9	4 – 2.6 µsec	To be optimized.
p5	1.5 µsec or calculated from cnst20	To be optimized.
d1	4s	For α-glycine.
l11=anavpt	4	2, 4, 8, 16 or 32.
o1p	10 or -1	To be optimized.
swh	1e6/2*(2*p9+10*p5)*0.6	To be corrected for proper scaling.
rg	16-64	
td	512	Up to 1024.
si	4k	
digmod	analog	
MASR	12-15 kHz	Depending on cycle time.

Table 22.2: PMLG Analog Mode

Parameter	Value	Comment
pulprog	dumboa	Runs on AV 3 instruments only.
pl12 sp1	for 100 kHz RF field up to 130 kHz	 To be optimized during setup.
spnam1	dumbo1_64	Set by xau dumbo.
p1	2.5 µsec	For 100 kHz RF field.
p8	1.2µsec	
p14	0.7 µsec	To be optimized.
cnst25	140	To be optimized.
p9	4 – 2.6 µsec	To be optimized.
p10	32 µsec or 24 µsec	Set by xau dumbo.
d1	4s	For α-glycine.
l11=anavpt	4	2, 4, 8, 16 or 32.
o1p	5	To be optimized.
swh	1e6/2*(2*p9+p10)*0.5	To be corrected for proper scaling.
rg	16-64	
td	512	Up to 1024.
si	4k	
digmod	analog	
MASR	10-12 kHz	Depending on cycle time.

Table 22.3: DUMBO, Analog Mode

22.6 Fine Tuning for Best Resolution

For fine tuning, the following parameters are important:

P9 sets the width of the observe window. The shorter it is, the better the resolution. However, the natural limit is the size of the sampling period and the dead time of the probe. Preamp and receiver play no significant role in the total dead time. A CP/MAS probe usually has a fairly narrow bandwidth (long dead time), so **p9** < 3µsec is only possible at frequencies 400 and higher. With **l11** 4-8, **p9** can be chosen shorter for better resolution, but at the cost of S/N. Sampling more data points during d9=**l11***0.1 µsec with larger values of **l11**, will increase S/N slightly but requires more time within the window, may require a longer p9 and therefore degrade resolution.

Since the decoupling bandwidths are not very large, **o1** should be close to resonance, especially for DUMBO. For PMLG, this is less critical. The power level for the shapes should be adjusted in steps of 0.2 dB. The splitting of the two high field lines (the protons in the –CH_2- are in-equivalent in the solid state) should be below the 50% level.

22.7 Fine Tuning for Minimum Carrier Spike

The tilt pulse **p14** and its phase (**cnst25**) determine the size of the carrier spike. Optimize both parameters alternately for minimum spike, and make sure the spike does not overlap with a resonance by selecting **o1** appropriately. Note: changing **o1** will lead to different values for **cnst25**.

22.8 Correcting for Actual Spectral Width

Since the sampling rate is governed by the multi-pulse sequence repetition rate, the foreground parameter **swh** has no real meaning. Once all tuning procedures are done, calculate the real spectral width **swh** according to the formula given in the parameter tables and run a new experiment. After FT, the spectrum should have an approximately correct spectral width. Calibrate the middle position between the two –CH2- peaks to 3.5 ppm, the NH3-peak should then be at about 7.5 ppm. Since the actual peak positions depend on the probe tuning, you will have to recalibrate for your sample using one or more known chemical shifts. If the peak separation is incorrect, change the status parameter **swh** by typing **s swh** and scaling it appropriately. Some pulse programs are written such that upon ased, the (approximately) correct sweep width is shown and can be set as an acquisition parameter.

22.9 Digital Mode Acquisition

Most parameters stay the same as adjusted in analog mode.

Parameter	Value	Comment
pulprog	dumbod or wpmlgd	AV 3 instruments only.
digmod dspfirm	digital sharp or medium	
aqmod	qsim or dqd	
swh	50000-10000	Depending on spectral range and o1.

Table 22.4: Parameters for Digital Mode

The correction for the scaling factor must be done after acquisition, changing the status parameter **swh** by typing **s swh** and dividing the value by the scaling factor (about 0.578 for WPMLG, 0.47 for wpmlgd2 and 0.5 for DUMBO). Some pulse programs are written such as to show the correct sweep width in **ased**, which can then be set appropriately as **s swh** before transform.

22.10 Examples

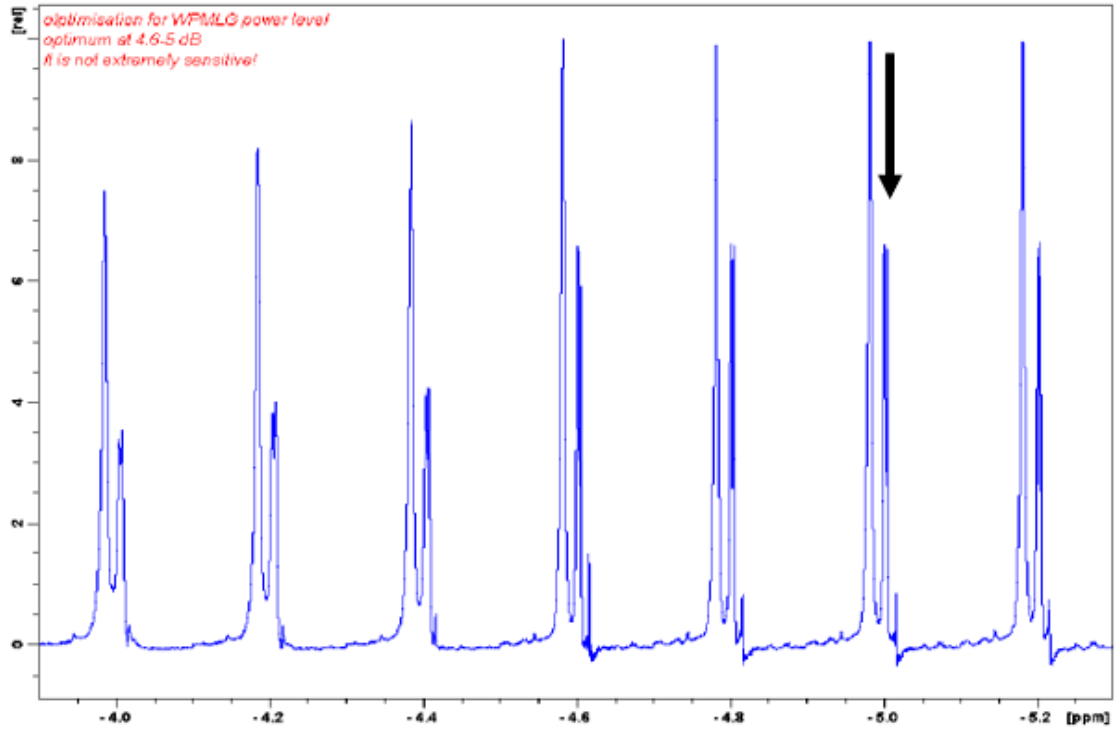


Figure 22.6: Optimizing sp1 for Best Resolution

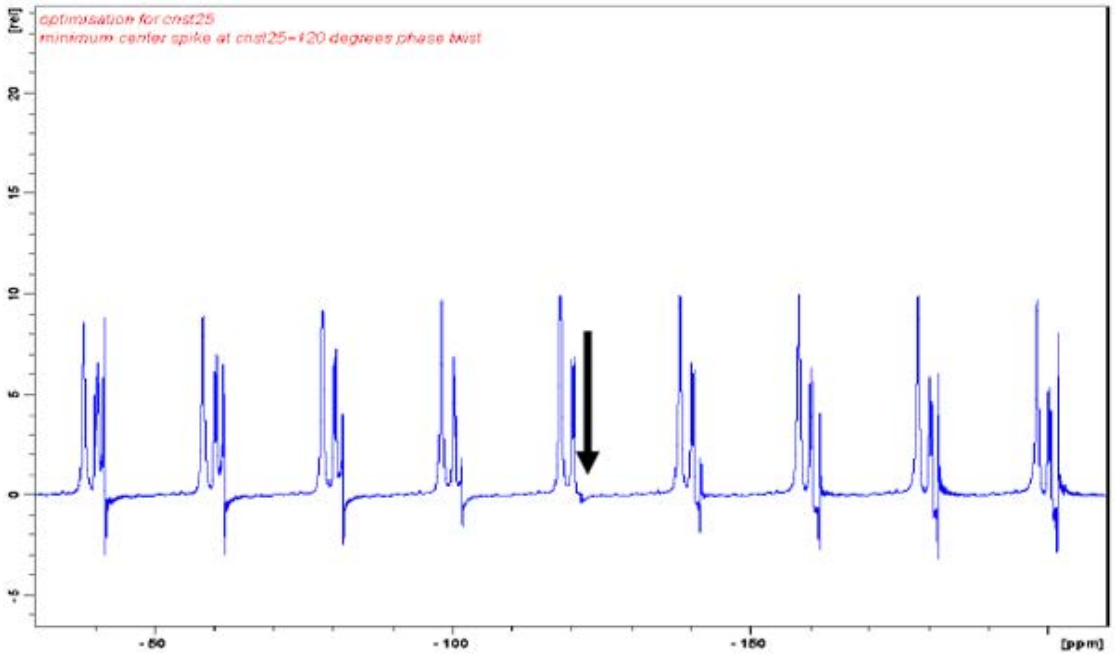


Figure 22.7: Optimizing cnst25 for Minimum Carrier Spike, Optimized at 120°C

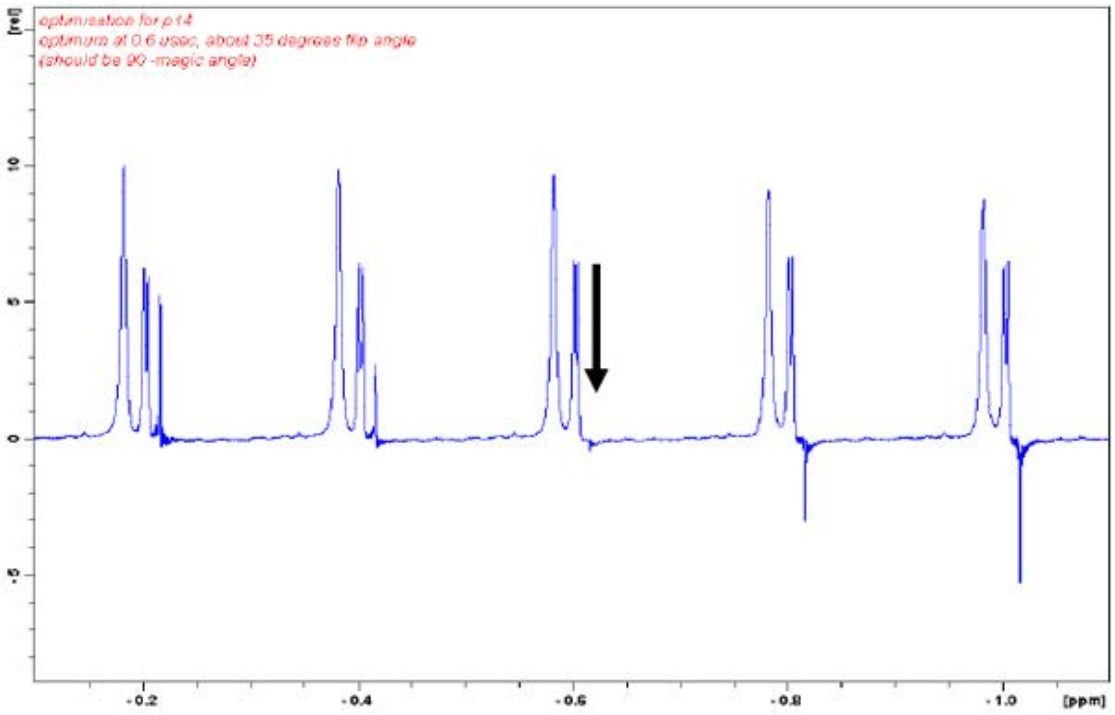


Figure 22.8: Optimizing p14 for Minimum Carrier Spike, Optimized at 0.6 µsec

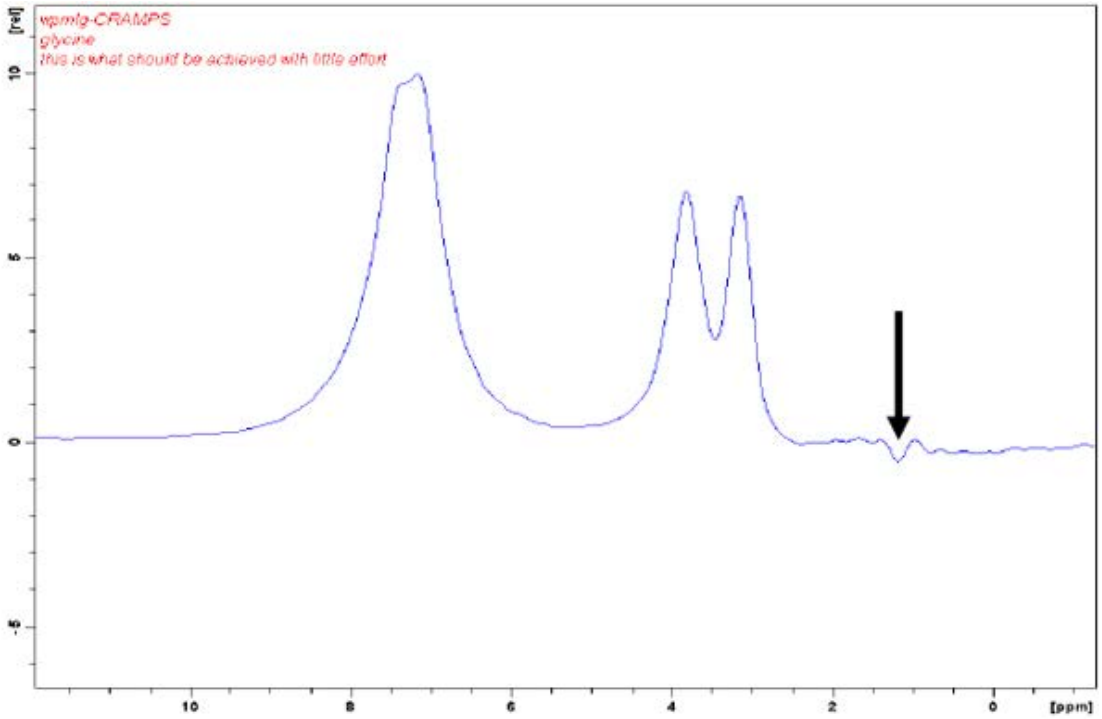


Figure 22.9: WPMLG-CRAMPS After Optimization, Digital Acquisition

23 Modified W-PMLG

Recently, a modified version of WPMLG was published by Leskes et al., which suppresses the carrier spike completely and therefore allows placing the carrier frequency **o1**, **o1p** arbitrarily. This is achieved by a 180 degree phase alternation between consecutive WPMLG-pulses. The magic angle tilt pulse is then not required anymore. This reduces setup time and enhances experimental possibilities significantly.

Reference:

M. Leskes, P.K. Madhu and S. Vega, *A broad-banded z-rotation windowed phase modulated Lee–Goldburg pulse sequence for 1H spectroscopy in solid-state NMR*, Chem. Phys. Lett. 447, 370 (2007).

23.1 Pulse Sequence Diagram for Modified W-PMLG

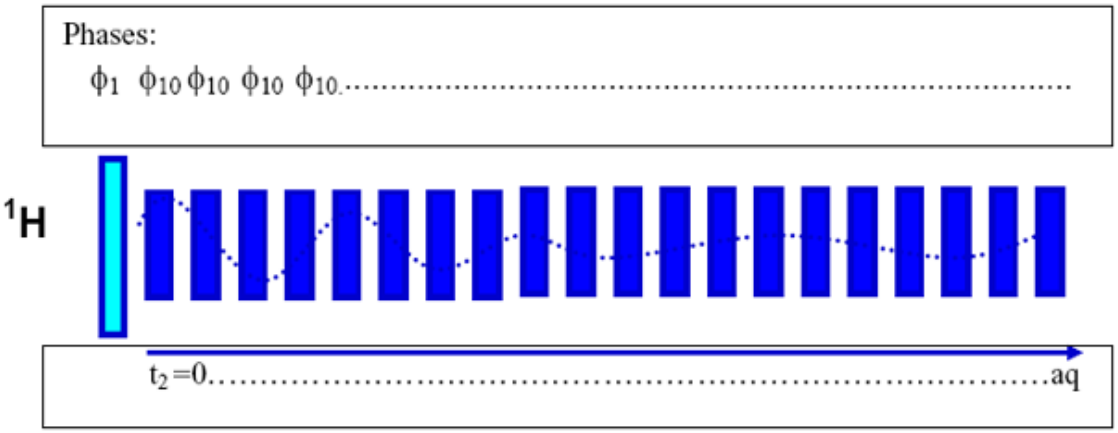


Figure 23.1: Pulse Sequence Diagram

Phases	RF Power Levels	Timing
φ1 = CYCLOPS, 1 2 3 0	pl12=set for around 100 kHz	p1 around 2.5 µsec
φ10 = 0 2	sp1: set for 100-130 kHz=cnst20	WPMLG: calculated from cnst20=RF field
φ31 = CYCLOPS, 0 1 2 3		

Table 23.1: Phrases, RF-Levels, Timings

23.2 Pulse Shapes for W-PMLG

PMLG-shapes: m3p, m3m, m5m, m5p

All these shapes perform similarly. M3p and m3m use 6 phases, m5p and m5m use 10 phases to generate the phase ramp. Obviously, 6 phases generate a phase ramp with less resolution, but shorter possible duration. With the timing resolution available on AV instruments, there is no need to prefer the coarse phase ramp. The letters m and p refer to the sense of phase rotation which is opposite between m and p. If probe tuning is not perfect, m or p may give different results depending on the carrier position. The overall added phases of 0 and 180 degrees on consecutive shape pulses are set by the phase program (phase list ph10).

##TITLE= m5p ##JCAMP-DX= 5.00 Bruker JCAMP library ##DATA TYPE= Shape Data ##ORIGIN= Bruker BioSpin GmbH ##OWNER= <hf> ##DATE= 2005/11/29 ##TIME= 14:47:39 ##$SHAPE_PARAMETERS= ##MINX= 1.000000E02 ##MAXX= 1.000000E02 ##MINY= 1.125000E01 ##MAXY= 3.487500E02 ##$SHAPE_EXMODE= None ##$SHAPE_TOTROT= 0.000000E00 ##$SHAPE_TYPE= Excitation ##$SHAPE_USER_DEF= ##$SHAPE_REPHFAC= ##$SHAPE_BWFAC= 0.000000E00 ##$SHAPE_BWFAC50= ##$SHAPE_INTEGFAC= 6.534954E-17 ##$SHAPE_MODE= 0 ##NPOINTS= 10 ##XYPOINTS= (XY...XY) 1.000000E02, 20.78 1.000000E02, 62.35 1.000000E02, 103.92 1.000000E02, 145.49 1.000000E02, 187.06 1.000000E02, 7.06 1.000000E02, 325.49 1.000000E02, 283.92 1.000000E02, 242.35 1.000000E02, 200.78 ##END	##TITLE= m3p ##JCAMP-DX= 5.00 Bruker JCAMP library ##DATA TYPE= Shape Data ##ORIGIN= Bruker BioSpin GmbH ##OWNER= <hf> ##DATE= 2005/11/29 ##TIME= 14:47:39 ##$SHAPE_PARAMETERS= ##MINX= 1.000000E02 ##MAXX= 1.000000E02 ##MINY= 1.125000E01 ##MAXY= 3.487500E02 ##$SHAPE_EXMODE= None ##$SHAPE_TOTROT= 0.000000E00 ##$SHAPE_TYPE= Excitation ##$SHAPE_USER_DEF= ##$SHAPE_REPHFAC= ##$SHAPE_BWFAC= 0.000000E00 ##$SHAPE_BWFAC50= ##$SHAPE_INTEGFAC= 6.534954E-17 ##$SHAPE_MODE= 0 ##NPOINTS= 6 ##XYPOINTS= (XY...XY) 1.000000E02, 214.64 1.000000E02, 283.92 1.000000E02, 353.21 1.000000E02, 173.2 1.000000E02, 103.92 1.000000E02, 34.64 ##END

23.3 Setup

Fine tuning is done in the same way as with the original sequence, except that the carrier is placed on a convenient position within the spectrum. There is no need to minimise the carrier spike, it should be all gone. Somewhat higher power is required for the wpmlg-shapes.

23.4 Parameter Settings for PMLG and DUMBO

Parameter	Value	Comment
pulprog	wpmlga2	Runs on AV 3 instruments only.
pl12 sp1	for 100 kHz RF field dto, set cnst20=100 000	 To be optimized during setup.
spnam1	m5m or m5p	
p1	2.5 µsec	As for 100 kHz RF field.
p8	1.2 µsec	
p9	4 – 2.6 µsec	To be optimized.
p5 not used	calculated from cnst20	To be optimized.
d1	4s	For α-glycine.
l11=anavpt	4	2, 4, 8, 16 or 32.
o1p	3 - 8	To be optimized.
swh	1e6/2*(2*p9+10*p5)*0.47	To be corrected for proper scaling.
rg	16-64	
td	512	Up to 1024.
si	4k	
digmod	analog	
MASR	12-15 kHz	Depending on cycle time.

Table 23.2: PMLG, Analog Mode

Parameter	Value	Comment
pulprog	dumboa2	Runs on AV 3 instruments only.
pl12 sp1	For 100 kHz RF field up to 130 kHz	 To be optimized during setup.
spnam1	dumbo1_64	Set by xau dumbo.
p1	2.5 µsec	For 100 kHz RF field.
p8	1.2µsec	
p9	4 – 2.6 µsec	To be optimized.
p10	32 µsec or 24 µsec	Set by xau dumbo.
d1	4s	For α-glycine.
l11=anavpt	4	2, 4, 8, 16 or 32.

o1p	5	To be optimized.
swh	1e6/2*(2*p9+p10)*0.5	To be corrected for proper scaling.
rg	16-64	
td	512	Up to 1024.
si	4k	
digmod	analog	
MASR	10-12 kHz	Depending on cycle time.

Table 23.3: DUMBO, Analog Mode

23.5 Fine Tuning for Best Resolution

Fine tuning is done by optimizing power levels, pulse widths and carrier offset as before, the carrier spike is gone, spikes at both sides may appear.

23.6 Correcting for Actual Spectral Width

The modified sequence has a slightly different scaling factor of 0.47.

23.7 Digital Mode Acquisition

Most parameters stay the same as adjusted in analogue mode.

Parameter	Value	Comment
pulprog	dumbod2 or wpmlgd2	AV 3 instruments only.
digmod dspfirm	digital sharp or medium	
aqmod	qsim or dqd	
swh	50000-10000	Depending on spectral range and o1.

Table 23.4: Parameters for Digital Mode

The correction for the scaling factor must be done after acquisition, changing the status parameter **swh** by typing **s swh** and dividing the value by the scaling factor (about 0.47 for WPMLG and 0.5 for DUMBO).

24 CRAMPS 2D

CRAMPS methods allow measurement of chemical shifts in the presence of strong homonuclear dipolar interactions. Therefore, CRAMPS-type sequences can be applied to measure chemical shifts of protons (where these sequences work most efficiently, and where fast spinning cannot easily be used). As an example, the proton-X heteronuclear chemical shift correlation experiment (see *Decoupling Techniques [▶ 79]*) uses FSLG to suppress homonuclear dipolar couplings between protons to resolve the proton chemical shifts. CRAMPS-type pulse sequences must be used in both dimensions if proton chemical shifts are to be correlated.

Two types of proton-proton correlation experiment will be described here:

1. Proton-proton shift correlation via spin diffusion (similar to the high resolution NOESY-experiment). In this case, the dipolar coupling between protons acts during the mixing period. The size of the off-diagonal cross peaks indicates the size of the dipolar coupling between the correlated sites.
2. Proton-proton DQ-SQ correlation (similar to the high resolution INADEQUATE) correlates proton chemical shifts with DQ-frequencies of dipolar coupled sites.

The modifications according to the *Modified W-PMLG [▶ 257]* are implemented in order to remove the carrier spike. Without a carrier spike, 2D experiments are much easier and faster to set up. Being able to set the carrier close to the desired spectral range, one can make the total acquired window smaller as well along F2 (using digital mode) as along F1.

References

P. Caravetti, P. Neuenschwander, R.R. Ernst, *Macromolecules*, 18, 119 (1985).

S.P. Brown, A. Lesage, B. Elena and L. Emsley, *Probing Proton-Proton Proximities in the Solid State: High-Resolution Two-Dimensional 1H-1H Double-Quantum CRAMPS NMR Spectroscopy*, J. Am. Chem. Soc. 126, 13230 (2004).

24.1 Proton-Proton Shift Correlation (spin diffusion)

The standard CRAMPS setup must be executed first (see chapters *CRAMPS 1D [▶ 247]* and *Modified W-PMLG [▶ 257]*). Any homonuclear dipolar decoupling scheme may be used, but in the following the experiment is described using windowed pmlg (w-PMLG). The reasons are the following:

1. At fast spin rates over 10 kHz, only w-PMLG and DUMBO work well. The sequence can easily be modified to use DUMBO, replacing the w-PMLG shapes by DUMBO-shapes and modifying the shape timing accordingly.
2. W-PMLG is easy to set up, since it is rather insensitive to power level missets and frequency offsets. When the experiment setup for the 1D experiment has been executed, no further setup is required for the 2D experiment. Start from the 1D experiment on your sample (the recommended setup sample is glycine) and generate a 2D data set by clicking on the symbol 1, 2 in the headline of the acquisition parameters.

24.2 Pulse Sequence Diagram

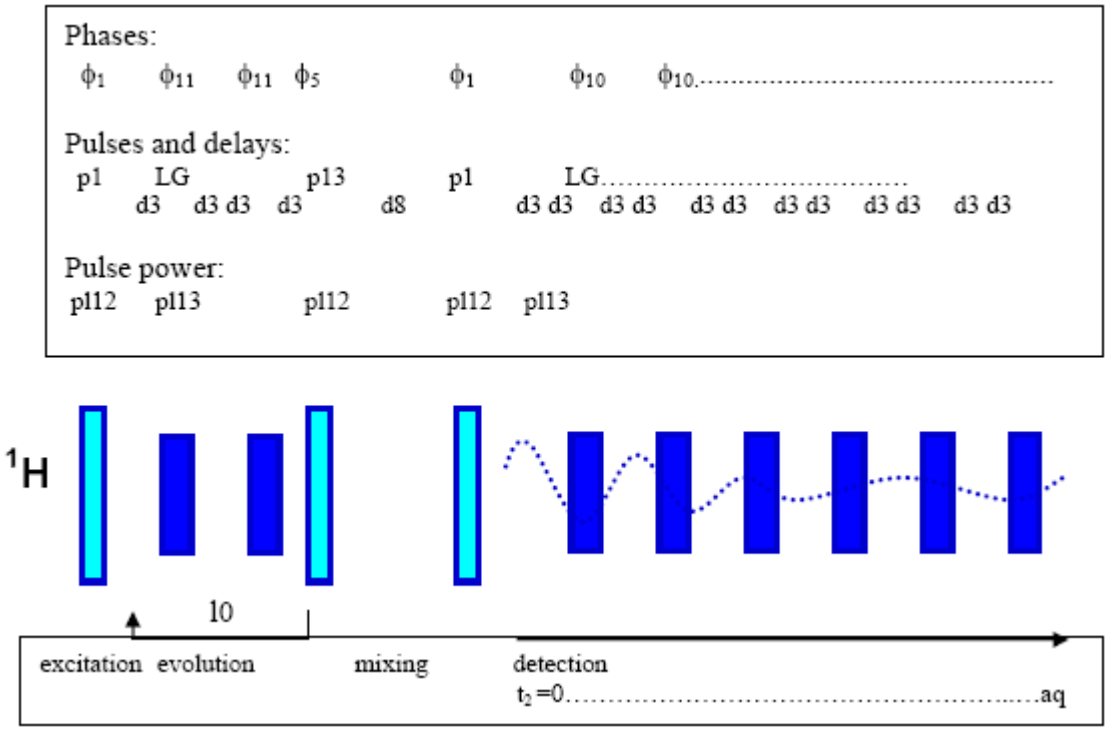


Figure 24.1: Pulse Sequence Diagram

This sequence is written in such a way that the windowed PMLG-unit is used both for detection and for the shift evolution along F1. This was done to minimize the setup requirements. In principle, a windowless sequence can be used as well and should give better resolution along F1. The power level for a windowless sequence is however usually slightly different from the windowed sequence, so this needs to be adjusted separately. Likewise, decoupling during t1 could be implemented using real frequency shifts as in the HETCOR sequence (see *chapter 5 [▷ 79]*). If a windowless sequence is incorporated, the windows d3 must of course be removed. A simple windowless FSLG-unit can be used, with a shape like lgs-2 or lgs-4 having duration of twice or 4 times the length of the w-pmlg pulse.

Parameter	Value	Comment
pulprog	wpmlg2d.	AV 3 instruments only, topspin 2.1 or later.
FnMODE	STATES-TPPI.	Any other method may be used with appropriate changes in ppg.
NUC1, NUC2	^{1}H.	
sw, swh along F1	Same as for F2.	Needs to be corrected before transform pulse program calculates approximate values to be set before transform (ased).
td	512-1k.	Depending on resolution.
1 td	128-256.	Depending on resolution.
spnam1	wpmlg1, m5m or m5p as in 1d.	DUMBO may be used with modified timing.
spnam2	lgs-2 or lgs-4 if used.	Set l3=2 or 4, depending on desired sw1 DUMBER-22 with modified timing.

Table 24.1: Acquisition Parameters

Phases	RF Power Levels	Timing
φ0 = 0, STATES-TPPI	pl12 = set for around 100 kHz	p1 around 2.5 µsec.
φ1 = CYCLOPS, 1 2 3 0	pl12	p1
φ5 = 2		
φ10 = 0 2	sp1,sp2: set for 100-130 kHz RF-field or pl13 for both set in ppg	WPMLG: calculated via cnst20. DUMBO: p10 set by xau dumbo.
φ11 = 0 2	dto.	dto.
φ31 = CYCLOPS, 0 1 2 3		d8 = desired mixing time, 50-1000 µs.

Table 24.2: Phases, RF-levels, and Timings

24.3 Data Processing

The spectral width in both dimensions assumes the absence of shift scaling. In order to account for the shift scaling effect of the sequence, one has to increase the spectral width by the scaling factor. Before doing the 2D-fourier transformation, type **s sw** to call the status parameters for both F2 and F1 and replace both values by <current value>/0.6. After **xfb**, the relative peak positions will be (approximately) correct, but the absolute peak positions must be corrected by calibrating a known peak position to the correct value. The pulse program is written such that the correctly scaled sweep widths are calculated and indicated upon ased. These values are set as status parameters before transform as indicated above.

Parameter	Value	Comment
mc2	STATES-TPPI	
wdw	QSINE	Slight-moderate resolution enhancement is usually required.
ssb	3 or 5	
si	2k -4k	
1 si	512 – 1k	

Table 24.3: Processing Parameters

24.4 Examples

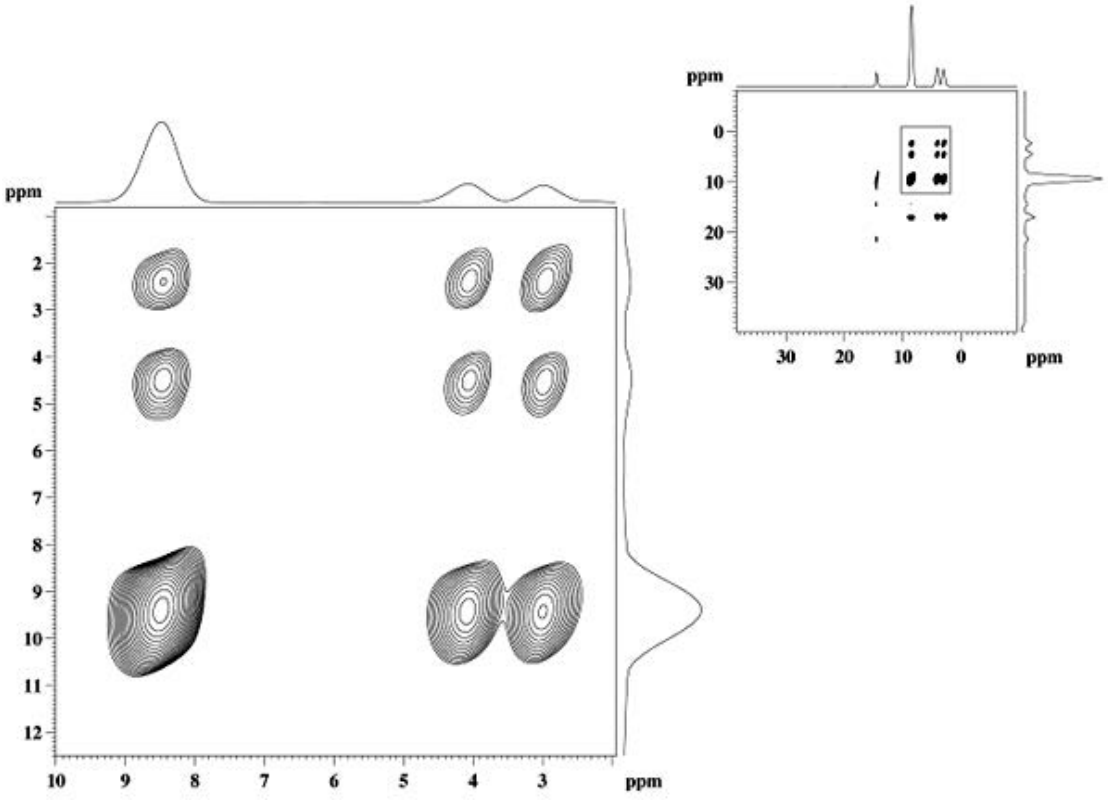


Figure 24.2: Setup and Test Spectrum of Alpha-glycine

The figure above shows the setup and test spectrum of alpha-glycine (note that glycine samples containing gamma glycine will show additional peaks!). The protons attached to the alpha-carbon are in-equivalent and strongly coupled. The cross peaks at 3 and 4 ppm will show at a mixing time as short as 50 µsec, the cross peaks to the NH_3-protons at 9 ppm require 200 -300 µsec to show. The mixing time here was 500 µsec. A sequence without carrier spike suppression was used here.

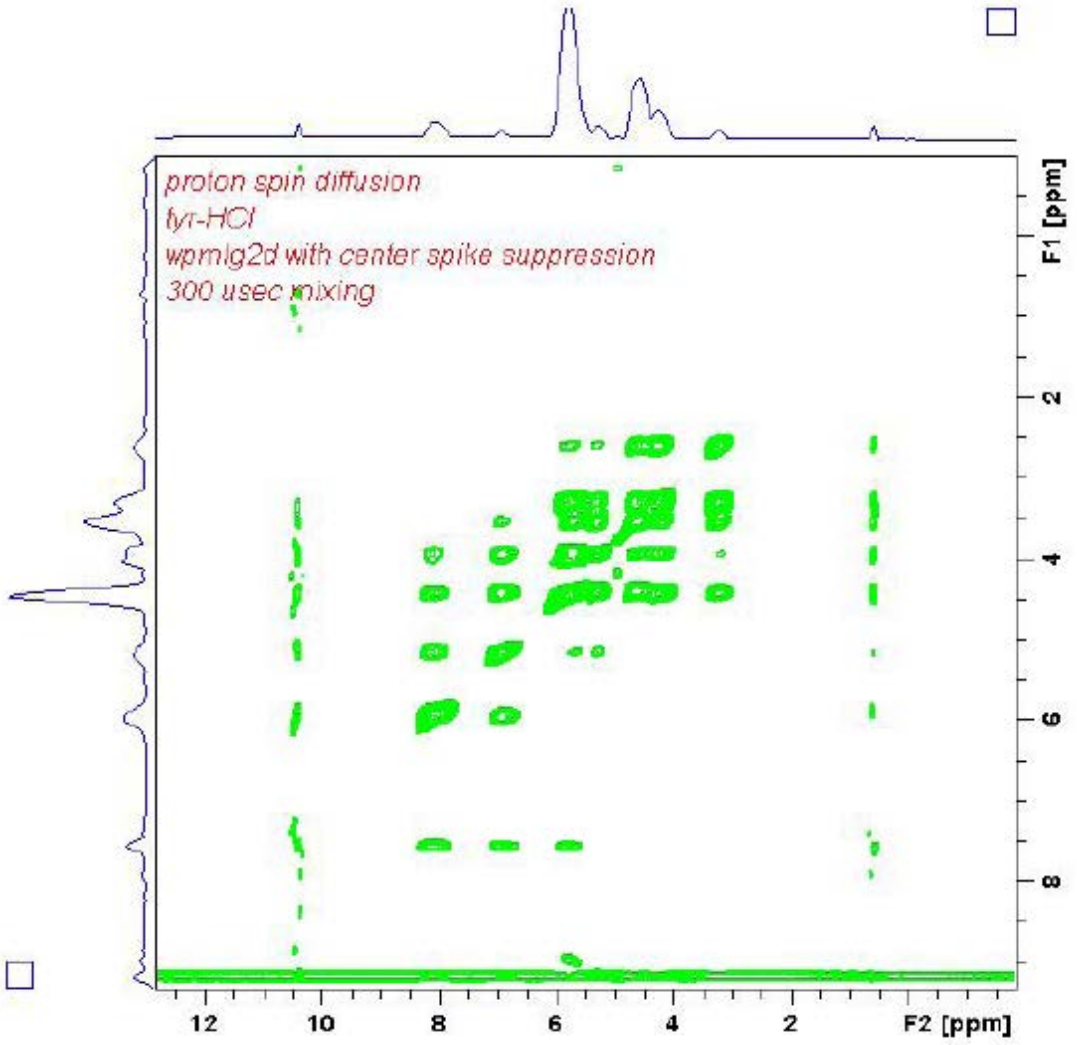


Figure 24.3: Spectrum of Tyrosine-hydrochloride

The mixing time was 300 µsec to show all connectivities. Full plot to show that smaller sweep widths can be chosen when the carrier can be conveniently placed within the spectrum.

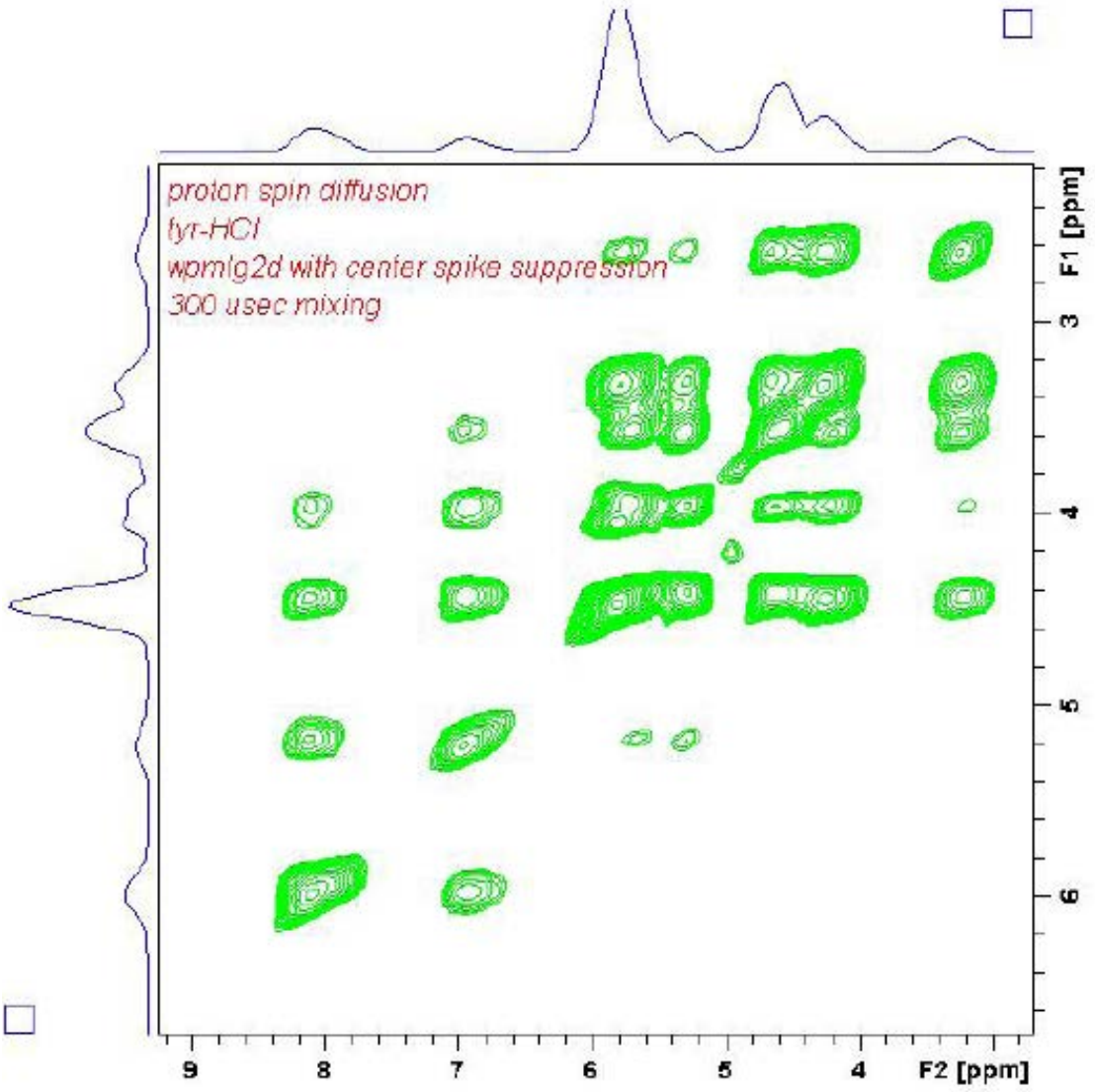


Figure 24.4: Expansion of the Essential Part of the Spectrum

24.5 Proton-Proton DQ-SQ Correlation

This experiment correlates proton shifts (F2) with double quantum frequencies (sum of shifts of the correlated sites). Double quantum transitions are excited and reconverted by a post-C7 or similar sequence.

24.6 Pulse Sequence Diagram

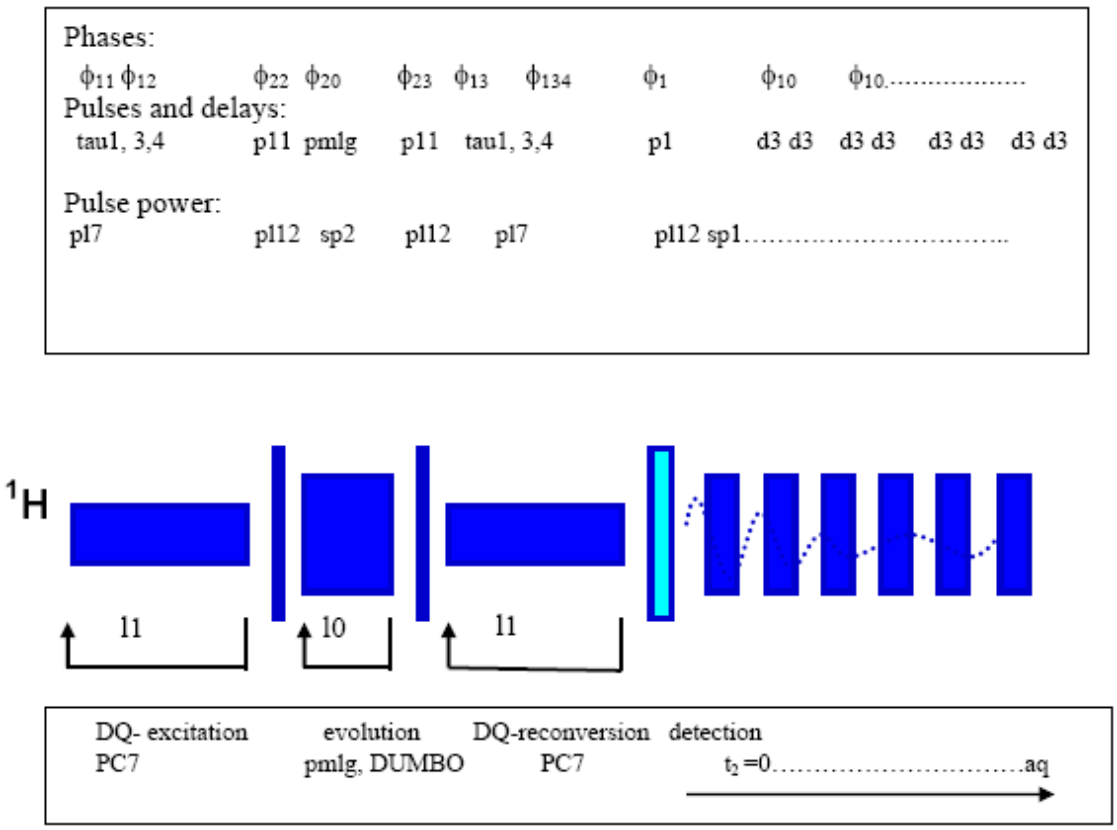


Figure 24.5: Pulse Sequence Diagram

When applied to X-nuclei like ^{13}C, the RF field during this sequence must be carefully matched to the 7-fold spin rate, since the dipolar couplings are small, and care must be taken that the excitation bandwidth of the sequence chosen covers the whole shift range of the X-nucleus. In the case of protons, this is rather forgiving, since the shift range to be covered is small, and the required power levels are easily achieved for protons. Usually it is enough to calculate the required power level from the spin rate and the known proton 90 degree pulse using the au program **calcpowlev**. Assume the spin rate is 14000 Hz and post-C7 is used. The required RF field is then 7*14000=98000 Hz. The known proton 90 degree pulse is 2.5 µsec=1/4*2.5e-6 =100000 Hz. Type **calcpowlev** and enter 100000, return, then enter 98000, return. The output will be “change power level by 0.18 dB”. The power level for the p-C7 sequence is therefore 0.18 dB to higher attenuation than what is required for a 2.5 µsec pulse.

Parameter	Value	Comment
pulprog	wpmlgdqsq	AV 3 instruments only, topspin 2.1 or later.
FnMODE	STATES-TPPI	Any other method may be used with appropriate changes in ppg.
NUC1, NUC2	^{1}H	
sw, swh along F1	same as for F2	Needs to be corrected before transform pulse program calculates approximate values upon **ased**.
td	512-1k	Depending on resolution.
1 td	128-256	Depending on resolution.
cnst31	spin rate, 10-15 000	Depending on available RF field.
l1	number of pc7-cycles	2-7 depending on dipolar coupling.
spnam1	m5m or m5p as in 1d setup	DUMBO may be used with modified timing.
spnam2	lgs-2 or lgs-4 if used.	Set l3=2 or 4, depending on desired sw1 DUMBER-22 with modified timing.

Table 24.4: Acquisition Parameters

Phases	Rf Power Levels	Timing
φ11,12 = POST-C7=φ13,14 φ11,12 incremented for DQ-evol. φ11,12 incremented for DQ-select	pl7 set for RF=7*spin rate	tau1,3,4 calculated from cnst31.
φ1 = CYCLOPS	pl12	p1
φ22 = 3	pl12	p11, ~45°
φ23 = 1	pl12	p11
φ10 = 0	sp1: set for 100-130 kHz RF-field	WPMLG: calculated via cnst20 DUMBO: p10 set by xau dumbo.
φ31 = DQ selection		

Table 24.5: Phases, RF-Levels and Timing

24.7 Data Processing

The spectral width in both dimensions assumes the absence of shift scaling. In order to account for the shift scaling effect of the sequence, one has to increase the spectral width by the scaling factor. Before doing the 2D-fourier transformation, type **s sw** to call the status parameters for both F2 and F1 and replace both values by <current value>/0.6. After **xfb**, the relative peak positions will be (approximately) correct, but the absolute peak positions must be corrected by calibrating a known peak position to the correct value.

Parameter	Value	Comment
mc2	STATES-TPPI	Or whatever used.
wdw	QSINE	Slight-moderate resolution enhancement is usually required.
ssb	3 or 5	
si	2k -4k	
1 si	512 – 1k	

Table 24.6: Processing Parameters

24.8 Examples

These spectra were both taken without the modification according to *Modified W-PMLG [▷ 257]*, so the offset is placed to the down field side and the spectrum width was chosen larger than necessary. The small plots show the full spectrum.

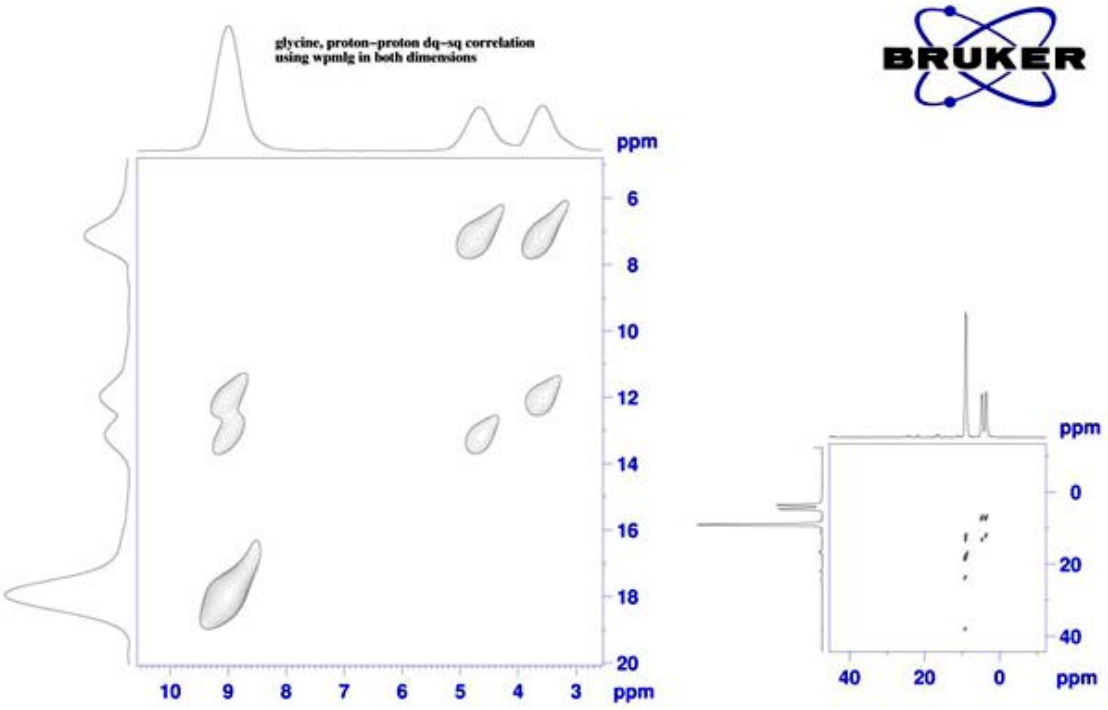


Figure 24.6: Glycine, Proton-Proton DQ-SQ Correlation Using WPMLG in Both Directions

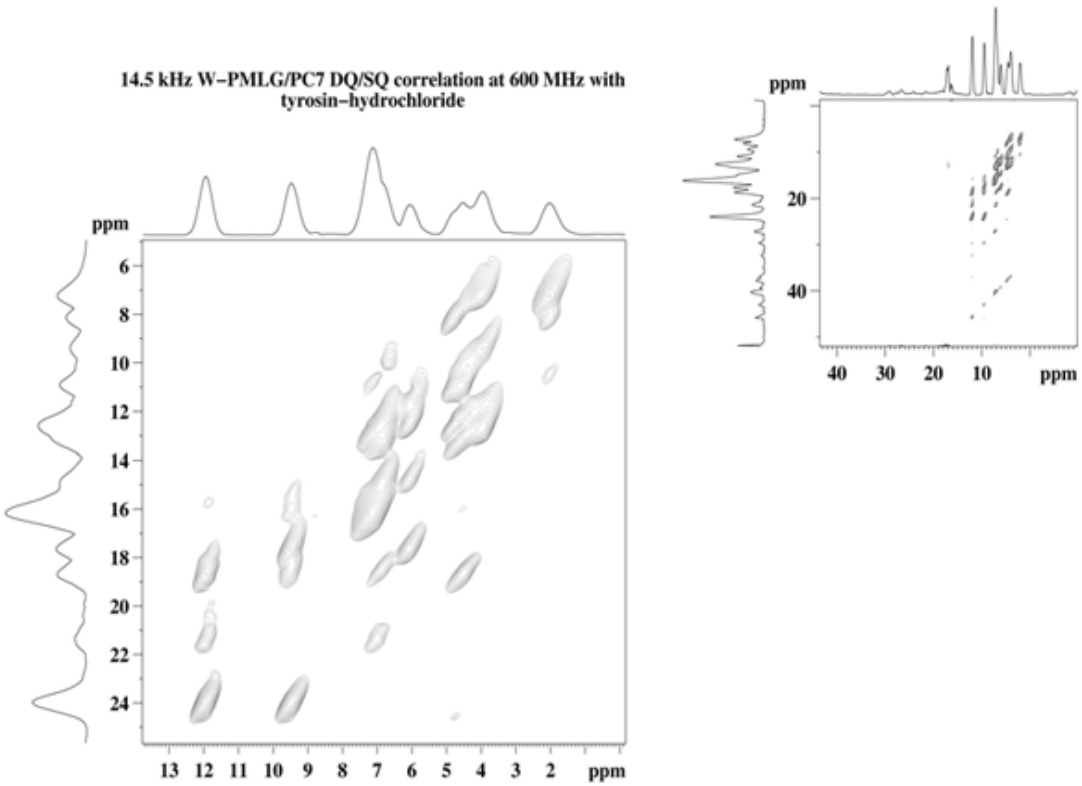


Figure 24.7: 14.5 kHz W-PMLG/PC7 DQ/SQ Correlation at 600 MHz with Tyrosine-Hydrochloride

25 Editing Techniques Using J-couplings

25.1 MAS-J-HMQC

In this chapter the MAS-J-HMQC experiment is described where the scalar J-coupling through the chemical bond is used for the transfer. In this experiment correlations are visible only for directly bonded pairs of unlike spins (mainly ^{1}H-^{13}C, but ^{31}P-^{27}Al has also been used) and can unambiguously be differentiated from short distance through space correlations which can be visible in the FSLG-HETCOR experiment (see *FSLG-HETCOR [▶ 105]*).

References

L. Emsley, D. Sakellariou, A. Lesage, S. Steuernagel, *Verfahren zum Verbessern der Auflösung in zweidimensionalen heteronuklearen Korrelationsspektren der Festkörper-NMR*, German Patent DE19834145C1, March 9, 2000.

A. Lesage, D. Sakellariou, S. Steuernagel, and L. Emsley, *Carbon-Proton Chemical Shift Correlation in Solid-State NMR by Through-Bond Multiple-Quantum Spectroscopy*, J. Am. Chem. Soc. 120, 13194 (1998).

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C. Coelho, T. Azais, L. Bonhomme-Coury, J. Maquet, D. Massiot and C. Bonhomme, *Application of the MAS-J-HMQC experiment to a new pair of nuclei {29Si,31P}: Si O(PO) and SiP2O7 polymorphs,* J. Magn. Reson. 179(1), 114-119 (2006).

25.1.1 Pulse Sequence Diagram for MAS-J-HMQC

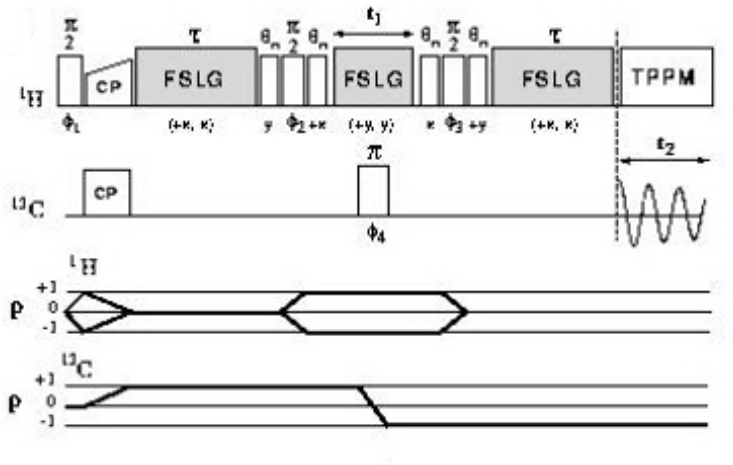

Figure 25.1: The Pulse Sequence and Coherence Transfer Pathway of the MAS-J-HMQC Experiment

Φ1 = 1 3	Φ2 = 1 1 3 3 (+States-TPPI)	Φ3 = 1
Φ4 = {0}*4 {1}*4 {2}*4 {3}*4	Φrec. = 0 2 2 0 2 0 0 2	

Note: The proton θ_m pulses around the 90° pulses are required to flip the proton magnetization from the tilted frame to the rotating frame. The 2nd pair of θ_m pulses can be taken out if the FSLG block during t_1 is along +x/-x, as is done in the current pulse program

25.1.2 Setting up MAS-J-HMQC

The MAS-J-HMQC sequence is in full analogy to the liquids 4-pulse HMQC experiment, i.e. excitation and generation of anti-phase magnetization, evolution of heteronuclear multiple quantum coherence, conversion back to anti-phase, refocusing and detection with decoupling. However, there are 3 main differences:

- Instead of protons like in liquids an X nucleus (mainly 13C) is detected.
- The initial 90° pulse (which should now be a 13C pulse) is replaced with a CP step from protons for the well known advantages of CP in the case of low abundant nuclei.
- During all τ delays and during evolution proton homonuclear decoupling is applied. In combination with MAS this means that – apart from chemical shifts – heteronuclear scalar J couplings are the only remaining interactions.

The recommended probe for this experiment is a 3.2 or 4 mm MAS probe which allow spinning frequencies between 10 and 20 (3.2 mm only) kHz. This is a range of spinning frequencies at which FSLG decoupling performs best. And it is fast enough to average residual heteronuclear dipole interactions. The maximum spinning of 7 kHz with a 7 mm probe is not sufficient. 1.3 or 2.5 mm probes are not recommended because FSLG does not benefit from spinning very fast and sample amounts in these probes don't allow sufficient sensitivity. Further details of FSLG are explained in other chapters (e.g. *FSLG-HETCOR [▶ 105]*).

Start from a data set with well adjusted cross polarization and proton decoupling at fairly high RF field. Dependent on the probe and the magnetic field 100 to 125 kHz is recommended. At magnetic fields below 400 MHz 80 to 100 kHz may also be sufficient.

If an FSLG-HETCOR experiment of your sample of interest has already been run you may use this sample for set-up right away. If you start from scratch we will recommend a well-known set-up sample, e.g. ^{13}C natural abundance (!) tyrosine hydrochloride. But any other well crystalline organic compound will do, of course a short proton T_1 is always preferred at this stage.

Spin the sample at any frequency between 10 and 15 kHz and tune the probe.

From the data set with the CP setup create a new dataset. Set pulprog **masjhmqc**. Type ased or click the pulse symbol in **eda**.

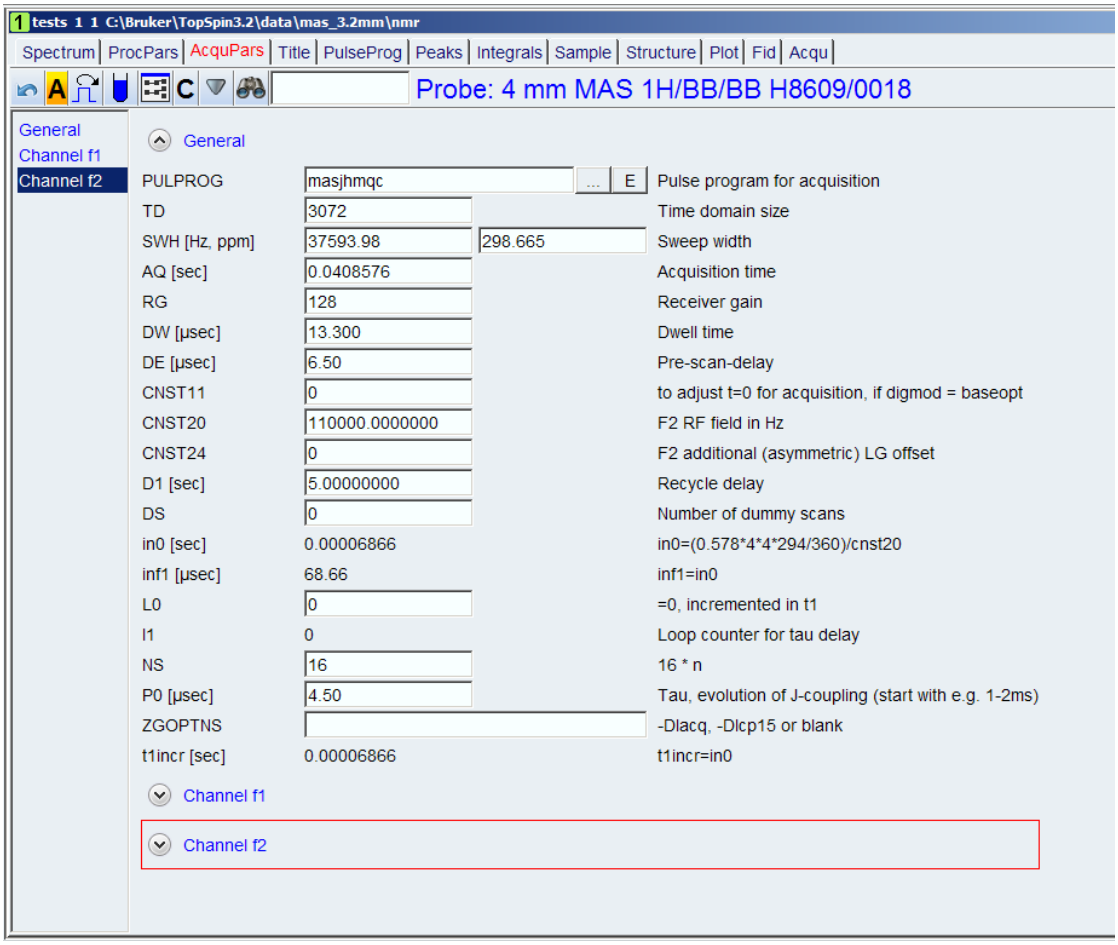

Figure 25.2: Display of ased in the case of 1D dataset

Performing **ased** will show all parameters which are essential for the acquisition, not all available parameters. In addition it performs calculations which are specified in the pulse program. Note that all parameters which are calculated are not editable, and will show only, if explicitly used during the main pulse program between ze and exit. The master parameter to set FSLG condition is **cnst20** which must equal the proton RF field, i.e. $1/\tau_{360°}$. If you have an older version of the masjhmqc pulse program this value can be entered directly in Hz, e.g. as 100000. From this the corresponding offsets (as **cnst22** and **cnst23**) and the duration of the pulse length (**p5**) are calculated within the included file lgcalc.incl. Since FSLG decoupling and MAS interfere with one another they must be synchronized. This can be achieved in such a way that an integer number of FSLG blocks (i.e. pairs of +x/-x 360°-off- resonance pulses) fit into one rotor period.

To avoid tedious calculations on a sheet of paper a line “cnst20=cnst31*1.63*l4” is included before the #include <lgcalc.incl> in a newer version of the pulse program. **cnst31** and **l4** now become the master parameters for the FSLG condition, where **cnst31** is the MAS spinning frequency and **l4** is the number of FSLG blocks per rotor period. **L4** must be set such that **cnst20** will be close to the desired RF field for FSLG. Possible values to solve this are shown in the figure below. It makes more sense to make the spinning frequency the freely adjustable parameter, because this is a sample dependent parameter in order to avoid overlap of spinning side bands and the performance of FSLG is comparable for quite a range of RF fields applied (e.g. between 100 and 120 kHz). To define the overall offset of FSLG decoupling **cnst24** can be used, usually it is set to 0. The actual power for RF is defined by **pl13**, which can be recalculated from the power of the 90° pulse using the **calcpowlev** AU program. If the value of **cnst20** comes out within 10% around the RF field for the 90° pulse **pl13** may as well be set equal to **pl12**.

Note: To recalculate RF fields with **calcpowlev** AU program you can type in RF field values directly (e.g. in kHz) instead of pulse lengths, but now the change in dB has the opposite sign (!).

spinning frequency [Hz] (*cnst31*)	RF field [Hz] (*cnst20*)	number of blocks per rotor period (*l4*)
10000	97800	6
10500	102690	6
11000	107580	6
11500	112470	6
12000	97800	5
12500	101875	5
13000	105950	5
13500	110025	5
14000	114100	5
14500	94540	4
15000	97800	4
15500	101060	4
16000	104320	4
16500	107580	4
17000	110840	4
17500	114100	4
18000	117360	4
18500	90465	3
19000	92910	3
19500	95355	3
20000	97800	3

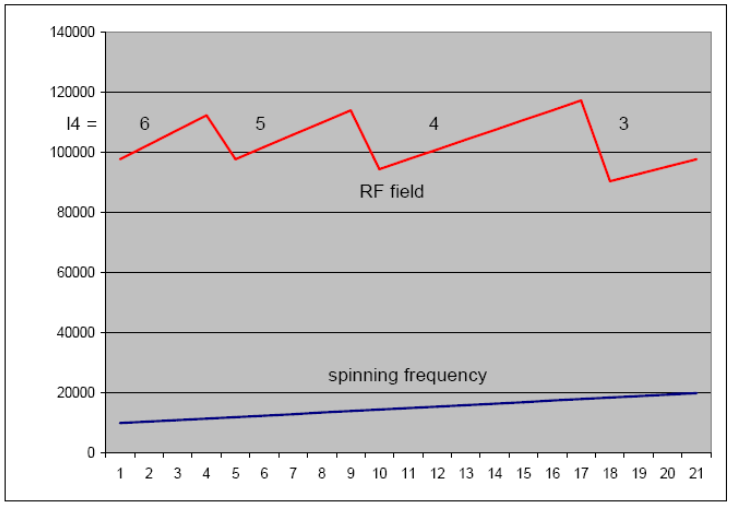

Figure 25.3: Various possible values of MAS spinning frequencies and applicable RF fields

In the figure above various possible values of MAS spinning frequencies and applicable RF fields are shown on the left for different number of rotor periods. Note that the spinning frequency (entered as **cnst31**) can be set freely to any value between 10 and 20 kHz; the RF field is calculated accordingly (taking into account **l4**). The graphics on the right shows that for a given number of blocks per rotor period the FSLG RF field must be increased when a higher the spinning frequency is used.

Before running the 2D experiment the τ values need to be optimized. This is quite essential because the overall sensitivity of the experiment depends on it. To define the duration of τ the loop counter **l1** is used which is calculated from **p0**. In a second step a delay called tau is recalculated and displayed in **ased** to see the precise value. The precise value of tau is always an integer multiple of rotor periods and therefore, **p0** can only be entered close to the desired value. To find the optimum value for tau **popt** can be used, optimizing in such a way that **p0** is incremented in steps of approx. 500 µs for the maximum peak intensities looking at the entire spectrum. Typically between 1.5 and 2.5 ms can be expected for most cases. In the figure below you can see that there won’t be a unique value for τ where all peaks in the spectrum achieve their maximum peak intensities so a compromise must be used. Note that all signals from quaternary carbons must be eliminated by the phase cycle.

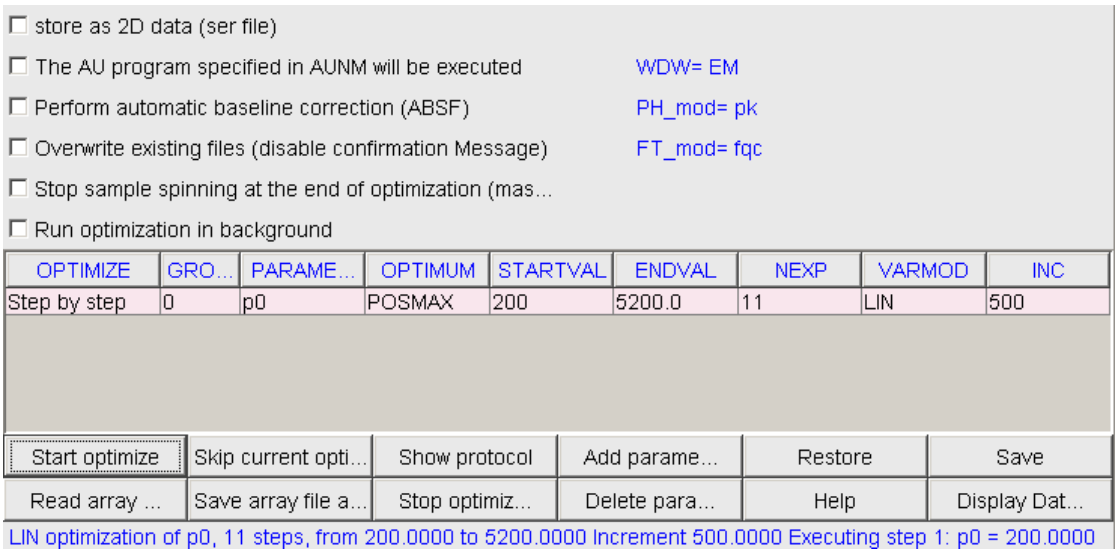

Figure 25.4: Display of popt for the optimization of the τ value

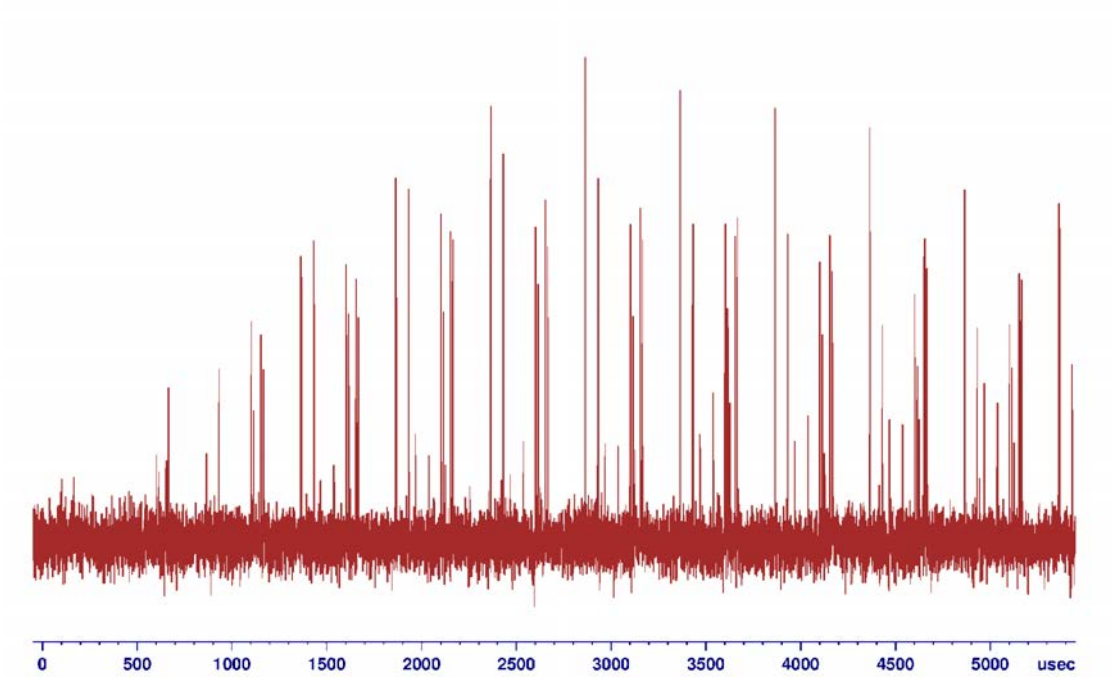

Figure 25.5: Result of popt parameter optimization procedure

Once the optimum value for **p0** is found the set-up is ready to run the 2D experiment. Create a new data set from the 1D set-up and convert to a 2D data set. Parameters should be set according to the figure below. Note that in contrary to the FSLG-HETCOR the t_1 increment is fixed to 4 FSLG blocks. For RF fields above 80 kHz this gives sufficient spectral range in the indirect dimension to avoid folding of peaks.

Note the spectral width in the indirect dimension is automatically calculated from the pulse program taken into account the theoretical scaling factor of 0.578. To cross check the corresponding F1 spectral width in ppm the value of **in_f1** displayed in **ased** must be entered manually into **IN_F**.

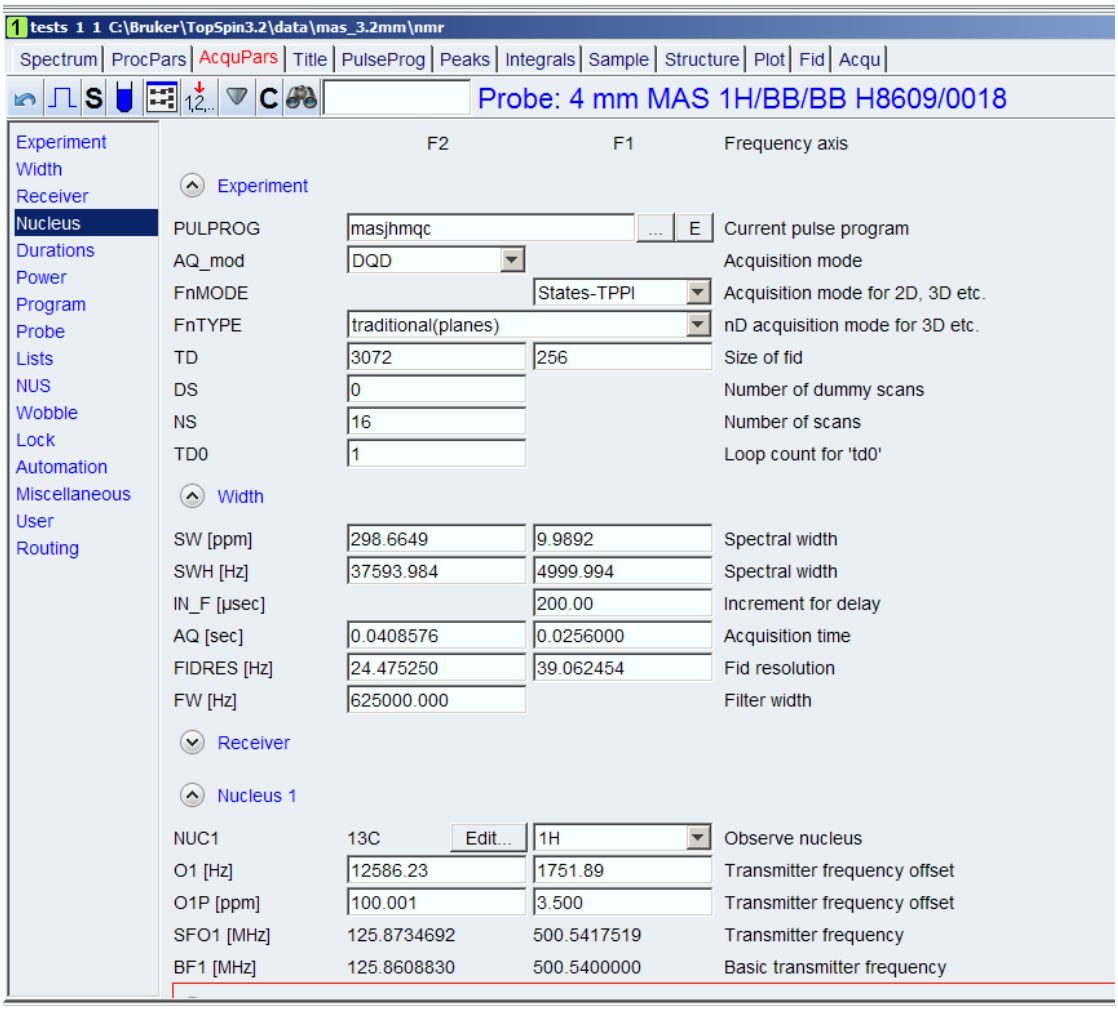

Figure 25.6: First part of display of eda in case of 2D data set

To check the spectral width in F1 the value of **in_f1** calculated in the pulse program and displayed in **ased** must be entered as **IN_F** here.

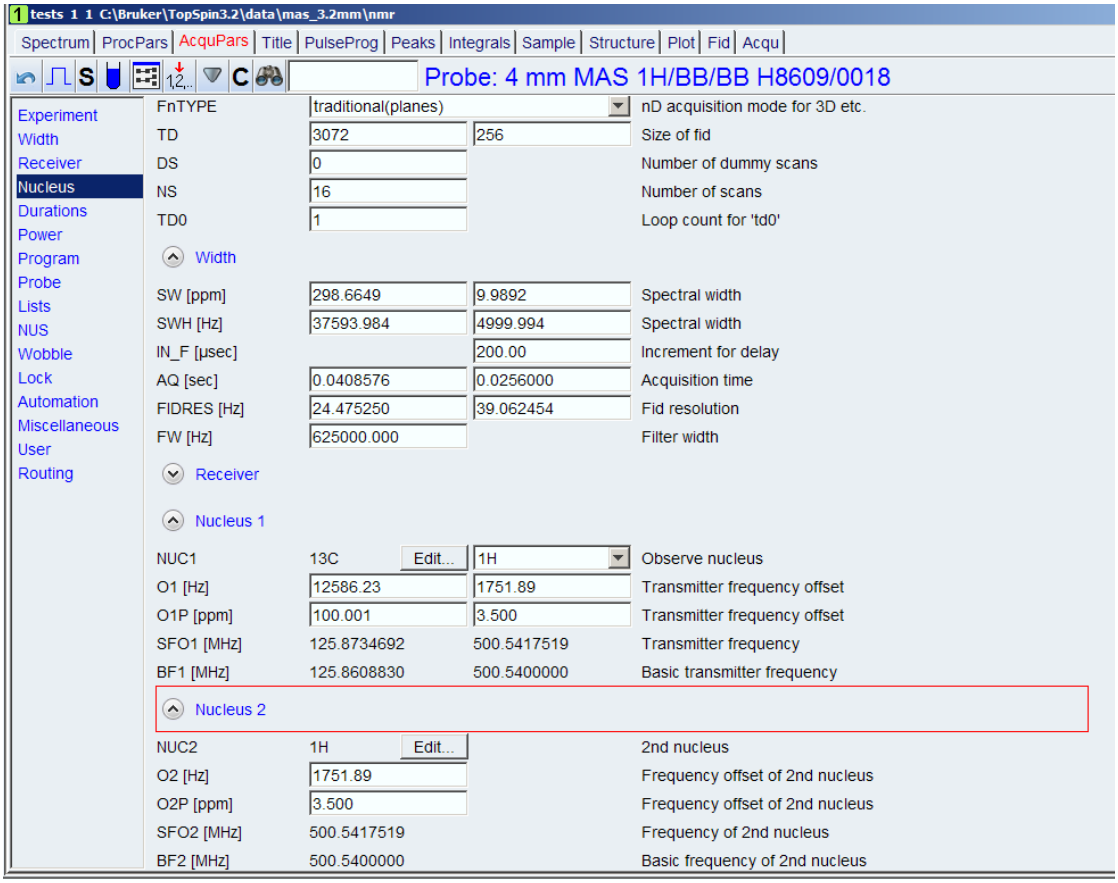

Figure 25.7: Second part of display of eda in case of 2D data set

Cross check that the right offsets are entered for ^{13}C and ^{1}H.

For processing settings should be the same as for FSLG-HETCOR. In the acquisition dimension some line broadening can be applied whereas in the indirect dimension some line narrowing to increase spectral resolution is recommended. The following figure shows the settings of the most relevant parameters.

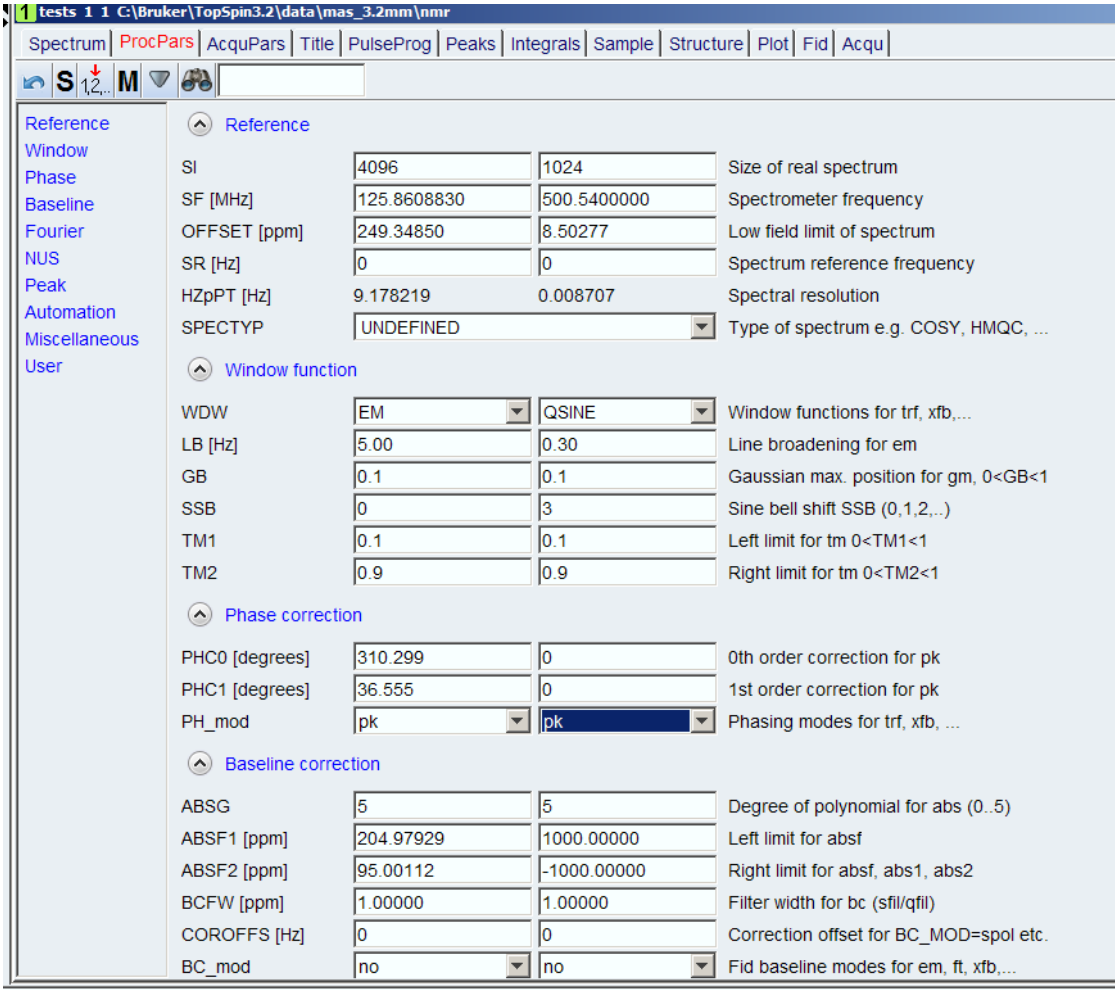

Figure 25.8: Display of edp processing parameters for 2D processing

In the following figure the MAS-J-HMQC of a tripeptide is shown.

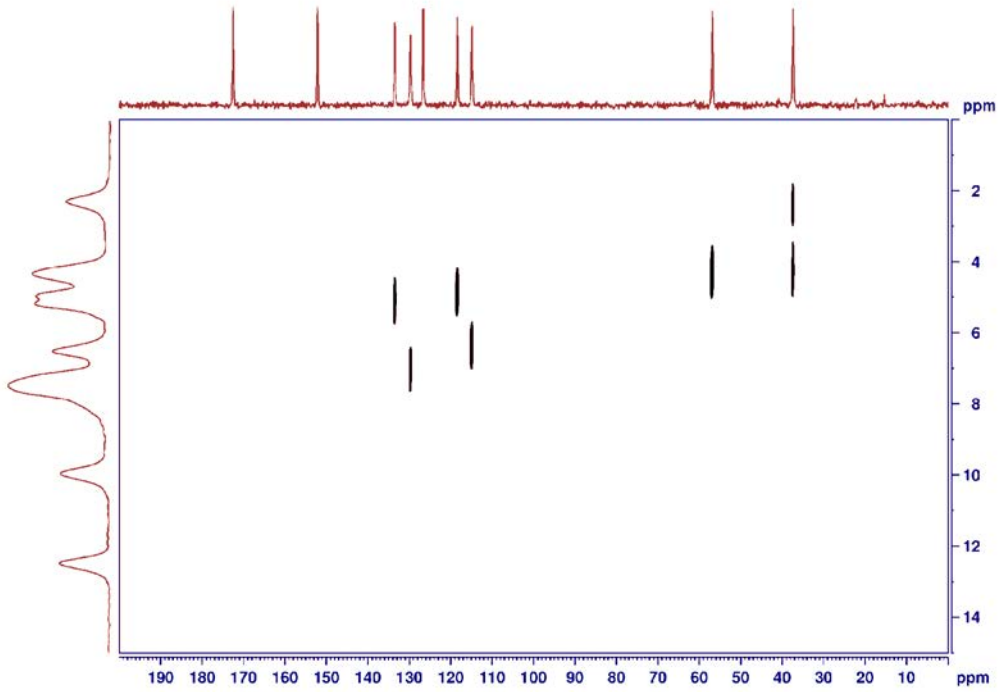

Figure 25.9: Plot of a 2D MAS-J-HMQC of a tyrosine HCl with the obtained and processed with the parameters given above

In the figure above, note that only signals from protonated carbons show correlations in the 2D map. 1D slice on top is regular CP/MAS spectrum, 1D slice on the left is the CRAMPS spectrum.

25.2 Solid State Attached Proton Test (sostapt)

In order to have access to scalar $^{13}C/^{1}H$ couplings in solids the much bigger dipole couplings need to be averaged. For the heteronuclear dipole couplings fast MAS (10 to 20 kHz) is in most cases sufficient to average out, but proton-proton homonuclear couplings must be averaged by manipulation of the protons in spin space. One very efficient and robust way of doing this is FSLG which has already been introduced for the FSLG-HETCOR experiment. The sostapt experiment described here takes advantage of the scalar C/H couplings and allows differentiating differently substituted carbon sites, i.e. CH from CH_2 and CH_3.

Reference

A. Lesage, S. Steuernagel, L. Emsley, *Carbon-13 Spectral Editing in Solid-State NMR Using Heteronuclear Scalar Couplings,* J. Am. Chem. Soc. 120(28), 7095-7100 (1998).

25.2.1 Pulse Sequence Diagram for Solid State Attached Proton Test

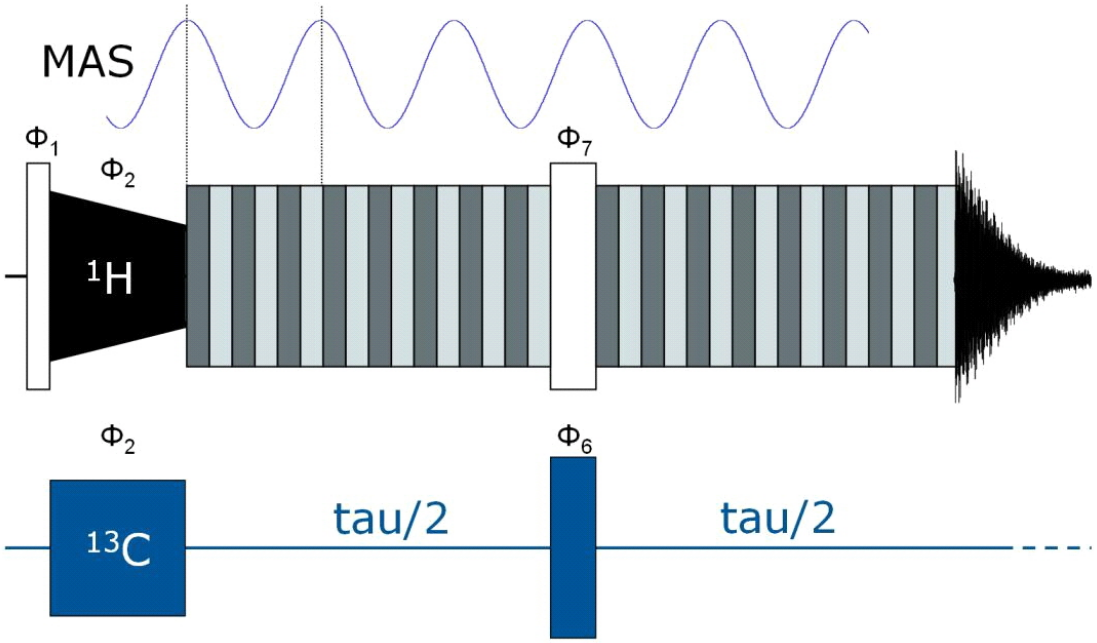

Figure 25.10: The pulse sequence and experiment

Φ1 = 1 3	Φ2 = 0	Φ7 = 1
Φ6 = 1 1 2 2 3 3 0 0	Φrec. = 0 2 2 0	

Note: Together with the sequence the rotation periods of the MAS rotor are shown. It indicates that an integer number of FSLG blocks [typically 4 to 6 ($+\Omega_{+x}$ / $-\Omega_{-x}$) blocks] should fit into a single rotor period.

25.2.2 Setting Up SOSTAPT

The solid state attached proton test experiment is analogous to the liquid state experiment. During to equally long τ delays ^{13}C-^{1}H J-coupling is allowed to evolve. The delays are separated by 180° pulses on each, ^{13}C and ^{1}H in order to refocus chemical shift evolution and to obtain ^{13}C signal intensities modulated with the scalar J-coupling. Again, like in MAS-J-HMQC residual heteronuclear dipole couplings are averaged by MAS; homonuclear proton-proton dipole couplings are averaged by FSLG during the τ delays. Since ^{13}C-^{13}C J-couplings cannot be averaged in this experiment and will disturb the modulation of the signal intensities a sample at natural abundance must be used. Site specific labelling without direct ^{13}C-^{13}C connectivities may work.

The recommended probe for this experiment is a 3.2 or 4 mm MAS probe which allow spinning frequencies between 10 and 20 (3.2 mm only) kHz. This is a range of spinning frequencies at which FSLG decoupling performs best. And it is fast enough to average residual heteronuclear dipole interactions. The maximum spinning of 7 kHz with a 7 mm probe is not sufficient. 1.3 or 2.5 mm probes are not recommended because FSLG does not benefit from spinning very fast and ample amounts in these probes don't allow sufficient sensitivity. Further details of FSLG are explained in other chapters (e.g. *FSLG-HETCOR [▷ 105]*).

Start from a data set with weak adjusted cross polarisation and proton decoupling at fairly high RF field. Dependent on the probe and the magnetic field 100 to 125 kHz is recommended. At magnetic fields below 400 MHz 80 to 100 kHz may also be sufficient.

In most cases we recommend that you start with a well-known set-up sample, e.g. ^{13}C natural abundance (!) tyrosine hydrochloride is possible, but iso- leucine may be better because it contains CH_3 as well. But any other well crystalline organic compound with different kinds of carbon substitution will do, of course a short proton T_1 is always preferred at this stage

Spin the sample at any frequency between 10 and 15 kHz and tune the probe.

From the data set with the CP setup create a new dataset. Set **pulprog** sostapt. Type **ased** or click the pulse symbol in **eda**.

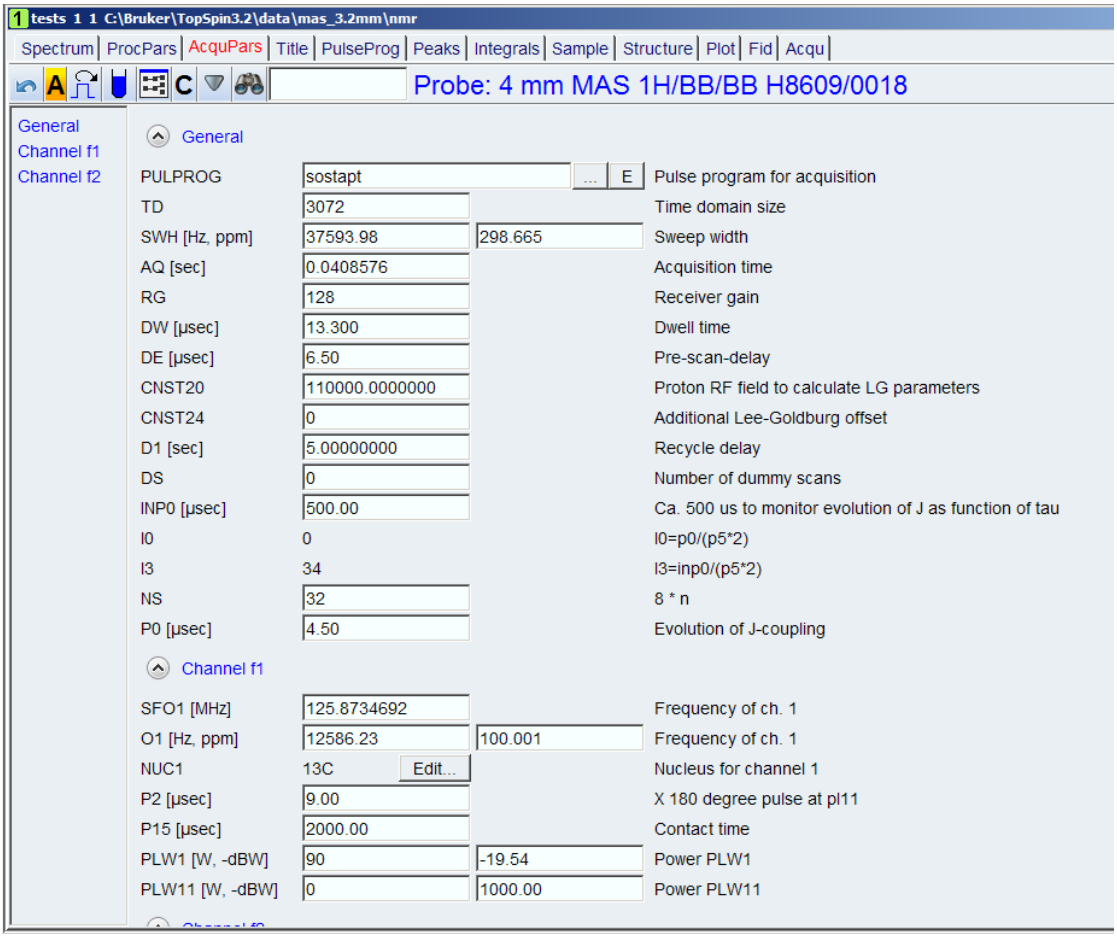

Figure 25.11: Display of general parameters in ased in the case of 1D

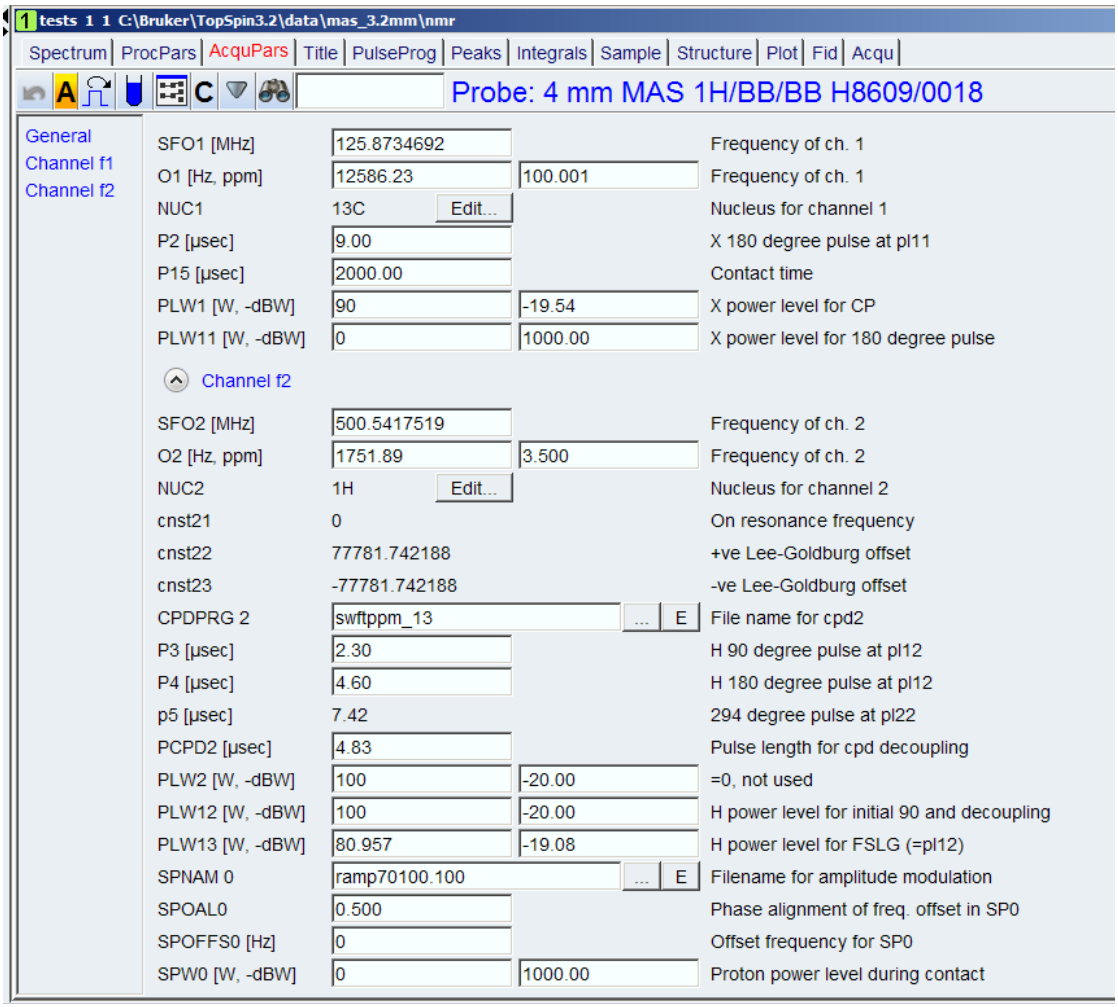

Figure 25.12: Display of channel dependent parameters in ased in the case of 1D data set

Parameter	Value	Comments
cnst11	0	= acqt0 (ref. baseopt)
cnst20	80000 to 120000	Proton RF field in kHz during FSLG
cnst31	10000 to 20000	MAS spin rate
Inf1	-	Only used for pseudo-2D
L0	= p0/(p5*2*l4)+1	Used to adjust τ delay
L3	0	Only used for pseudo-2D
L4	4 to 6	Must be set such that proton RF field calculated as cnst20 is 80 to 120 kHz
P0	2.5 to 7 ms	τ delay; depends on sample and which CHx need to be seen; see explanation below
f1 channel:		
P15	Typ. 2 ms	Contact time
P2	Typ. 8 µs	X 180° pulse
Pl1		X power level for CP
Pl11		X power level for 180°
f2 channel:		
Cpdprg2	Tppm15, spinal64, or swftppm	Cpd sequence for decoupling during observe
P3	Typ. 2.5 µs	H 90° pulse
P4	Typ. 5 µs	H 180° pulse
Cnst22/cnst23/p5		Calculated from pulse program
Pcpd2	Typ. 4.8 µs	Pulse length used in cpd sequence
Pl12		H power level for 90° and 180°
Pl13		H power level for FSLG, either =pl12 or recalculated
Sp0		H power level for CP
Spnam0	Ramp…	Shape used for protons during contact

Table 25.1: Acquisition parameters for sostapt sequence.

Acquisition parameters for sostapt sequence. Values for power levels need to be taken from the CP set-up and are left blank here because they depend on the available spectrometer hardware.

The idea of the experiment is now to run a series of experiments with different τ values and follow the signal intensities as τ is increased in order to monitor the different $\cos^n(\pi J\tau)$ for the different CH_n moieties. Note that signals from quaternary carbons will remain positive at all τ values (n=0) but decrease in intensity due to T_2. CH will show the slowest decay and eventually invert sign; CH_3 behave similarly but the decay is much faster and negative intensities show only at much longer τ delays. CH_2 intensities never invert. The strong coupling to 2 protons usually prevents signal intensities to come back due to relaxation. Furthermore, they are most sensitive to the adjustment of FSLG. This can be realized in such a way that a few 1D spectra are taken for which suitable τ values are set. An example of such a result is shown in the figure below. It can be seen that for a τ value of 4.5 ms the CH_3 signals are still null while it takes 6 ms to be back at to noticeable negative intensity. For CH there is less negative intensity when increasing τ value from 4.5 to 6 ms. Quaternary carbon signals remain at fairly large intensity because there is no dephasing because of J-coupling and additional relaxation from imperfections in FSLG set-up. Signals from CH_2 usually drop to 0 but don't manage to come back to positive intensity.

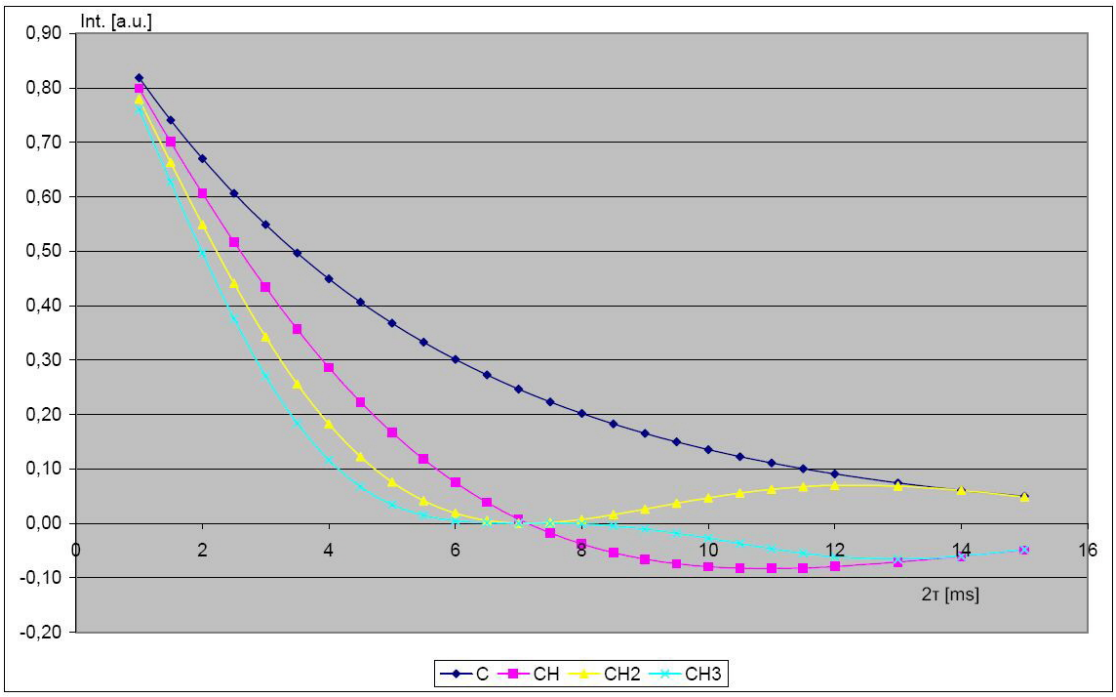

Figure 25.13: Relative signal intensities for C (blue), CH (pink), CH2 (yellow) and CH3 (turquoise) as function of 2τ value [ms]

Simulation is done for 100 Hz line width and 70 Hz effective coupling. Note that the CH_3 intensities remain very weak for 2τ delays between approx. 5.5 and 9 ms. CH_2 intensity comes back to noticeable values only for 2τ delays of more than 10 ms for which FSLG is very critical.

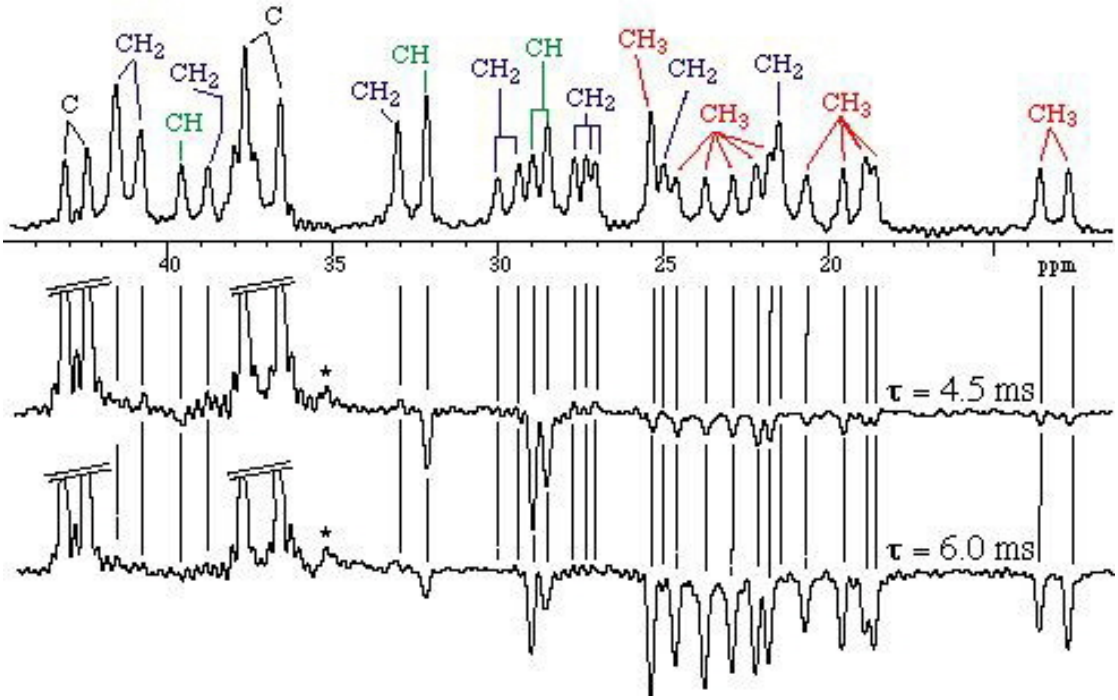

Figure 25.14: Assignment of 13C signals in cholesteryl acetate obtained with the parameters explained above

In the figure above T delays of 4.5 and 6 ms were used, respectively.

To run a pseudo-2D spectrum with a series of spectra with regularly increasing τ values you can proceed as follows: from the 1D set-up create a new data set. In **eda** change to 2D mode and set **FnMode** to QF. Type **ased** or click on the corresponding symbol.

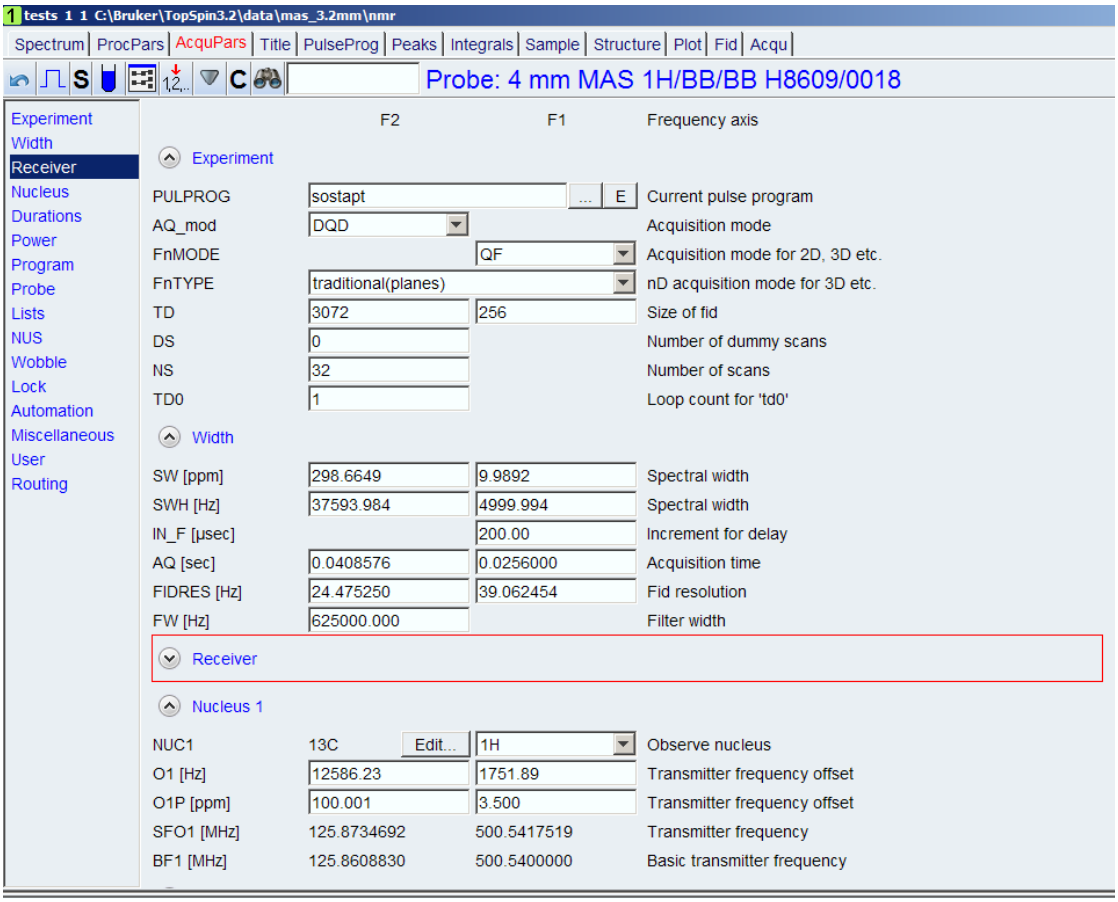

Figure 25.15: Display of eda for sostapt experiment in case of pseudo-2D experiment

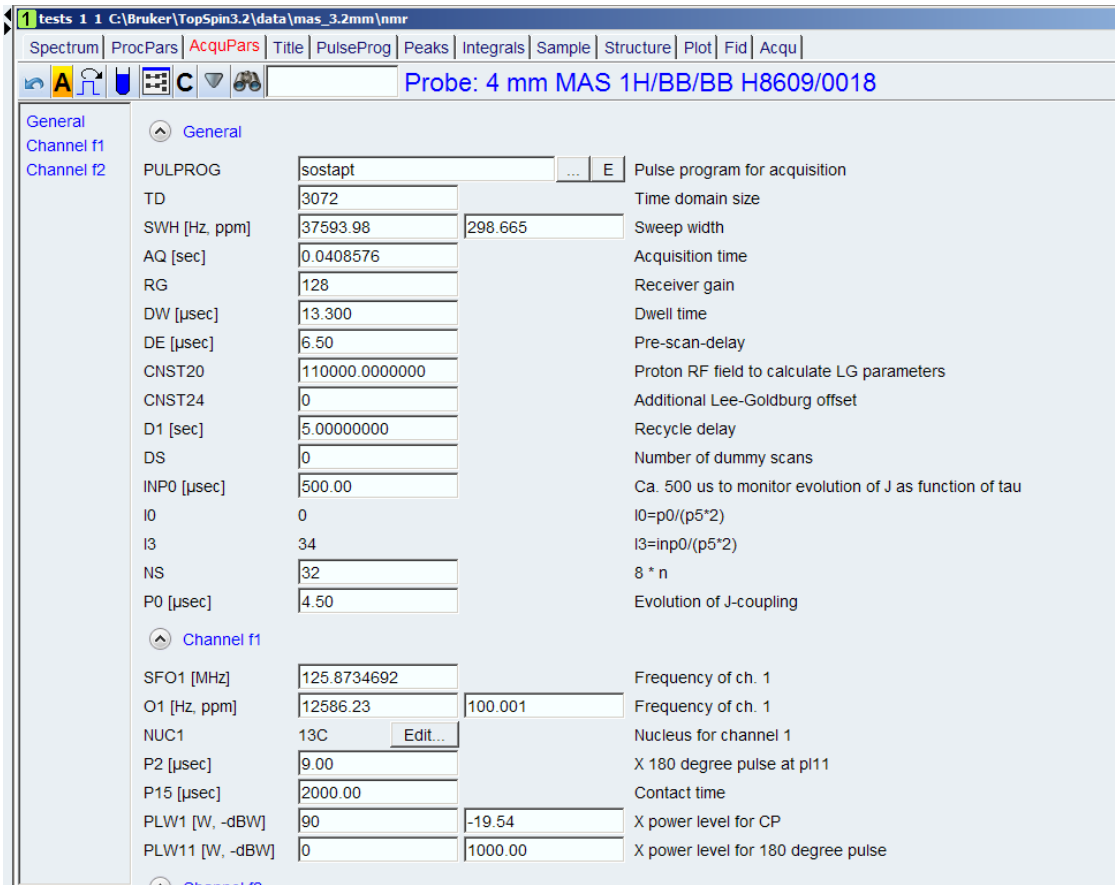

Figure 25.16: Display of general parameters in ased for the 2D setup

Parameter	Value	Comments
Inf1	= p5*4*l4*l0*l3	Calculated, total time increment is 2τ
L0	= p0/ (p5*2*l4)+1	Used to adjust τ delay
L3	1	Counter for τ increment in rotor periods
P0		Approx. τ, for time increment this value is multiplied by l3, make it 1 to 10 rotor periods (depends on spinning frequency).

Table 25.2: Additional acquisition parameters for sostapt sequence in case of a pseudo 2D experiment

26 Appendix

26.1 Form for Laboratory Logbooks

The form on the following pages may be printed and filled out by every user using the instrument to trace eventual problems and provide information for the next user.

Another copy may be printed for every user‘s own laboratory notebook. One form per probe used should be filled out. The following form serves as an example:

Operator	**used from:**		**to:**						
HF	10.12.07 10.00h		13.12.07 18.00h						
Probe:	**shim file:**		**B_0-field:**						
4mm triple C/N/H	triple4.hf		390 SR -360.14						
Sample:	**Experiment:**		**pulses**	**(us)**	**pl(n)**	**(dB/watt)**		**ok?**	**S/N:**
Glycine	CP S/N test	F1 ^{13}C	p90: contact: mix: else:	3.5 2m - -	pl11 pl1	3 3.5	150 120	y	100
		F2 ^{1}H	p90: contact: decouple: mix: else:	2.5 2m 4.6 - -	pl12 sp0 pl12	4 5	120 100		
		F3	p90: contact: decouple: mix: else	- - - 					

stored under filename: glycine-4 /opt/topspin/reference 1 1

comments: spinning at 10 kHz ok, mains pressure at 6 bars, linewidth α-C 50 Hz, O2= 1500,spinal decoupling

Operator	used from:	to:						
Probe:	shim file:	B_0-field:						
		SR:						
Sample:	**Experiment:**	**pulses**	**us**	**pl(n)**	**dB**	**watt**	**ok?**	**S/N:**
	F1	p90: contact: mix: else:						
	F2	p90: contact: decouple: mix: else:						
	F3	p90: contact: decouple: mix: else						
Sample:	**Experiment:**							

stored under filename:
comments:

2. Pulse program cpopt:

```
;cpopt (TopSpin 2.1)

;single pulse excitation, acquisition without decoupling

;Avance III version
;parameters:
;d1 : recycle delay
;p3 : proton excitation pulse length as in cp
;pl12 : decoupling/excitation power level for cp
;spnam0 : usual shape for cp
;sp0 : usual shape power for cp contact

;$COMMENT=single pulse excitation, acquisition without decoupling
;$CLASS=Solids
;$DIM=pseudo-2D
;$TYPE=direct excitation
;$SUBTYPE=relaxation measurement
;$OWNER=Bruker
;cnst11 : to adjust t=0 for acquisition, if digmod = baseopt
"acqt0=1u*cnst11"

1 ze
2 d1
  (p3 pl12 ph1):f1
  (p15:sp0 ph10):f1
  go=2 ph31
  wr #0
exit

ph1= 0 2
ph10= 1
ph31= 0 2
```

3. Power conversion table:

power conversion table				
probe: 4mm triple				
nucleus/frequency	p90 (us)	RF-field (kHz)	power (W)	remarks
^{1}H/				
^{19}F/				
^{15}N/				
^{15}N/				
^{29}Si/				
^{13}C/				
^{13}C/				
^{119}Sn/				
^{31}P/				

27 Contact

Manufacturer

Bruker BioSpin GmbH
Silberstreifen 4
D-76287 Rheinstetten
Germany
http://www.bruker.com

WEEE DE43181702

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